

Estimating Compressibility by Maximal-mass Compact Star Observations

Balázs E. Szigeti^{1 2}, G. G. Barnaföldi²

¹Eötvös Loránd University

²Wigner Research Centre for Physics

Eur.Phys.J.ST 229 (2020) 22-23, 3605-3614

Pub.Astr.Soc.Aust., Volume 35, e019, (2018)

Eur.Phys.J.ST 229, 22-23, 3615-3628 (2020)



Extended Walecka-model ($\sigma - \omega$ model)

- **Scalar meson** $U_i(\sigma) = \gamma_3\sigma^3 + \gamma_4\sigma^4$

- Electron kinetic and mass
- Kinetic terms: $\Psi = (\Psi_p, \Psi_n)$

$$\mathcal{L} = \bar{\Psi}(i\partial - m_N + g_\sigma\sigma - g_\omega\omega + g_\rho\rho^2\tau_a)\Psi + \bar{\Psi}_e(i\partial - m_e)\Psi_e + \frac{1}{2}\sigma(\partial^2 + m_\sigma^2)\sigma - U_i(\sigma) - \frac{1}{4}\omega^{\mu\nu}\omega^{\mu\nu} + \frac{1}{2}m_\omega^2\omega^\mu\omega_\mu - \frac{1}{4}\rho^{a\mu\nu}\rho_{\mu\nu}^a + \frac{1}{2}m_\rho^2\rho_\mu^a\rho_\mu^a$$

- Vector mesons

$$\omega_{\mu\nu} = \partial_\mu\omega_\nu - \partial_\nu\omega_\mu \quad \rho_{\mu\nu}^a = \partial_\mu\rho_\nu^a - \partial_\nu\rho_\mu^a + g_\rho\epsilon^{abc}\rho_\nu^b\rho_\mu^c$$



Extended Walecka-model ($\sigma - \omega$ model)

- Scalar meson $U_i(\sigma) = \gamma_3\sigma^3 + \gamma_4\sigma^4$
- **Electron kinetic and mass**
- Kinetic terms: $\Psi = (\Psi_p, \Psi_n)$

$$\mathcal{L} = \bar{\Psi}(i\cancel{\partial} - m_N + g_\sigma\sigma - g_\omega\cancel{\omega} + g_\rho\cancel{\rho}^2\tau_a)\Psi + \bar{\Psi}_e(i\cancel{\partial} - m_e)\Psi_e + \frac{1}{2}\sigma(\partial^2 + m_\sigma^2)\sigma - U_i(\sigma) - \frac{1}{4}\omega^{\mu\nu}\omega^{\mu\nu} + \frac{1}{2}m_\omega^2\omega^\mu\omega_\mu - \frac{1}{4}\rho^{a\mu\nu}\rho_{\mu\nu}^a + \frac{1}{2}m_\rho^2\rho_\mu^a\rho_\mu^a$$

- Vector mesons

$$\omega_{\mu\nu} = \partial_\mu\omega_\nu - \partial_\nu\omega_\mu \quad \rho_{\mu\nu}^a = \partial_\mu\rho_\nu^a - \partial_\nu\rho_\mu^a + g_\rho\epsilon^{abc}\rho_\nu^b\rho_\mu^c$$



Extended Walecka-model ($\sigma - \omega$ model)

- Scalar meson $U_i(\sigma) = \gamma_3\sigma^3 + \gamma_4\sigma^4$
- Electron kinetic and mass
- **Kinetic terms:** $\Psi = (\Psi_p, \Psi_n)$

$$\mathcal{L} = \bar{\Psi}(i\cancel{\partial} - m_N + g_\sigma\sigma - g_\omega\cancel{\omega} + g_\rho\cancel{\rho}^2\tau_a)\Psi + \bar{\Psi}_e(i\cancel{\partial} - m_e)\Psi_e + \frac{1}{2}\sigma(\partial^2 + m_\sigma^2)\sigma - U_i(\sigma) - \frac{1}{4}\omega^{\mu\nu}\omega^{\mu\nu} + \frac{1}{2}m_\omega^2\omega^\mu\omega_\mu - \frac{1}{4}\rho^{a\mu\nu}\rho_{\mu\nu}^a + \frac{1}{2}m_\rho^2\rho_\mu^a\rho_\mu^a$$

- Vector mesons

$$\omega_{\mu\nu} = \partial_\mu\omega_\nu - \partial_\nu\omega_\mu \quad \rho_{\mu\nu}^a = \partial_\mu\rho_\nu^a - \partial_\nu\rho_\mu^a + g_\rho\epsilon^{abc}\rho_\nu^b\rho_\mu^c$$



Extended Walecka-model ($\sigma - \omega$ model)

- Scalar meson $U_i(\sigma) = \gamma_3\sigma^3 + \gamma_4\sigma^4$
- Electron kinetic and mass
- Kinetic terms: $\Psi = (\Psi_p, \Psi_n)$

$$\mathcal{L} = \bar{\Psi}(i\cancel{\partial} - m_N + g_\sigma\sigma - g_\omega\cancel{\omega} + g_\rho\cancel{\rho}^2\tau_a)\Psi + \bar{\Psi}_e(i\cancel{\partial} - m_e)\Psi_e + \frac{1}{2}\sigma(\partial^2 + m_\sigma^2)\sigma - U_i(\sigma) - \frac{1}{4}\omega^{\mu\nu}\omega^{\mu\nu} + \frac{1}{2}m_\omega^2\omega^\mu\omega_\mu - \frac{1}{4}\rho^{a\mu\nu}\rho_{\mu\nu}^a + \frac{1}{2}m_\rho^2\rho_\mu^a\rho_\mu^a$$

- **Vector mesons**

$$\omega_{\mu\nu} = \partial_\mu\omega_\nu - \partial_\nu\omega_\mu \quad \rho_{\mu\nu}^a = \partial_\mu\rho_\nu^a - \partial_\nu\rho_\mu^a + g_\rho\epsilon^{abc}\rho_\nu^b\rho_\mu^c$$



Mean-Field Approximation

- Mean-field approximation at zero temperature and finite chemical potential
- Tree-level: $\omega = \omega_0$ and $\rho_0^3 = \rho$
- β -equilibrium: $\mu_n = \mu_p + \mu_e$
- Baryon number and electric charge conservation
- **Free-energy:**

$$f_T = f_F(m_N - g_\sigma\sigma, \mu_p - g_\omega\omega + g_\rho\rho) + f_F(m_N - g_\sigma\sigma, \mu_n - g_\omega\omega - g_\rho\rho) + f_f(m_e, \mu_e) + \frac{1}{2}m_\sigma^2\sigma^2 + U_i(\sigma) - \frac{1}{2}m_\omega^2\omega^2 - \frac{1}{2}m_\rho^2\rho^2 \quad (2)$$

- **Fermionic pressure:**

$$f_F(T, m, \mu) = -2T \int \frac{d^3k}{(2\pi)^3} \ln \left[1 + e^{-\beta(E_k - \bar{\mu})} \right]$$



Data Fit

The free parameters of the model are determined by using nuclear saturation data [Gle97].

The values used to fit the model

- (1) **Binding energy:** $B = -16.3 \text{ MeV}$
- (2) Saturation density: $n_0 = 0.153 \text{ fm}^{-3}$
- (3) Nucleon Landau mass $m_L = 0.8854 m_N$
- (4) Compressibility: $K = 240 \text{ MeV}$
- (5) Asymmetry energy: $a_{sym} = 32.5 \text{ MeV}$

$$K = 9n^2 \frac{\partial^2}{\partial n^2} \left(\frac{\epsilon}{n} \right) \quad m_L = \sqrt{k_f^2 + m_{N,eff}^2} \quad a_{sym} = \frac{1}{2} \frac{\partial^2}{\partial t^2} \left(\frac{\epsilon}{n} \right) \Big|_t$$



(4)

Data Fit

The free parameters of the model are determined by using nuclear saturation data [Gle97].

The values used to fit the model

- (1) Binding energy: $B = -16.3 \text{ MeV}$
- (2) **Saturation density:** $n_0 = 0.153 \text{ fm}^{-3}$
- (3) Nucleon Landau mass $m_L = 0.8854 m_N$
- (4) Compressibility: $K = 240 \text{ MeV}$
- (5) Asymmetry energy: $a_{sym} = 32.5 \text{ MeV}$

$$K = 9n^2 \frac{\partial^2}{\partial n^2} \left(\frac{\epsilon}{n} \right) \quad m_L = \sqrt{k_f^2 + m_{N,eff}^2} \quad a_{sym} = \frac{1}{2} \frac{\partial^2}{\partial t^2} \left(\frac{\epsilon}{n} \right) \Big|_t$$



(4)

Data Fit

The free parameters of the model are determined by using nuclear saturation data [Gle97].

The values used to fit the model

- (1) Binding energy: $B = -16.3 \text{ MeV}$
- (2) Saturation density: $n_0 = 0.153 \text{ fm}^{-3}$
- (3) **Nucleon Landau mass** $m_L = 0.8854 m_N$
- (4) Compressibility: $K = 240 \text{ MeV}$
- (5) Asymmetry energy: $a_{sym} = 32.5 \text{ MeV}$

$$K = 9n^2 \frac{\partial^2}{\partial n^2} \left(\frac{\epsilon}{n} \right) \quad m_L = \sqrt{k_f^2 + m_{N,eff}^2} \quad a_{sym} = \frac{1}{2} \frac{\partial^2}{\partial t^2} \left(\frac{\epsilon}{n} \right) \Big|_t$$



(4)

Data Fit

The free parameters of the model are determined by using nuclear saturation data [Gle97].

The values used to fit the model

- (1) Binding energy: $B = -16.3 \text{ MeV}$
- (2) Saturation density: $n_0 = 0.153 \text{ fm}^{-3}$
- (3) Nucleon Landau mass $m_L = 0.8854 m_N$
- (4) **Compressibility:** $K = 240 \text{ MeV}$
- (5) Asymmetry energy: $a_{sym} = 32.5 \text{ MeV}$

$$K = 9n^2 \frac{\partial^2}{\partial n^2} \left(\frac{\epsilon}{n} \right)$$

$$m_L = \sqrt{k_f^2 + m_{N,eff}^2}$$

$$a_{sym} = \frac{1}{2} \frac{\partial^2}{\partial t^2} \left(\frac{\epsilon}{n} \right) \Big|_t$$



(4)

Data Fit

The free parameters of the model are determined by using nuclear saturation data [Gle97].

The values used to fit the model

- (1) Binding energy: $B = -16.3 \text{ MeV}$
- (2) Saturation density: $n_0 = 0.153 \text{ fm}^{-3}$
- (3) Nucleon Landau mass $m_L = 0.8854 m_N$
- (4) Compressibility: $K = 240 \text{ MeV}$
- (5) Asymmetry energy: $a_{sym} = 32.5 \text{ MeV}$

$$K = 9n^2 \frac{\partial^2}{\partial n^2} \left(\frac{\epsilon}{n} \right) \quad m_L = \sqrt{k_f^2 + m_{N,eff}^2} \quad a_{sym} = \frac{1}{2} \frac{\partial^2}{\partial t^2} \left(\frac{\epsilon}{n} \right) \Big|_t$$



(4)

Data Fit

The free parameters of the model are determined by using nuclear saturation data [Gle97].

The values used to fit the model

- (1) Binding energy: $B = -16.3 \text{ MeV}$
- (2) Saturation density: $n_0 = 0.153 \text{ fm}^{-3}$
- (3) Nucleon Landau mass $m_L = 0.8854 m_N$
- (4) Compressibility: $K = 240 \text{ MeV}$
- (5) **Asymmetry energy: $a_{sym} = 32.5 \text{ MeV}$**

$$K = 9n^2 \frac{\partial^2}{\partial n^2} \left(\frac{\epsilon}{n} \right) \quad m_L = \sqrt{k_f^2 + m_{N,eff}^2} \quad a_{sym} = \frac{1}{2} \frac{\partial^2}{\partial t^2} \left(\frac{\epsilon}{n} \right) \Big|_t$$



(4)

Mass-Radius Diagrams

- Static, spherically symmetric space-time
- Tolman-Oppenheimer-Volkoff equations (TOV)

$$\frac{dp(r)}{dr} = -\frac{G\epsilon(r)m(r)}{r^2} \left[1 + \frac{p(r)}{\epsilon(r)} \right] \left[1 + \frac{4\pi r^3 p(r)}{m(r)} \right] \left[1 - \frac{2Gm(r)}{r} \right]^{-1} \quad (5)$$
$$\frac{dm(r)}{dr} = 4\pi r^2 \epsilon(r)$$

- We used core-crust model for the calculation
- Low density region: well known BPS nuclear EoS
- The calculated R in this case corresponds to the radius of the neutron star core



Maximal Mass Stars

- Determine the mass and radius of the maximal mass stars (MMS)
- Most relevant parameter: the Landau mass m_L . ($\delta m_L^{10x} > \delta K$)
- Linear dependence M_{maxM}, R_{MaxM} on the m_L
- Landau mass was optimized with fixed K, a_{sym} .

$$M_{maxM} [M_{\odot}] = 5.418 - 0.0043 m_L [\text{MeV}] ,$$
$$R_{maxM} [\text{km}] = 19.04 - 0.0104 m_L [\text{MeV}] .$$

After fixing the m_L by the MMS observations we can also obtain linear, one-parameter dependence on the parameter K .

$$M_{maxM} [M_{\odot}] = 1.940 + 0.000880 K [\text{MeV}] ,$$
$$R_{maxM} [\text{km}] = 9.248 + 0.00718 K [\text{MeV}] .$$



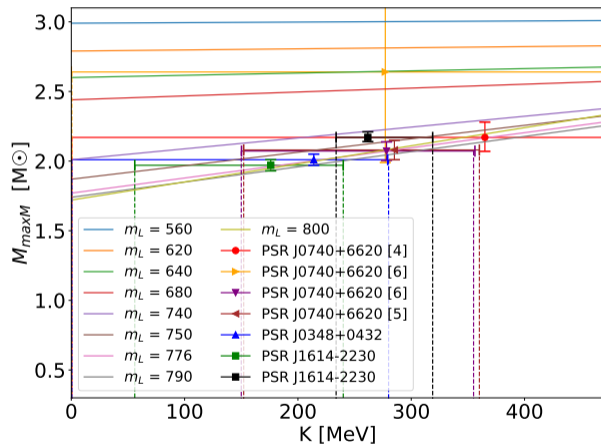
Observation Data

Ref.	Pulsar	$M_{maxM} [M_{\odot}]$	$m_L [\text{MeV}]$	$K [\text{MeV}]$	$R_{maxM} [\text{km}]$
[J0740]	PSR J0740+6620	$2.17^{+0.11}_{-0.10} *$	$748.39^{+63.3}_{-57.2}$	$351.8^{+115}_{-84.5}$	$11.25^{+1.06}_{-1.04}$
[J0348]	PSR J0348+0432	$2.01^{+0.04}_{-0.04} *$	$785.25^{+20.0}_{-20.3}$	$206.4^{+42.7}_{-20.5}$	$10.87^{+0.82}_{-0.80}$
[J1614]	PSR J1614-2230	$1.97^{+0.04}_{-0.04} *$	$794.47^{+20.1}_{-20.4}$	$170.0^{+15.5}_{-20.9}$	$10.77^{+0.82}_{-0.80}$
[NICER]	PSR J0740+6620	$2.07^{+0.07}_{-0.07} *$	$778.14^{+15.3}_{-15.5}$	$278.2^{+60.9}_{-60.8}$	$11.01^{+0.46}_{-0.47}$
[NICER]	PSR J0740+6620	$2.64^{+1.98}_{-0.65}$	639.42^{+159}_{-125}	277.3^{+257}_{-178}	$12.39^{+1.30}_{-0.98} *$
[XMNS]	PSR J0740+6620	$2.08^{+0.07}_{-0.01} *$	$769.12^{+16.9}_{-16.9}$	$285.1^{+54.8}_{-54.8}$	$11.06^{+0.41}_{-0.41}$

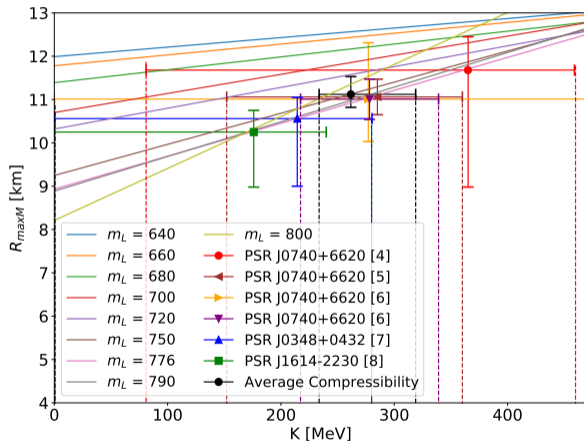
Table 1: The Landau mass, m_L and compressibility, K values calculated via eq. (??)- from measured pulsar mass data, assuming that these are maximal-mass neutron stars. Maximal radius of the maximal-mass neutron star is also calculated form eq. (??). [Ant+13; al20; Dem10]



K dependence of the mass of MMS



K dependence of the radius of MMS



Landau mass values:

Remark I.

The range of $550 \text{ MeV} < m_L < 800 \text{ MeV}$ is compatible with effective nucleon mass given by the various equation of state at saturation [Kl +06].

Remark II.

- Extended Walecka Model: $m_L = 776.0^{+38.5}_{-84.9} \text{ MeV}$
- Bayesian Analysis I. [Alv+20a]: (PSR J0030+0451 + PSR 0740+6620 + GW170817)
 - $m_L = 750^{+15}_{-15} \text{ MeV}$
- Bayesian Analysis II. [Alv+20b]: (PSR J0740+6620 + PSR J0348+0432 + GW170817 + PSR J0030+0451)
 - $m_L = 750^{+15}_{-15} \text{ MeV}$



Summary

Input:

- Extended Walecka Model
- Mean-Field Approximation

Method:

- Tollman-Oppenheimer-Volkoff equation
- Using the linear relations between the microscopic and macroscopic parameters
- Observation data

Results:

- Landau mass: $m_L = 776.6_{-84.2}^{+40.1}$ MeV
- Compressibility: $K = 287.3_{-29.2}^{+59.1}$ MeV



References

Thank you for your time. Questions?

- [al20] Cromatie H. T. Fonseca E. Ransom S.M. et al. “Relativistic Shapiro delay measurements of an extremely massive millisecond pulsar”. Í: *Nature Astronomy* 4.72–76 (2020).
- [Alv+20a] David Alvarez-Castillo o.fl. “Studying the Landau mass parameter of the extended sigma-model for neutron star matter”. Í: *Physics of Particles and Nuclei* 51 (apr. 2020), bls. 725–729.
- [Alv+20b] David Alvarez-Castillo o.fl. “Studying the parameters of the extended $\sigma - \omega$ model for neutron star matter”. Í: *Eur. Phys. J. Special Topics* 229 (2020), bls. 3615–3628.
- [Ant+13] John Antoniadis o.fl. “A Massive Pulsar in a Compact Relativistic Binary”. Í: *Science* 340.6131 (2013). ISSN: 0036-8075. DOI: 10.1126/science.1233232. eprint: <https://science.sciencemag.org/content/340/6131/1233232.full.pdf>. URL: <https://science.sciencemag.org/content/340/6131/1233232>.



References (cont.)

- [Dem10] Ransom S. et al Demorest P. Pennucci T. “A two-solar-mass neutron star measured using Shapiro delay”. Í: *Nature* 467 (2010), bls. 1081–1083.
- [Gle97] N. K. Glendenning. *Compact stars: Nuclear physics, particle physics, and general relativity*. 1997.
- [Klä+06] T. Klähn o.fl. “Constraints on the high-density nuclear equation of state from the phenomenology of compact stars and heavy-ion collisions”. Í: *Phys. Rev. C* 74 (3 sep. 2006), bls. 035802. DOI: 10.1103/PhysRevC.74.035802. URL: <https://link.aps.org/doi/10.1103/PhysRevC.74.035802>.

