## Measurement of electric charge dependent splitting of directed flow in STAR experiment at RHIC

Részecskefizika Seminar **Department Of High Energy Physics, Wigner Research Centre for Physics** Hungarian Academy of Sciences

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STAR

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## Quark Gluon Plasma (QGP)



- $\bullet$  Later on the universe began to cool down then nucleons, low mass light nuclei and eventually the matter around us are formed
- (QGP)
- To understand early universe => Need to study QGP

• A few microseconds after BigBang, the universe consisted of a hot soup of quarks and gluons

• The deconfined state of quarks and gluons, existed in the early universe - Quark-Gluon Plasma

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## How to create QGP?

- Quarks and gluons are confined inside hadrons How do we deconfine them?
- Certain conditions like extremely high pressure (P) and temperature (T) can deconfine them



- Under such conditions, the hadrons get close to each other so that quarks and gluons inside them can fly around freely in an extended volume
- Hadrons to QQP transition takes place at T ~ 10^4 times temperature of the core of the sun
- Big particle collider produces such extremely high T for a short period of time by colliding ions at relativistic energies







## **Relativistic colliders**

## **Relativistic Heavy Ion Collider (RHIC) at BNL**



Energy/(proton mass) ~ 500 Colliding energy ~ Few hundreds of GeV

Ideal places to study QGP

## Large Hadron Collider (LHC) at CERN

Energy/(proton mass) ~ 14000 **Colliding energy ~ Few TeV** 





## **Relativistic Heavy Ion Collider**

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## The future (>2025) : The Electron Ion collider



P Tribedy, Rutgers Nuclear Physics Seminars, 2018

## The future (>2025) : The Electron Ion collider





## )n Ior

## **Relativistic Heavy Ion Collider**





## The future (>2025) : The Electron Ion collider





## )n Ior

## Heavy Ion Collisions : QGP at the Lab

- RHIC circulates heavy nuclei (Au) almost at speed of light and smash them together
- Shortly after the collisions (~ 1 fm/c), an enormous amount of energy is released into a tiny volume
- As a result, a QGP medium possibly form
- The QGP expands and cools down rapidly, and when T falls below a critical T, quarks and gluons form bound states - Hadronization
- Then particles get captured in the detectors placed around collision point





## Event display of Au+Au collisions at STAR detector



## Au+Au 200 GeV Event# 1007 Run# 17172038



Ashik Ikbal, Particle Physics Seminar, Wigner RCP, Budapest



Au

## What do we want to know from Au+Au collisions ?



Au+Au 200 GeV Event# 1007 Run# 17172038







Ashik Ikbal, Particle Physics Seminar, Wigner RCP, Budapest



Au

(100 GeV/A)

particles



# Something interesting in particle emission pattern



• Particle emission w.r.t transverse plane of collisions in not isotropic

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# Something interesting in particle emission pattern



resembles an ellipse

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## Less particles







Initial state collision geometry is like an almond in position space

- Large internal pressure along minor axis than major axis of the position space ellipse => momentum space ellipse
- Position space anisotropy translated into momentum space anisotropy Elliptic flow

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## **Preferred particle emission in one direction**



- Particle emission might be emitted preferably in one direction
- flow



• This anisotropy describes sideward motion of the emitted particles - One directional => Directed

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## Characterization of anisotropic flow (v<sub>n</sub>)

- Angular distribution is characterized by Fourier series
- Express particle angular distributions in Fourier series,

$$E\frac{d^3N}{dp^3} = \frac{d^2N}{2\pi p_T dp_T dy} \Big(1 + 2\sum_{n=1}^{+\infty} v_n \cos n\Big)$$

where  $v_n = \langle \cos n(\phi - \Psi_{RP}) \rangle$ 

 $\Psi_{RP} \rightarrow Reaction Plane angle$ XZ -> Reaction Plane (RP)

- $V_n \longrightarrow$  flow harmonics
- $V_1 \longrightarrow$  directed flow,  $V_2 \longrightarrow$  elliptic flow,  $V_3 \longrightarrow$  triangular flow, and so on ...



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## What if there is non central collisions?



**Spectators** 

- Non-central HICs (Non-zero impact parameter)
- Charged spectator nuclei produce electric currents (like two parallel current carrying wires in opposite directions
- The currents can produce magnetic fields
- Magnetic fields due to these two sources add up



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## Estimates of the produced magnetic field

## A crude estimate of the magnetic field (using Biot-Savart Law): A crude estimate of the magnetic field (using Biot-Savart Law): A crude estimate of the magnetic field (using Biot-Savart Law): A crude estimate of the magnetic field (using Biot-Savart Law): A crude estimate of the magnetic field (using Biot-Savart Law): A crude estimate of the magnetic field (using Biot-Savart Law): A crude estimate of the magnetic field (using Biot-Savart Law): A crude estimate of the magnetic field (using Biot-Savart Law): A crude estimate of the magnetic field (using Biot-Savart Law):

 $-eB_y \sim 40m_\pi^2 \sim 10^{18}$  Gauss (At RHIC Au+Au collisions,  $\sqrt{s_{NN}} = 200$ GeV, b = 5 fm, t = 0)

- Strongest magnetic field ever produced in the universe
- Field has observable effects on properties of produced particles, such as anisotropic flow





Earth ~0.5 Gauss

STAR magnet ~5000 Gauss





Neutron Star (Magentar) ~ 10<sup>14</sup> Gauss



Heavy ion collisions ~ 10<sup>18</sup> Gauss



# Directed flow $(v_1)$ and splitting $(\Delta v_1)$

Directed flow (v<sub>1</sub>) describes sideward motion of particles and can be measured with particle rapidity



- Splitting:  $\Delta v_1 = v_1 (q1, S1) v_1 (q2, s2)$



## $\odot$ EM field can have observable consequences on v<sub>1</sub> splitting between charged particles

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## EM field drives splitting



- Assume a non-central HIC  $(b \neq 0)$
- ${\ensuremath{\, \bullet }}$  Beam direction:  $\hat{z}$  , Impact parameter:  $\hat{x}$
- Reaction plane (RP): xz
- Or Charged spectators produce magnetic
   Ansatz in the second field -  $\vec{B} \perp RP$

![](_page_17_Figure_6.jpeg)

![](_page_17_Figure_10.jpeg)

![](_page_17_Figure_11.jpeg)

# EM field drives splitting - Hall effect

Hall Effect -1 ⊙≡Â

![](_page_18_Figure_2.jpeg)

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![](_page_18_Picture_3.jpeg)

- Lorentz force pushes positively and negatively charged particles in opposite directions
- Generated current  $\perp B, \vec{u}$ => Hall effect

![](_page_18_Picture_9.jpeg)

# EM field drives splitting - Faraday and Coulomb effect

- Spectators fly away,  $\vec{B}$  decays down fast
- Time varying  $\vec{B}$  induces  $\vec{E}$  field => 0 Faraday effect
- Charged spectators also generate 0 **Coulomb field**

![](_page_19_Picture_4.jpeg)

# EM field drives splitting - Hall, Faraday and Coulomb effect

![](_page_20_Figure_1.jpeg)

- Faraday, Coulomb and Hall are competing effects
- Net effect of Faraday, Hall and Coulomb affects v<sub>1</sub> and splitting between particles and antiparticles
- Direction of  $v_1$  for positive particles shown by dashed arrows (when Faraday+Coulomb > Hall)
- $\bullet$  Direction of v<sub>1</sub> for negative particles the other way around
- EM field drives  $v_1$  splitting ( $\Delta v_1$ ) between particles and anti-particles
- Output Can we measure this splitting?

![](_page_20_Figure_10.jpeg)

![](_page_20_Figure_11.jpeg)

![](_page_20_Figure_12.jpeg)

![](_page_20_Picture_13.jpeg)

# Splitting ( $\Delta v_1$ ): Challenge in measurements (Transport)

- The u, d quarks can be transported from beam rapidity
- Transported quarks suffer a lot more interactions than produced quarks
- Transported quarks have different v<sub>1</sub> than produced quarks
- There is already a  $v_1$  splitting between quarks (transported) and anti-quarks (produced)
- This splitting interferes with the EM field driven splitting becomes difficult to isolate

![](_page_21_Picture_6.jpeg)

# Splitting ( $\Delta v_1$ ): Interplay between transport and EM field

![](_page_22_Figure_1.jpeg)

- This  $\Delta v_1$ -slope difference between proton and anti-proton is negative (Hydro+EMF expectation)
- These two effects convolute Overall effect changes the sign
- This splitting acts as a background effect for EM-field-driven splitting This background should be subtracted

 $\bullet \Delta v_1$ -slope difference between transported proton and anti-proton is positive (UrQMD expectation)

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![](_page_22_Picture_9.jpeg)

# Splitting ( $\Delta v_1$ ): An approach to subtract transported quark effect

- In experiment, it is impossible to distinguish between produced and transported u and d quarks
- Avoid particles containing u, d quarks
- Use only produced particles (only produced constituent quarks)  $= ar{u}, ar{d}, ar{s}, ar{s}$   $= K^-, ar{p}, ar{\Lambda}, \phi, ar{\Xi}^+, \Omega^- ext{ and } ar{\Omega}^+$
- With these particles, make a clean case to measure EM fielddriven-splitting
- Output Compare the combinations with same mass at the constituent level
- Apply and test Coalescence-inspired sum same  $y - p_T/n_q$  space, with  $n_q \rightarrow$  cons Constituent quarks  $q_i$  –

A. Ikbal, D. Keane, P. Tribedy, Phys. Rev. C 105, 014912 (2022)

rule: 
$$v_1(hadron) = \sum v_1^i(q_i)$$
,  
stituent quarks )

![](_page_23_Figure_9.jpeg)

![](_page_23_Picture_12.jpeg)

![](_page_23_Picture_13.jpeg)

![](_page_23_Picture_14.jpeg)

# Splitting ( $\Delta v_1$ ): Testing Coalescence sum rule

## Combine particles and make identical quark combinations

![](_page_24_Figure_2.jpeg)

• Charge difference,  $\Delta q = 0$  and strangeness difference,  $\Delta S = 0$ 

A. Ikbal, D. Keane, P. Tribedy, Phys. Rev. C 105, 014912 (2022)

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![](_page_24_Picture_8.jpeg)

# Splitting: Combination with non-zero $\Delta q$ and $\Delta S$

Combine particles and make non-identical quark combinations, same mass at the constituent level

![](_page_25_Figure_2.jpeg)

• Charge difference,  $\Delta q = 4/3$  and strangeness difference,  $\Delta S = 2$ 

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![](_page_25_Picture_8.jpeg)

## Making more combinations

• Combinations having same or nearly same quark mass but different  $\Delta q$  and  $\Delta S =>$ No transported quark effect

Index	Quark Mass	Charge	Strangeness	Expression
1	$\Delta m = 0$	$\Delta q = 0$	$\Delta S = 0$	$[\bar{p}(\bar{u}\bar{u}\bar{d}) + \phi(s\bar{s})] - [K(\bar{u}s) + \bar{\Lambda}(\bar{u}d\bar{s})]$
2	$\Delta m pprox 0$	$\Delta q = 1$	$\Delta S = 2$	$\left[\bar{\Lambda}(\bar{u}\bar{d}\bar{s})\right] - \left[\frac{1}{3}\Omega^{-}(sss) + \frac{2}{3}\bar{p}(\bar{u}\bar{u}\bar{d})\right]$
3	$\Delta m pprox 0$	$\Delta q = rac{4}{3}$	$\Delta S = 2$	$\left[\overline{\Lambda}(\overline{u}\overline{d}\overline{s})\right] - \left[K(\overline{u}s) + \frac{1}{3}\overline{p}(\overline{u}\overline{u}\overline{d})\right]$
4	$\Delta m = 0$	$\Delta q = 2$	$\Delta S = 6$	$\left[\overline{\Omega}^+(\overline{s}\overline{s}\overline{s}\overline{s}) ight] - \left[\Omega^-(sss) ight]$
5	$\Delta m pprox 0$	$\Delta q = \frac{7}{3}$	$\Delta S = 4$	$[\overline{\Xi}^+(\overline{d}\overline{s}\overline{s})] - [K(\overline{u}s) + \frac{1}{3}\Omega(sss)]$

Only 5 combination differences among many are independent

• Two degenerate combinations in  $\Delta S = 2$  - Good cross check

• Measure splitting with  $\Delta q$  and  $\Delta S$ , though they are correlated

A. Ikbal, D. Keane, P. Tribedy, Phys. Rev. C 105, 014912 (2022)

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![](_page_26_Picture_11.jpeg)

![](_page_26_Picture_12.jpeg)

## **Towards measurements: STAR detector and datasets**

• TPC+TOF for PID: TPC measures dE/dx of tracks ( $|\eta| < 1$ ,  $0 < \phi < 2\pi$ ) and TOF measures time of flight ( $|\eta| < 0.9$ )

• EPD ( $2.1 < |\eta| < 5.1$ ) or ZDC ( $|\eta| > 6.3$ ) for event plane reconstruction

## **Datasets analyzed:**

• At  $\sqrt{s_{NN}} = 27$  GeV Au+Au at BES-II, and  $\sqrt{s_{NN}} = 200 \text{ GeV Au+Au}$ collisions

![](_page_27_Picture_5.jpeg)

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![](_page_27_Picture_8.jpeg)

## **Coalescence sum rule at Au+Au @ 27 GeV**

![](_page_28_Figure_2.jpeg)

Sum rule with identical quark combinations

- $\Delta v_1$  slope (with y) ~ 10^{-4}
- Sum rule holds within measured uncertainties

 $v_1[K(\bar{u}s)] + v_1[\bar{\Lambda}(\bar{u}s\bar{d})] \stackrel{?}{=} v_1[\bar{p}(\bar{u}\bar{u}\bar{d})] + v_1[\phi(s\bar{s})]$ 

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![](_page_28_Picture_10.jpeg)

## Splitting at non-zero $\Delta q$ and $\Delta S$ (27 GeV)

![](_page_29_Figure_2.jpeg)

- $\Delta v_1$  increases at larger y and pT/ng
- Significant non-zero slope (with y) for  $\Delta q = 4/3$ ,  $\Delta S = 2$
- AMPT has the opposite trend No EM field in AMPT

## $v_1[K(\overline{u}s)] + \frac{1}{3}v_1[\overline{p}(\overline{u}\overline{u}d)]$

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![](_page_29_Picture_10.jpeg)

![](_page_29_Picture_11.jpeg)

## **Splitting with charge and strangeness**

![](_page_30_Figure_1.jpeg)

- $\circ \Delta v_1$  slope (fit constrained to origin) increases with  $\Delta q$  and ΔS
  - Splitting increases going from  $\sqrt{s_{NN}} = 200$  to 27 GeV
- AMPT can not explain the data (Nayak et al., Phys. Rev. C 100, 054903 (2019))
  - PHSD(+EMF) can describe the data within errors, but EMF is not the sole difference between these two models

![](_page_30_Figure_8.jpeg)

![](_page_30_Figure_9.jpeg)

![](_page_30_Figure_10.jpeg)

![](_page_30_Figure_11.jpeg)

![](_page_30_Picture_12.jpeg)

## Summary

- Heavy ion collisions are the tool to create QGP primordial matter
- Strong EM field is produced in non-central collisions and can affect the directed flow splitting
- directed flow splitting - free from the transported quark effect
- uncertainties
- Produced EM field can lead to the splitting

![](_page_31_Picture_7.jpeg)

Discussed how to measure charge ( $\Delta q$ ) and strangeness ( $\Delta S$ ) dependent

• Measured splitting increases with  $\Delta q$  and  $\Delta S$ , stronger in lower collision energy

PHSD+EM field calculations can describe the charge-dependent splitting within

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![](_page_31_Picture_16.jpeg)