

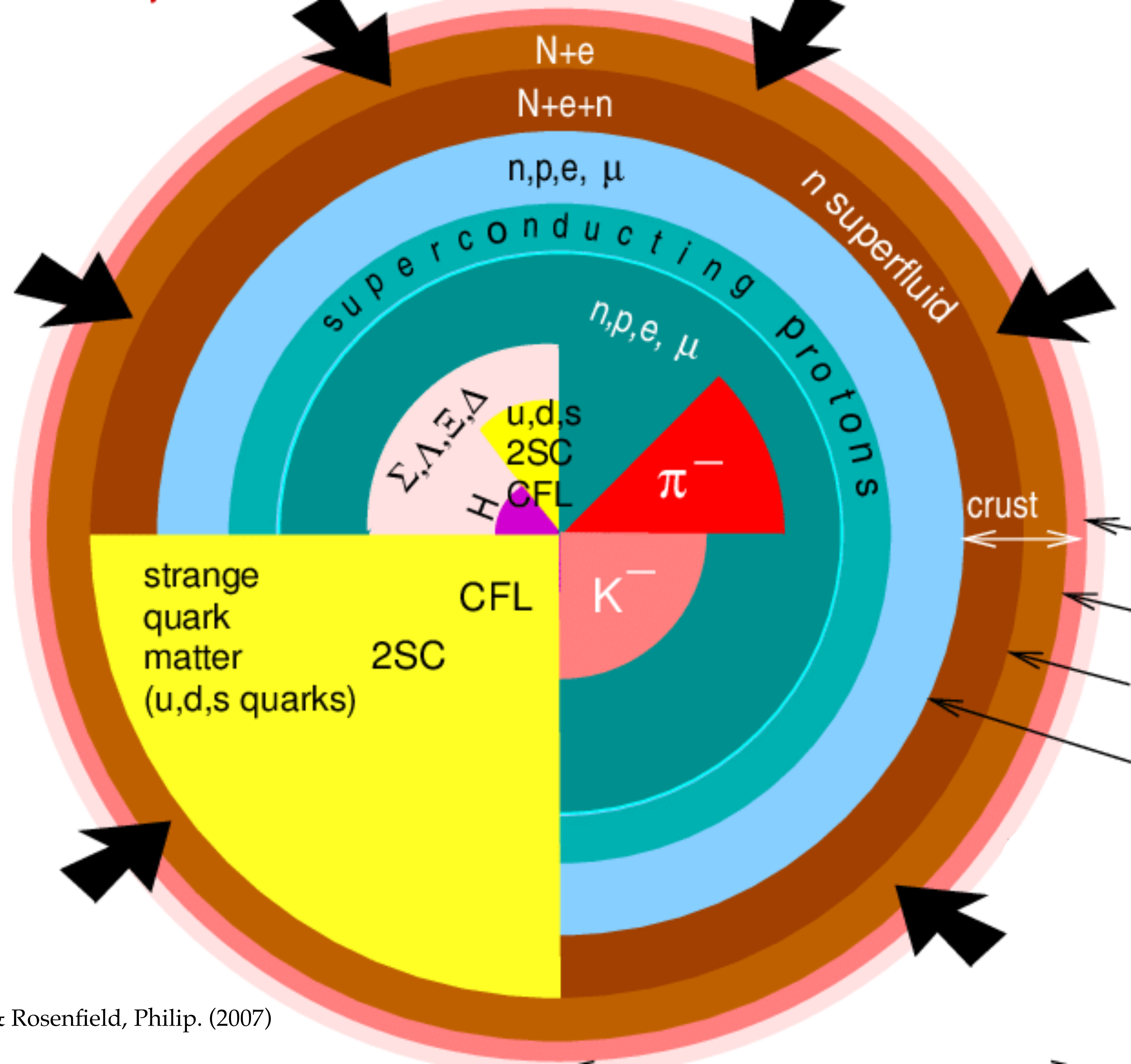
Universal relations in astrophysics

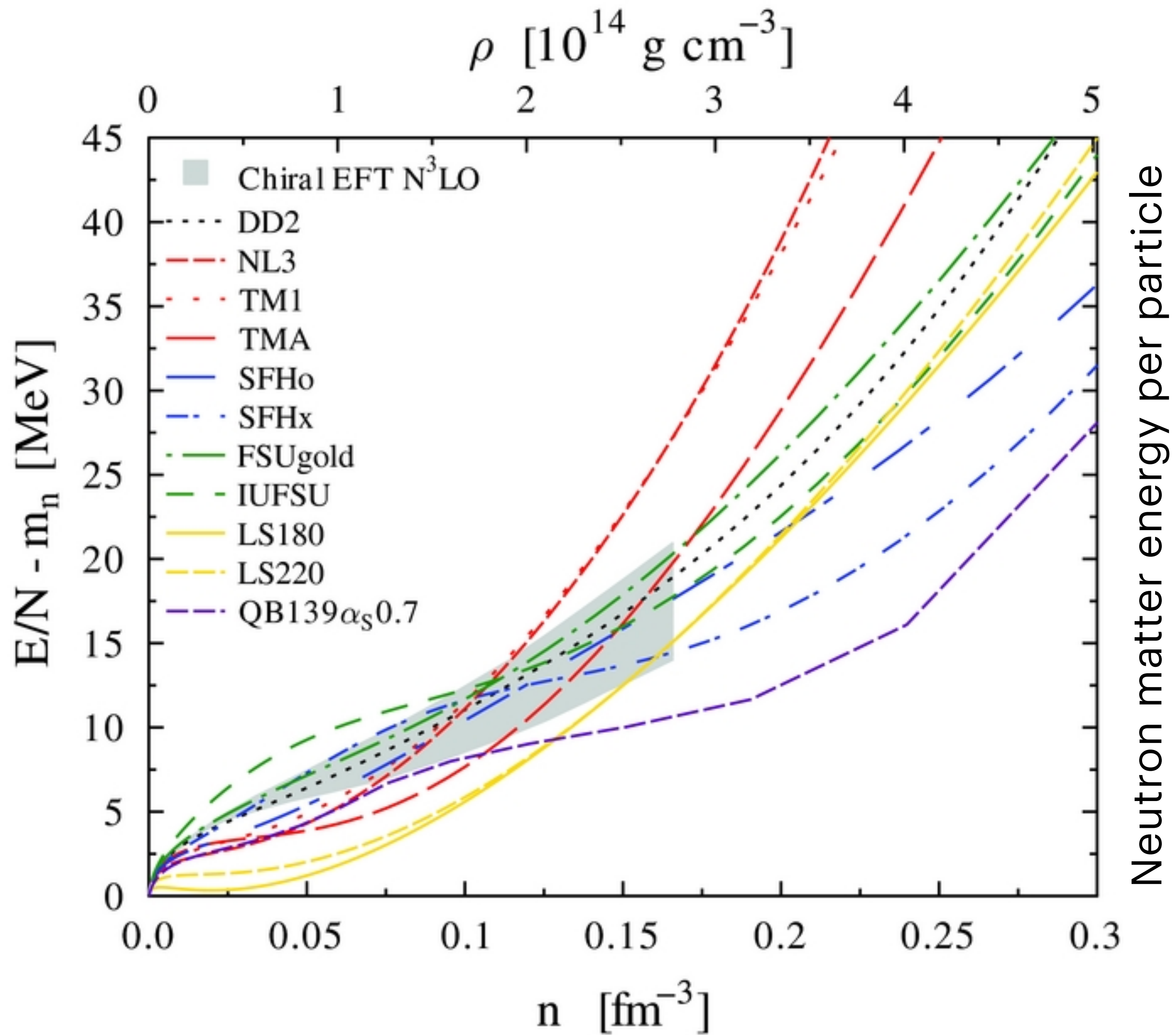
Wigner institute seminar, Budapest

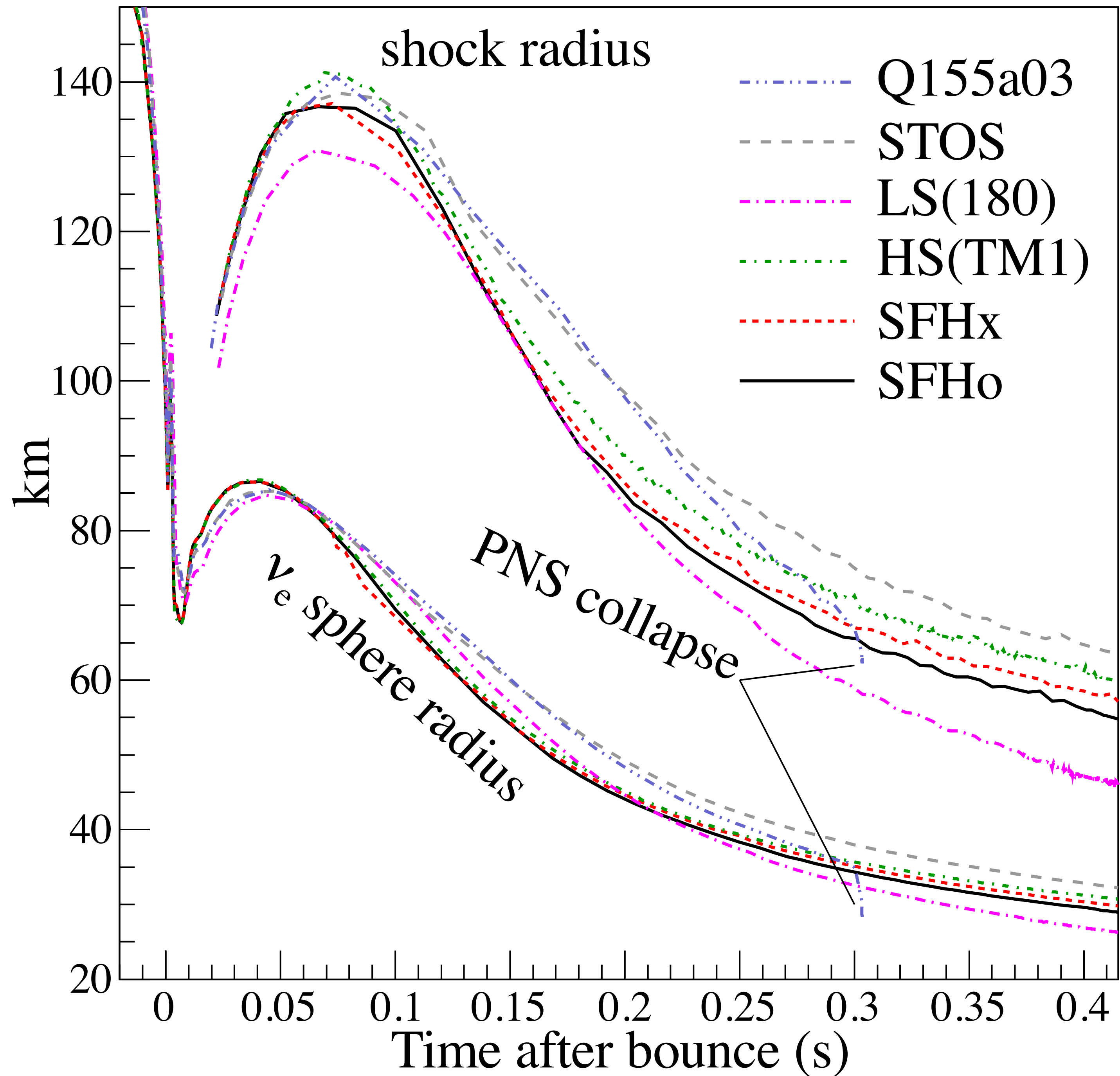
Noshad Khosravi Largani
University of Wroclaw
17 June 2022

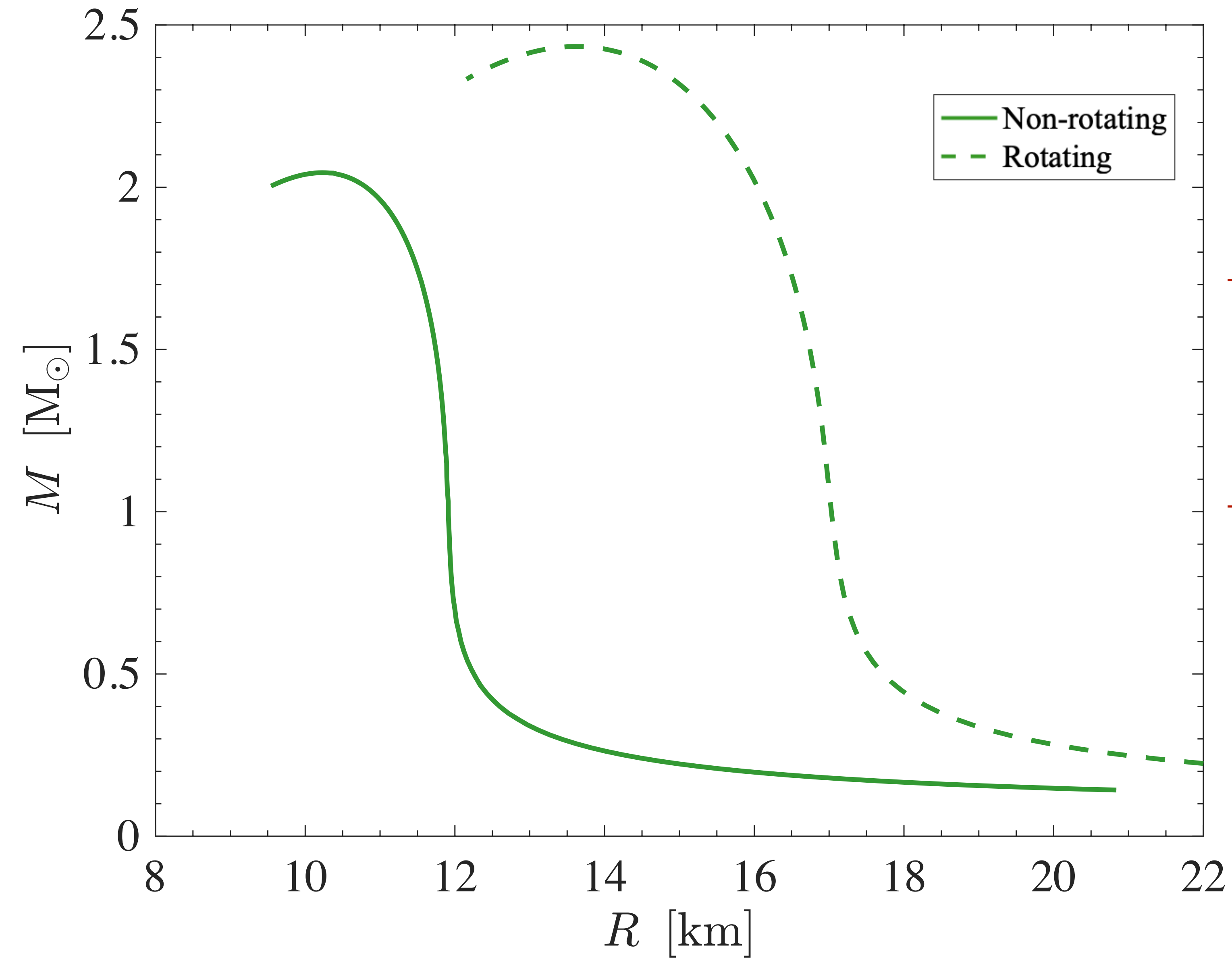
1. Rapidly Rotating Neutron Stars

arXiv: 2112.10439v1









$$\frac{dp}{dr} = - \frac{(\epsilon + p/c^2)G(m + 4\pi r^3 p/c^2)}{r^2(1 - 2Gm/rc^2)}$$

$$\frac{dm}{dr} = 4\pi r^2 \epsilon$$

$$x_p = x_e$$

$$\mu_n - \mu_p = \mu_e = \mu_{\mu}$$

$$p(\epsilon)$$

$$ds^2 = -e^{\gamma+\rho} dt^2 + e^{2\alpha} (dr^2 + r^2 d\theta^2) + e^{\gamma-\rho} r^2 \sin^2 \theta (d\phi - \omega dt)^2$$

$$u^{\mu} = \frac{e^{-(\gamma+\rho)}}{\Gamma} (1, 0, 0, \Omega)$$

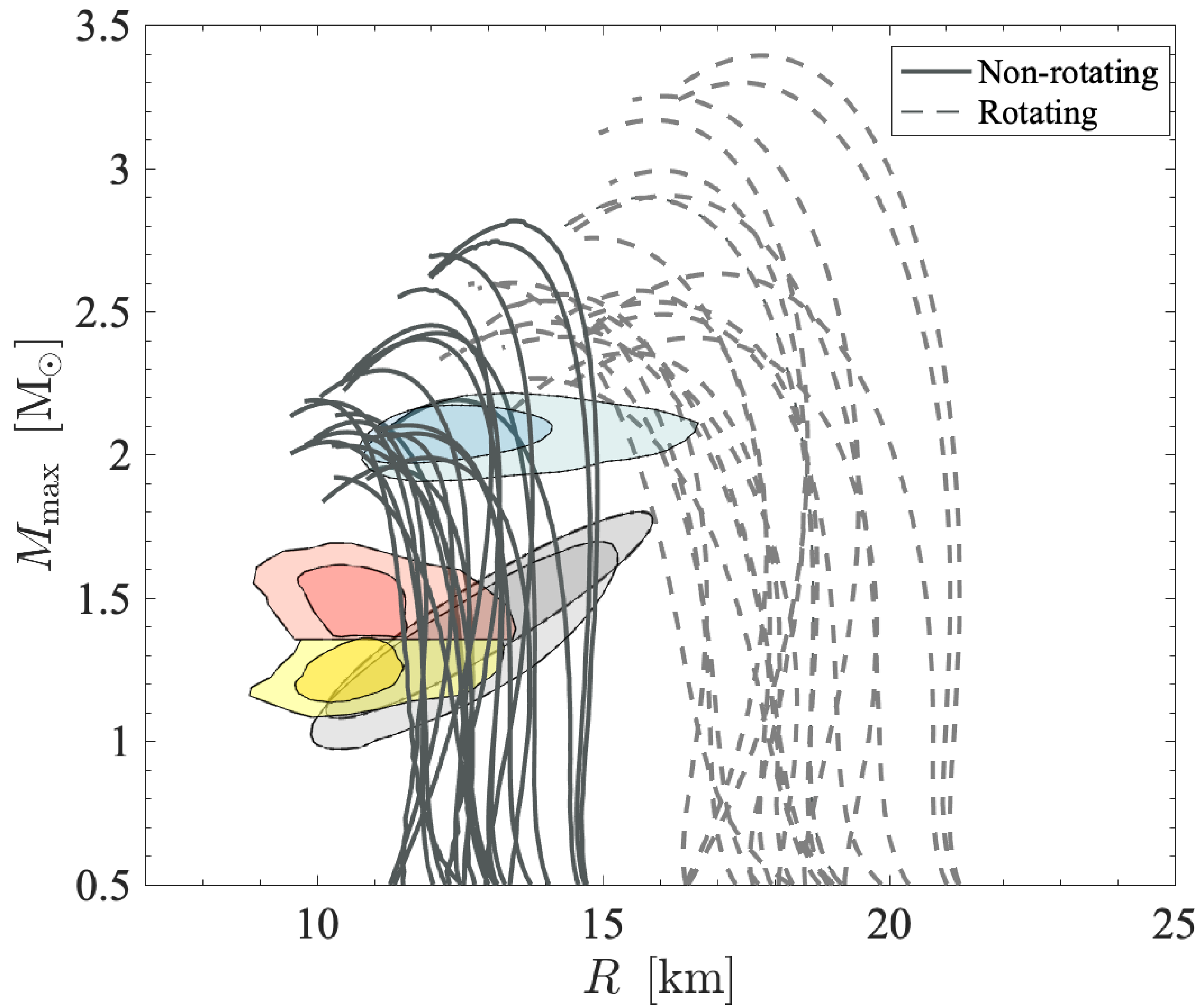
$$\Gamma = \sqrt{1 - v^2}$$

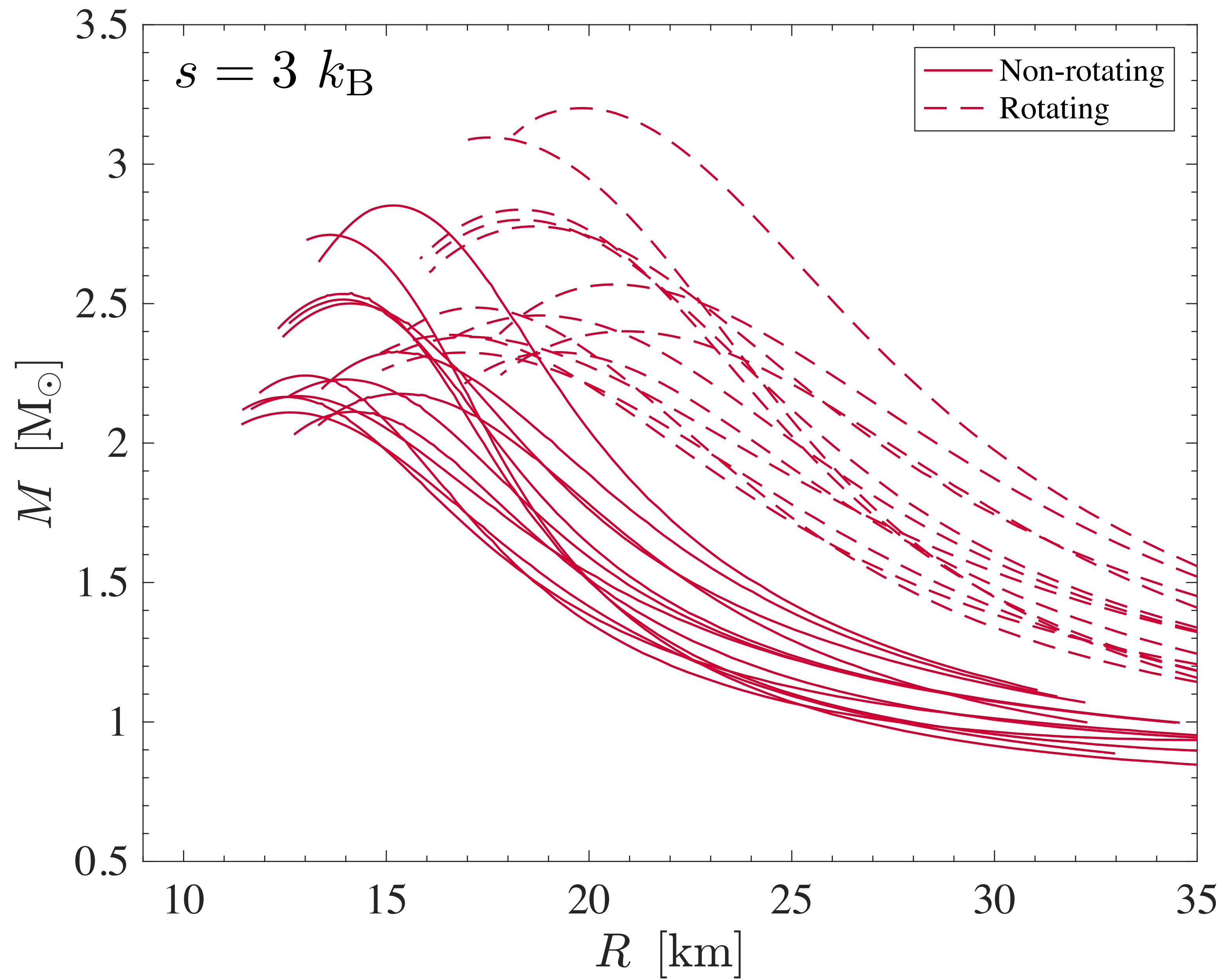
Cook G. B. Et al.(1994), Astrophys. J.

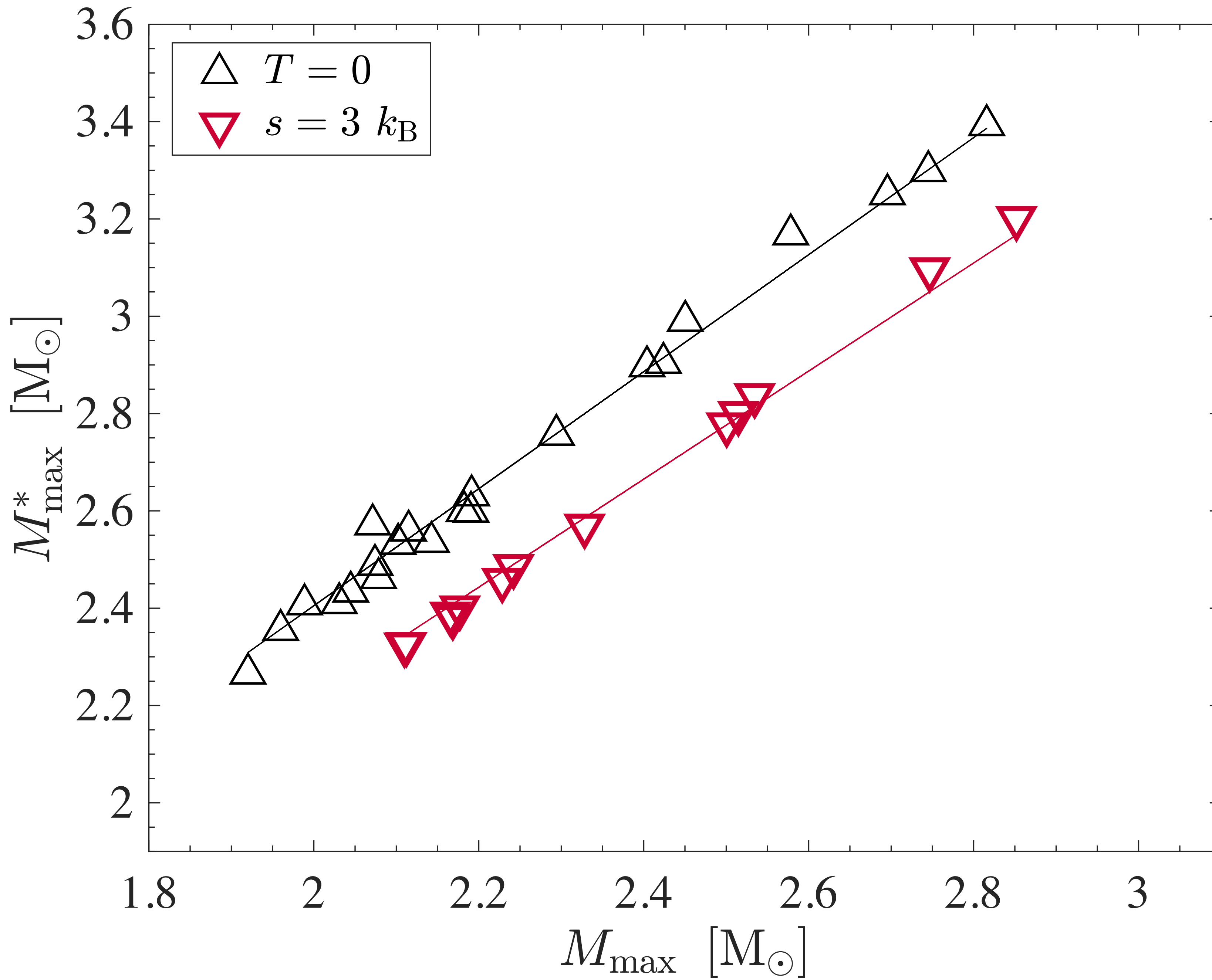
$$v = (\Omega - \omega) r \sin \theta e^{-\rho}$$

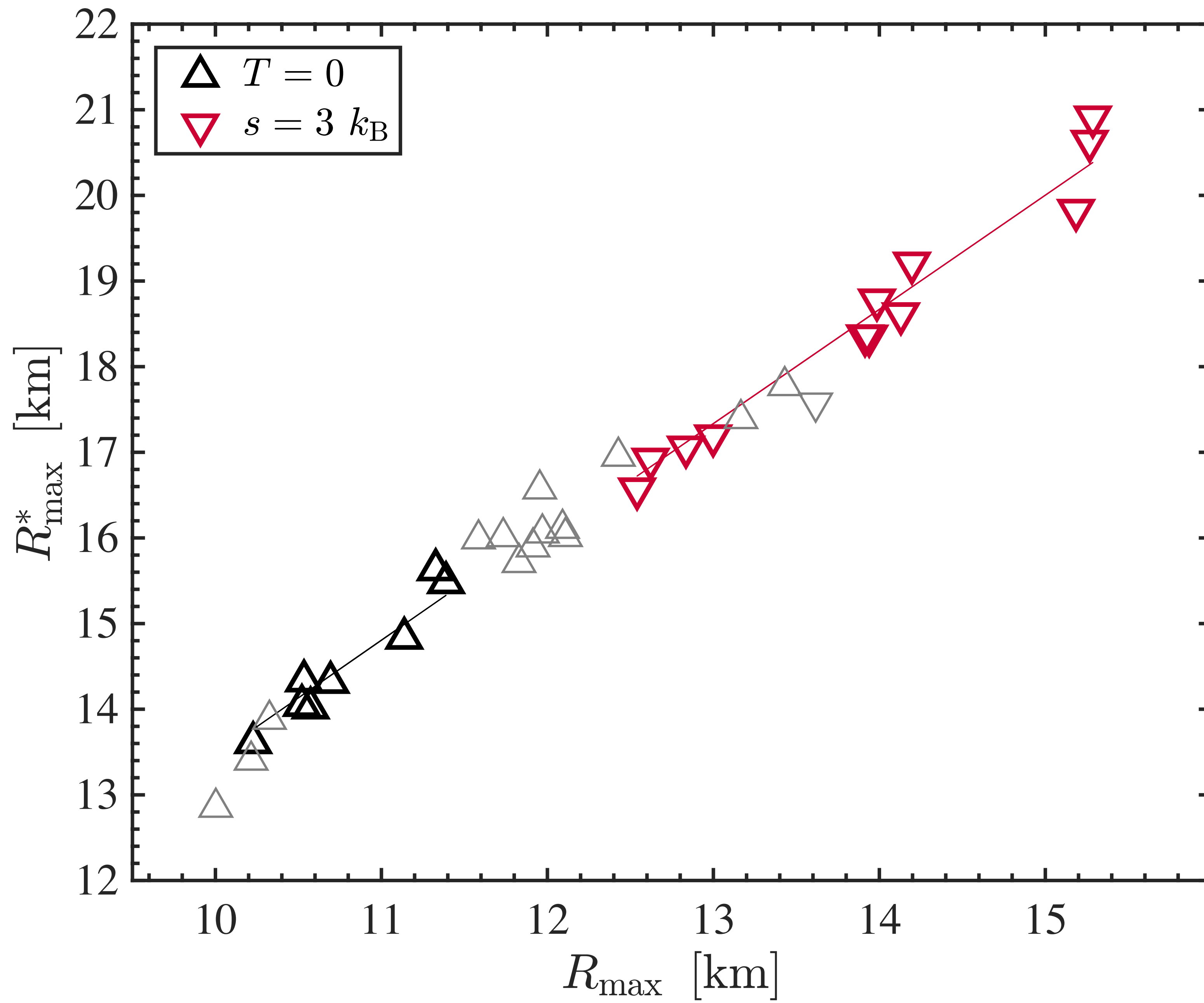
Numerical tool : RNS

Stergioulas et al. (1995), Astrophys. J.









$$T = 0$$

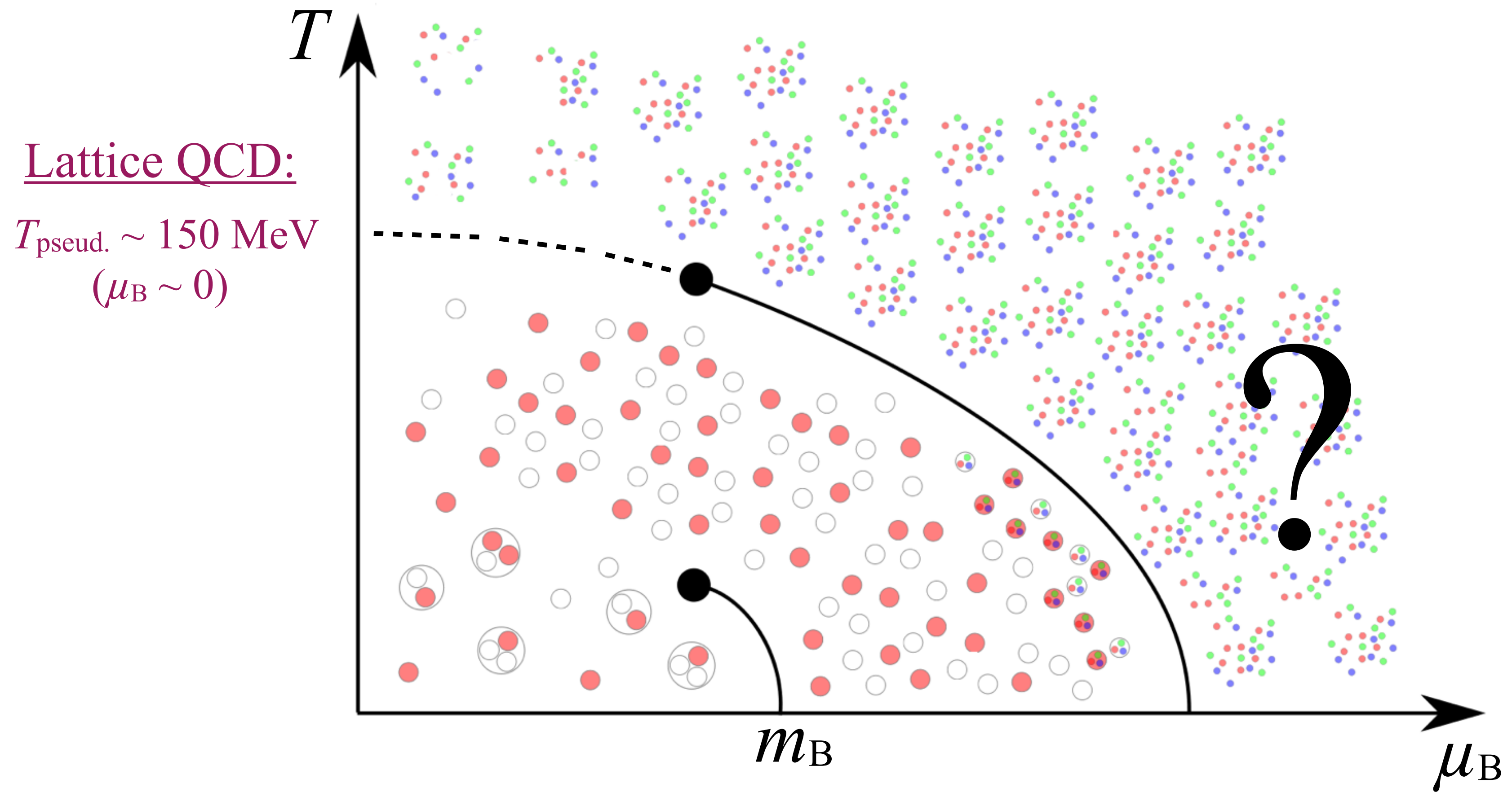
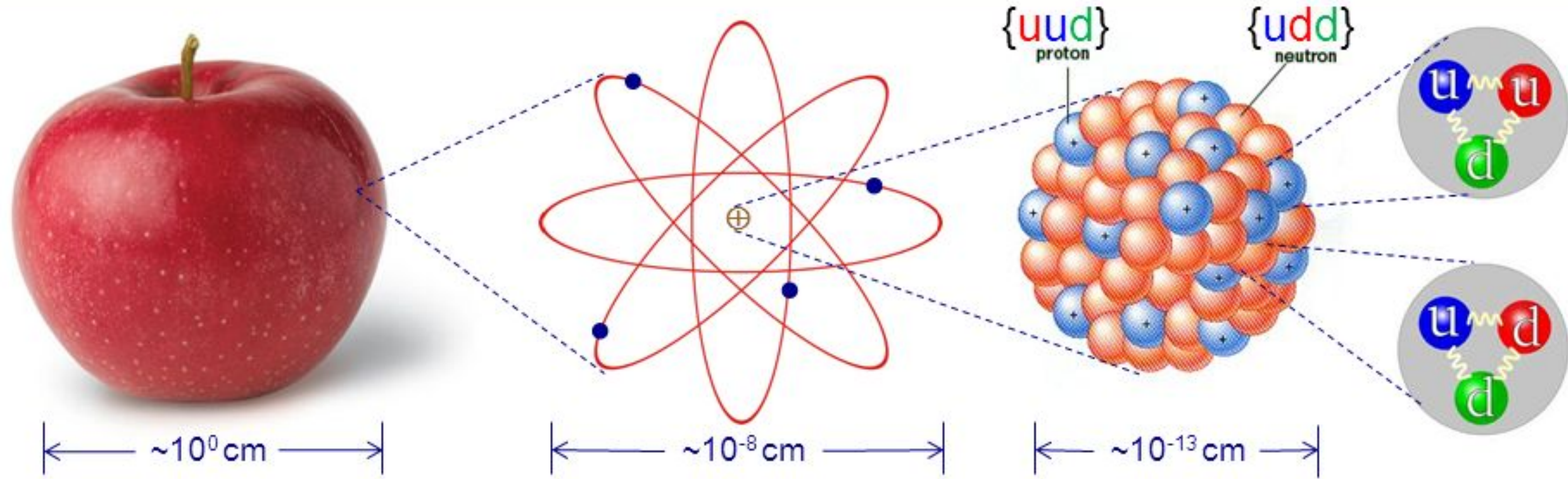
$$\frac{M_{\max}^*}{M_{\max}} = 1.200 \pm 0.0160$$

$$\frac{R_{\max}^*}{R_{\max}} = 1.346 \pm 0.0160 \quad (\text{Hadronic})$$

$$s = 3 k_B, Y_L = 0.3$$

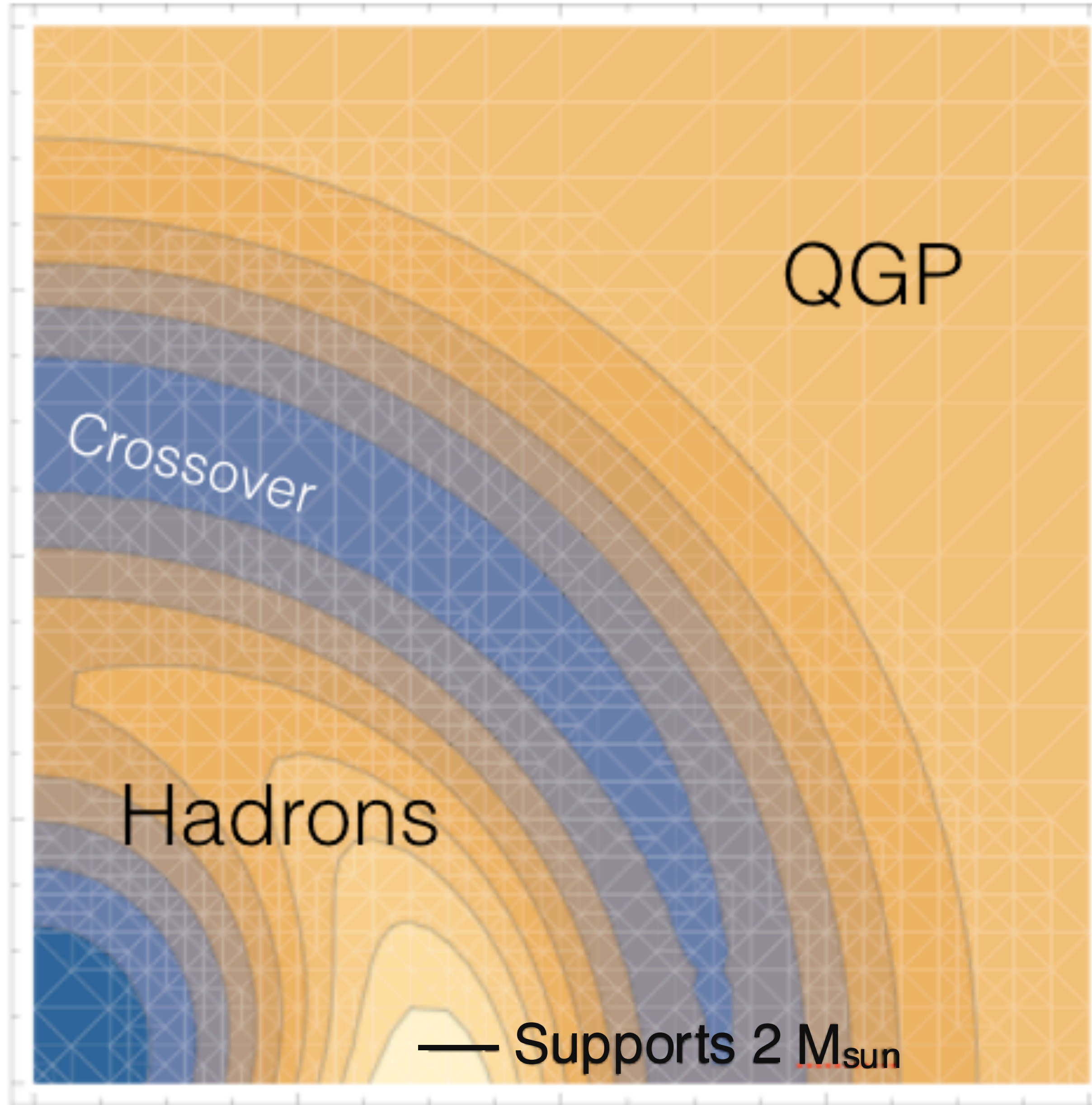
$$\frac{M_{\max}^*}{M_{\max}} = 1.109 \pm 0.0055$$

$$\frac{R_{\max}^*}{R_{\max}} = 1.334 \pm 0.0125 \quad (\text{Hadronic})$$



PSR J0348+0432: $2.01 \pm 0.04 M_{\odot}$

T



c_s^2

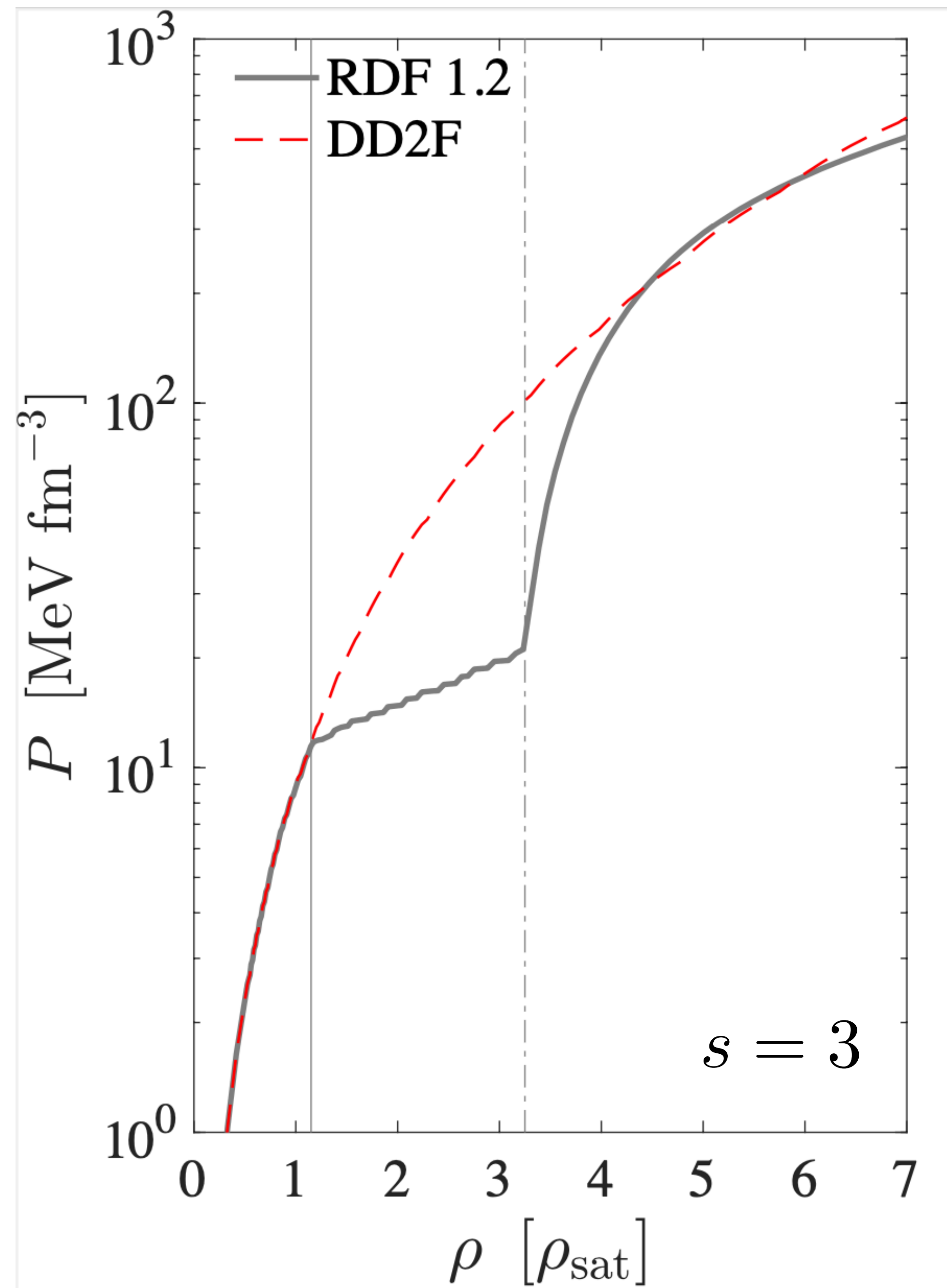
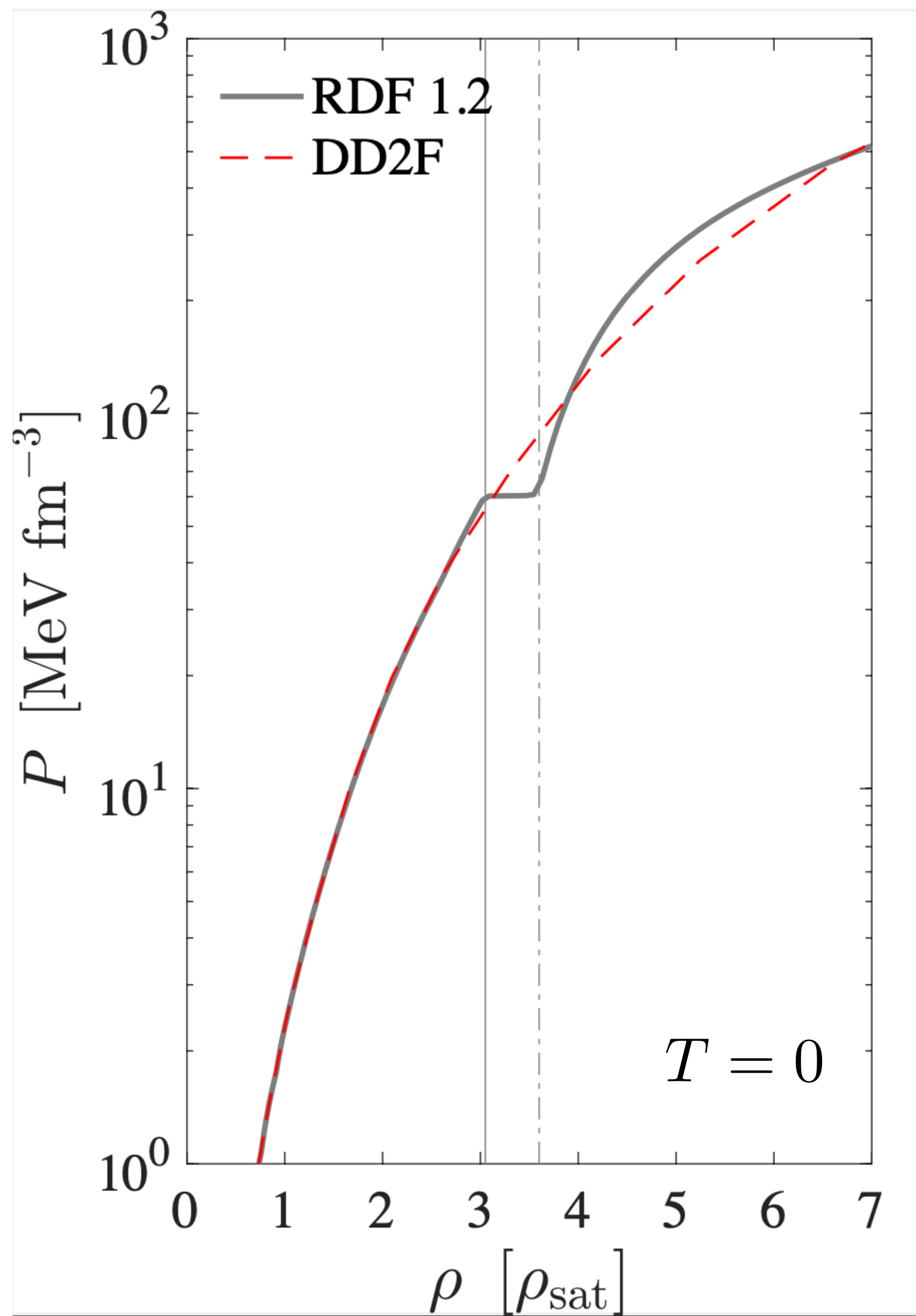
$> \frac{1}{3}$

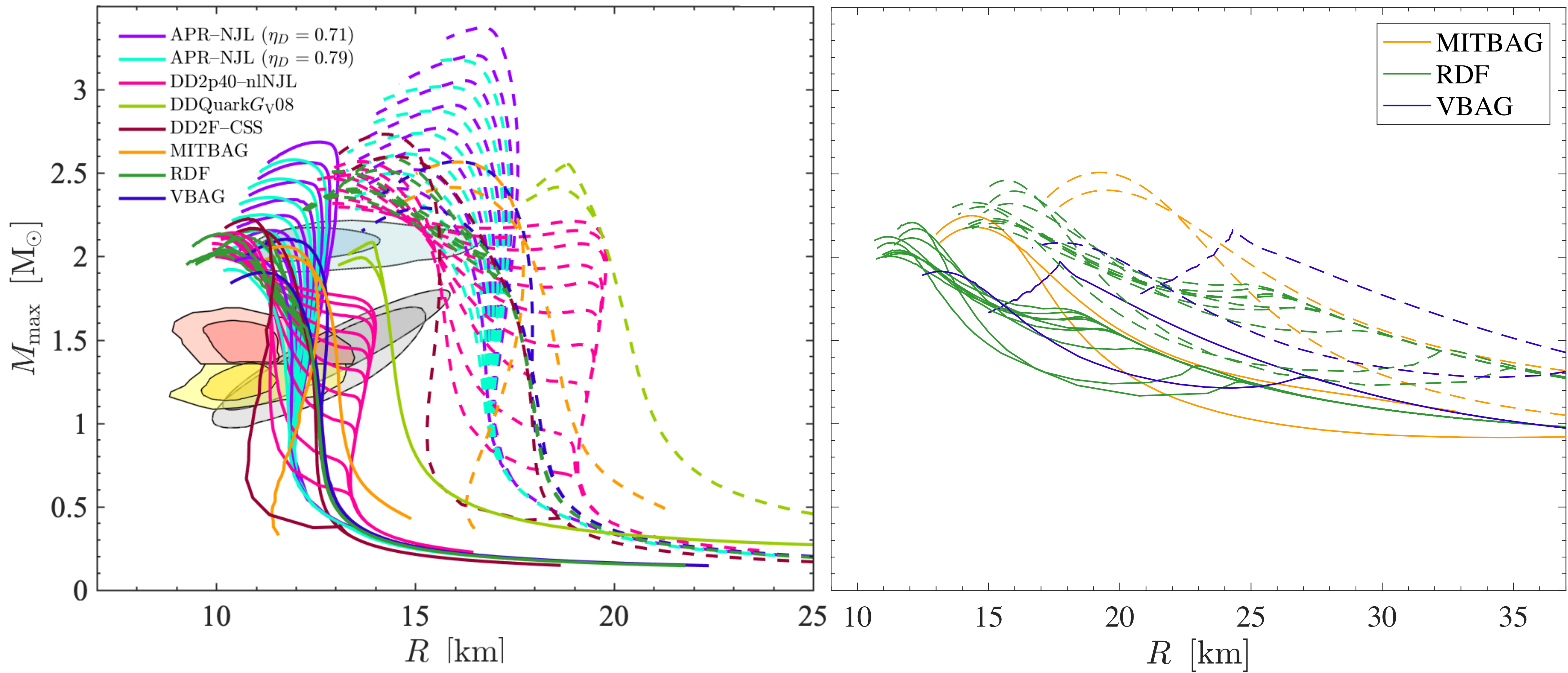
$\frac{1}{3}$

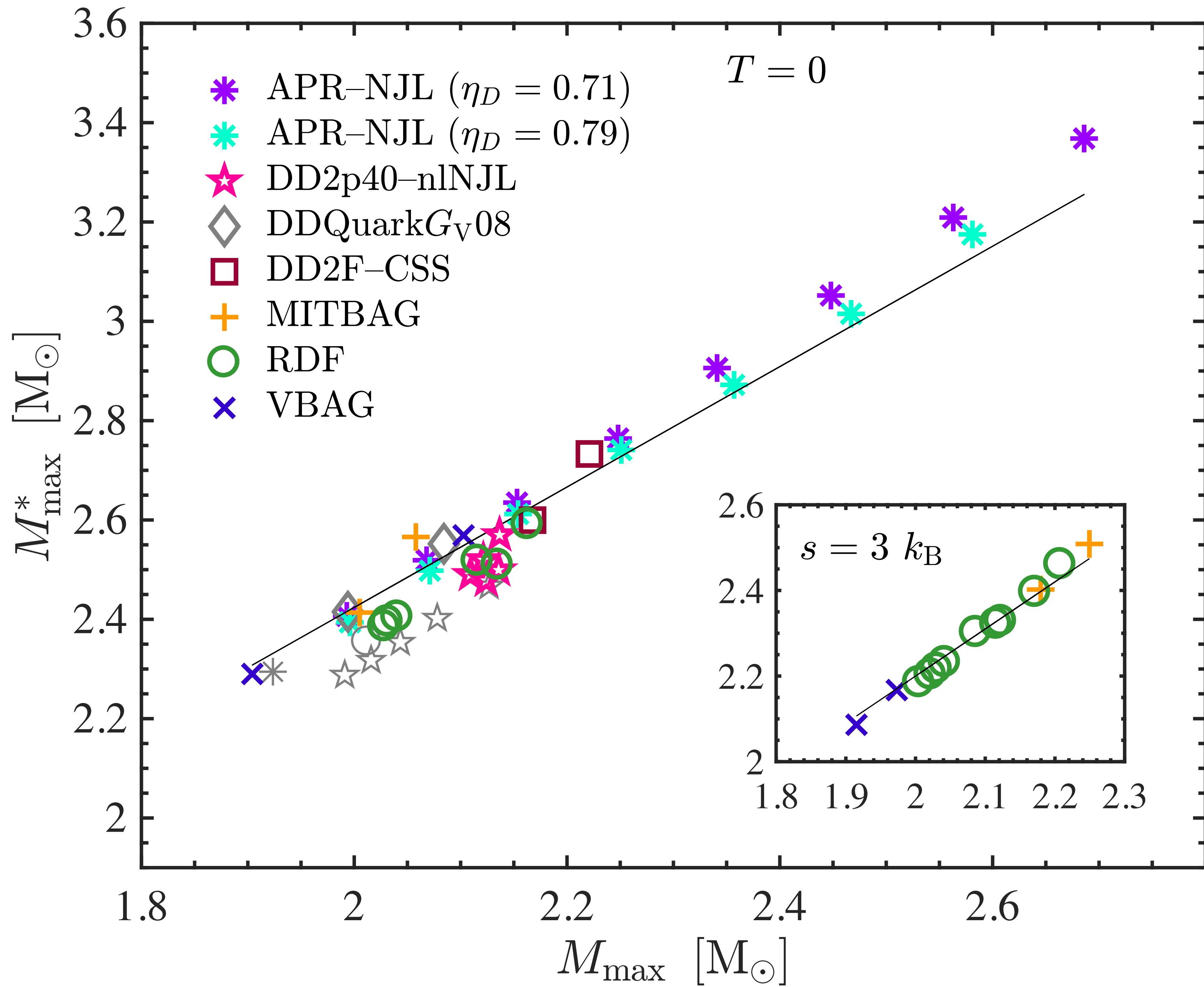
$< \frac{1}{3}$

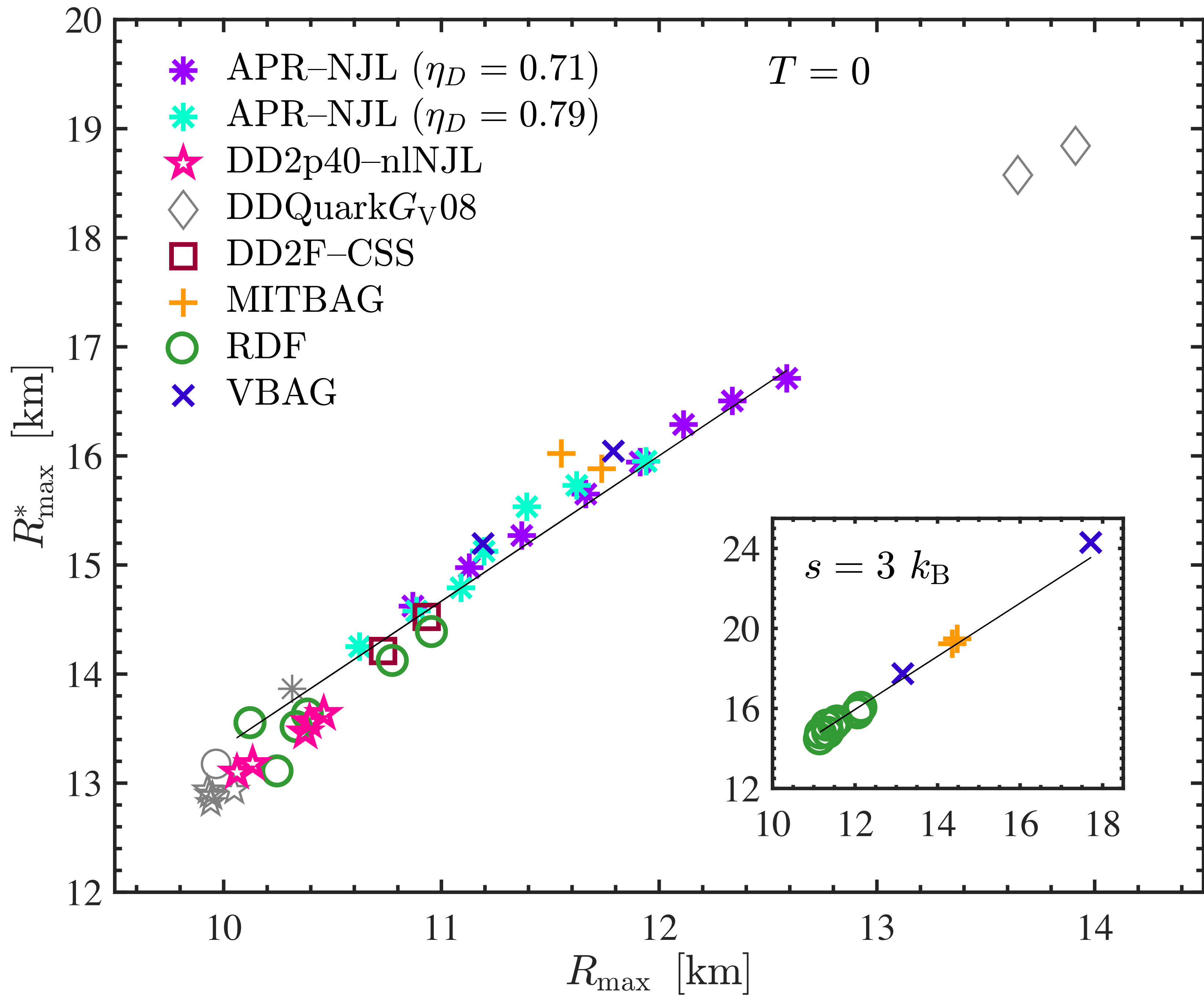
μ

— Supports $2 M_{\text{sun}}$









$$T = 0$$

$$\frac{M_{\max}^*}{M_{\max}} = 1.200 \pm 0.0160$$

$$\frac{R_{\max}^*}{R_{\max}} = 1.346 \pm 0.0160 \quad (\text{Hadronic})$$

$$\frac{M_{\max}^*}{M_{\max}} = 1.212 \pm 0.0090$$

$$\frac{R_{\max}^*}{R_{\max}} = 1.334 \pm 0.0085 \quad (\text{Hybrid})$$

$$s = 3 k_B, Y_L = 0.3$$

$$\frac{M_{\max}^*}{M_{\max}} = 1.109 \pm 0.0055$$

$$\frac{R_{\max}^*}{R_{\max}} = 1.334 \pm 0.0125 \quad (\text{Hadronic})$$

$$\frac{M_{\max}^*}{M_{\max}} = 1.100 \pm 0.0055$$

$$\frac{R_{\max}^*}{R_{\max}} = 1.329 \pm 0.0160 \quad (\text{Hybrid})$$

2. Core Collapse Supernovae

Supernova explosions of massive stars triggered by the QCD phase transition

PRL **102**, 081101 (2009)

PHYSICAL REVIEW LETTERS

week ending
27 FEBRUARY 2009

Signals of the QCD Phase Transition in Core-Collapse Supernovae

I. Sagert,¹ T. Fischer,³ M. Hempel,¹ G. Pagliara,² J. Schaffner-Bielich,² A. Mezzacappa,⁴
F.-K. Thielemann,³ and M. Liebendörfer³

¹*Institut für Theoretische Physik, Goethe-Universität, Max-von-Laue-Str. 1, 60438 Frankfurt am Main, Germany*

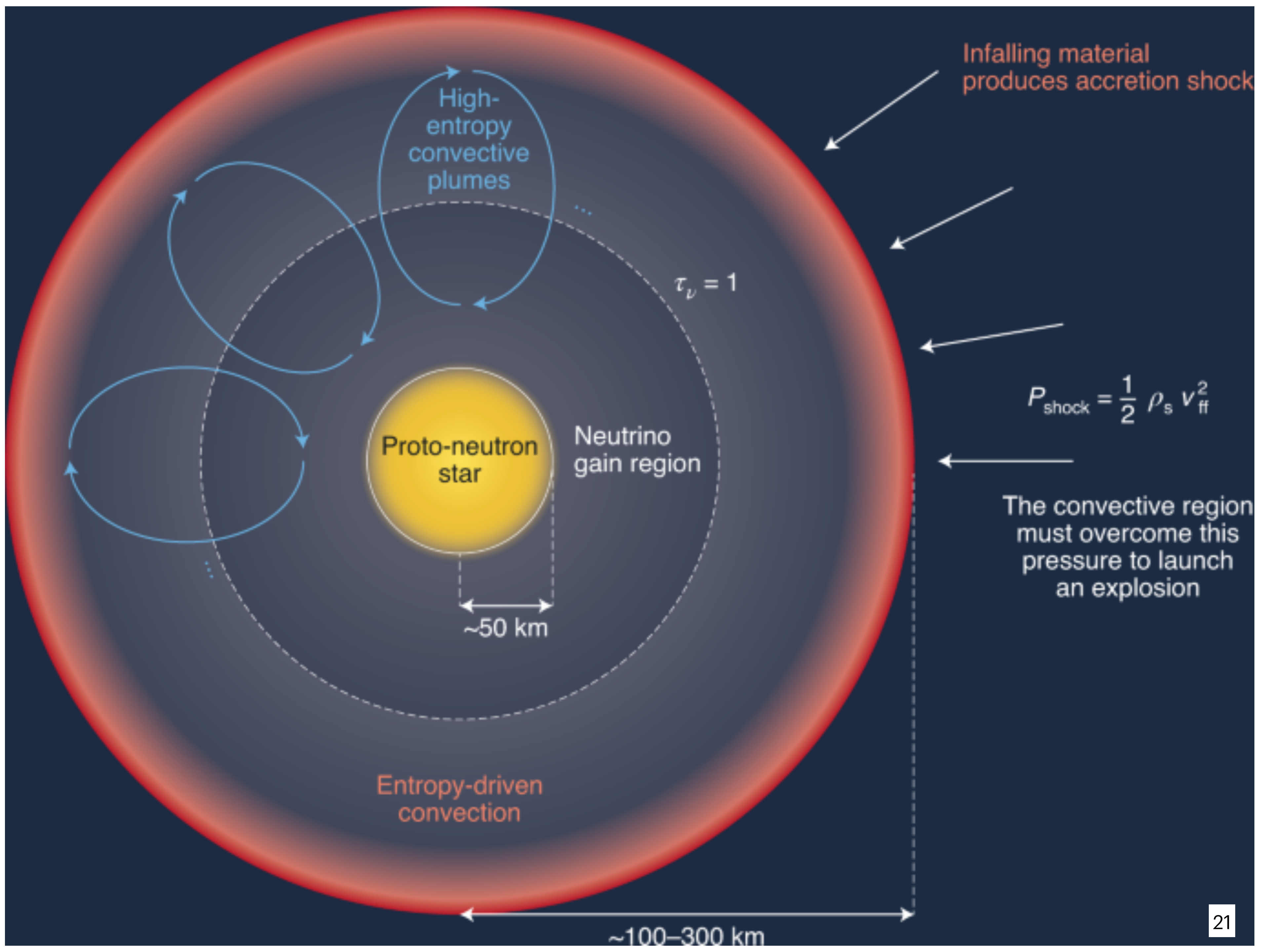
²*Institut für Theoretische Physik, Ruprecht-Karls-Universität, Philosophenweg 16, 69120 Heidelberg, Germany*

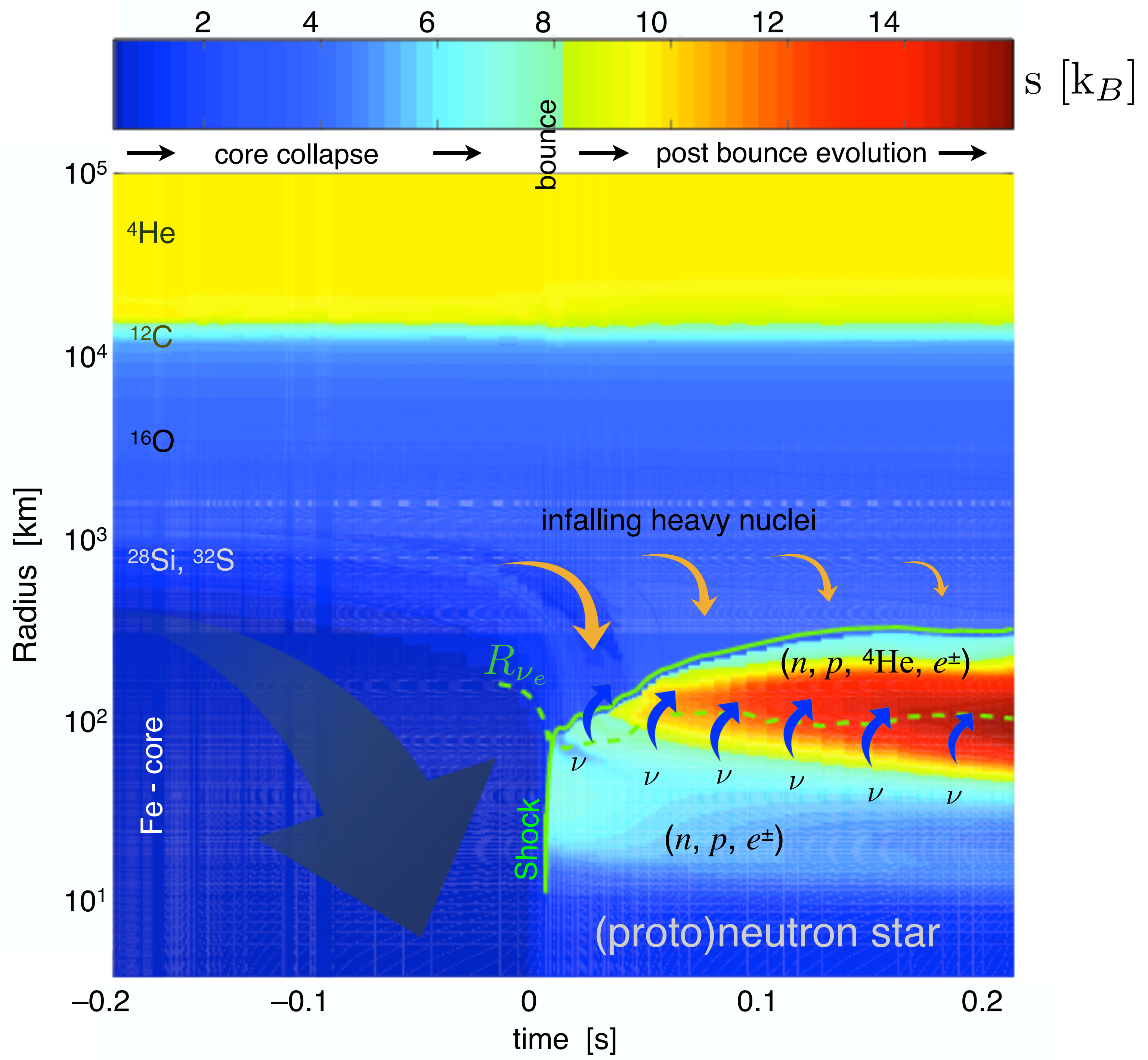
³*Department of Physics, University of Basel, Klingelbergstr. 82, 4056 Basel, Switzerland*

⁴*Physics Division, Oak Ridge National Laboratory, Oak Ridge, Tennessee 37831, USA*

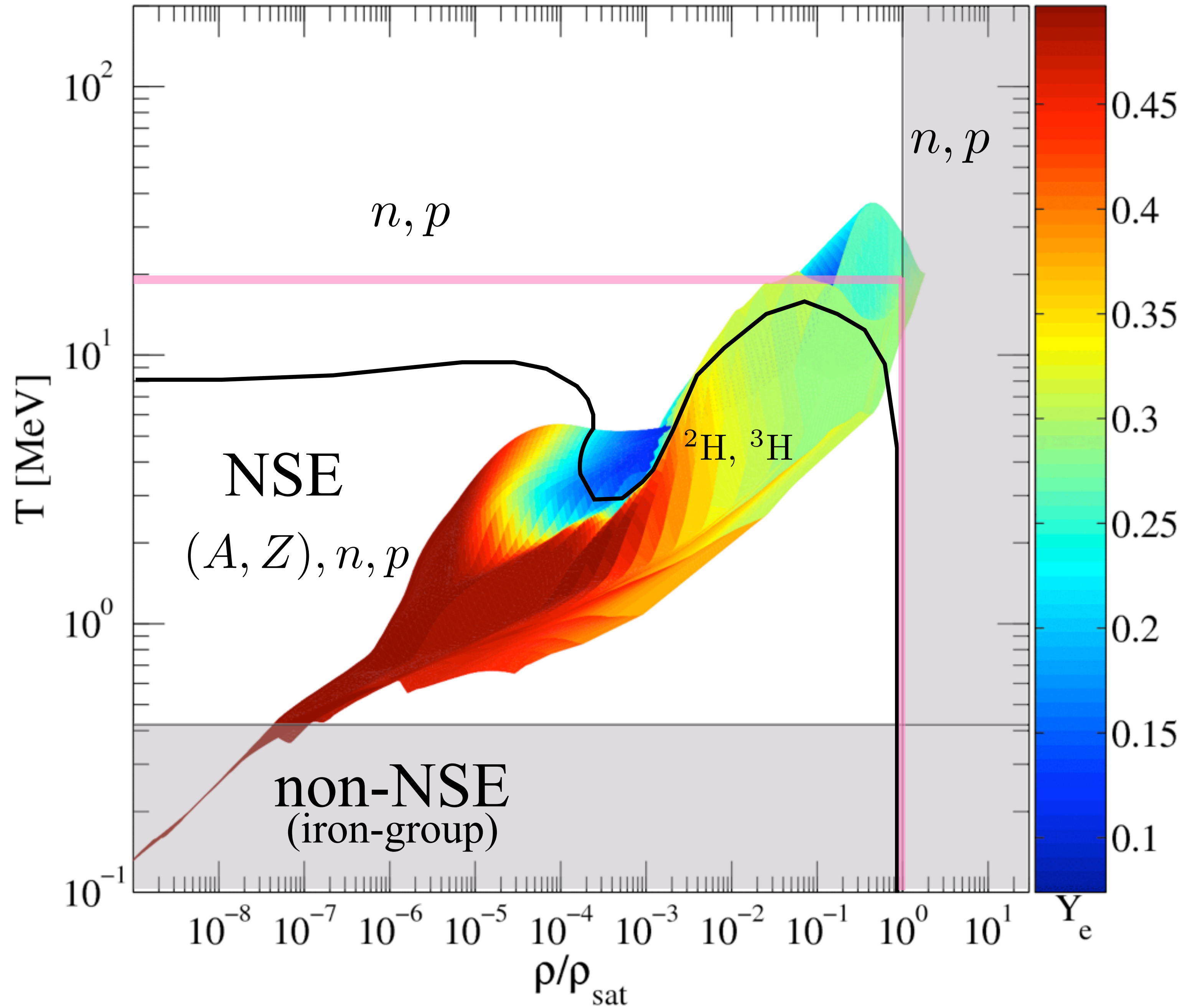
(Received 12 August 2008; published 26 February 2009)

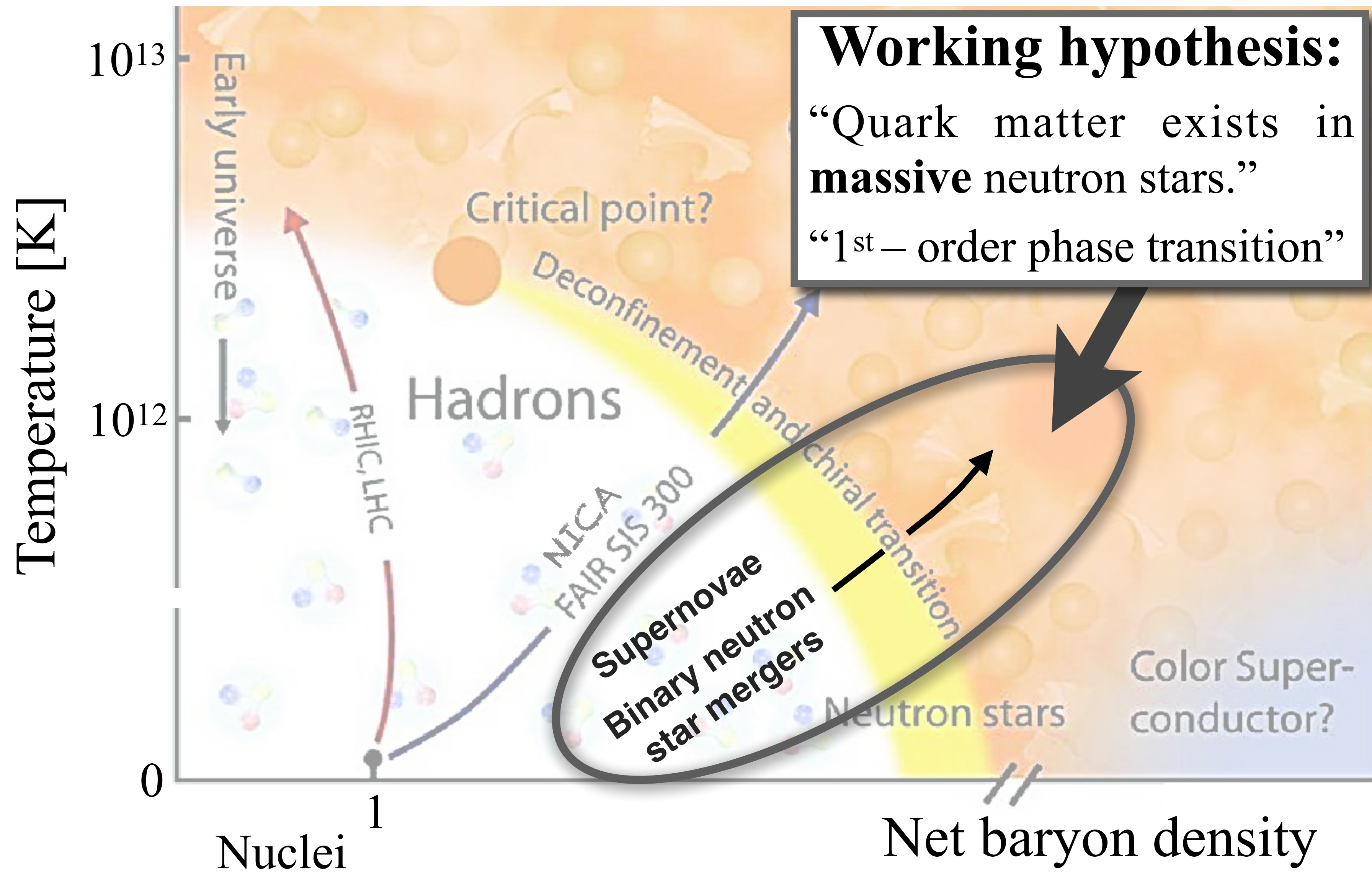
We explore the implications of the QCD phase transition during the postbounce evolution of core-collapse supernovae. Using the MIT bag model for the description of quark matter, we model phase transitions that occur during the early postbounce evolution. This stage of the evolution can be simulated with general relativistic three-flavor Boltzmann neutrino transport. The phase transition produces a second shock wave that triggers a delayed supernova explosion. If such a phase transition happens in a future galactic supernova, its existence and properties should become observable as a second peak in the neutrino signal that is accompanied by significant changes in the energy of the emitted neutrinos. This second neutrino burst is dominated by the emission of antineutrinos because the electron degeneracy is reduced when the second shock passes through the previously neutronized matter.



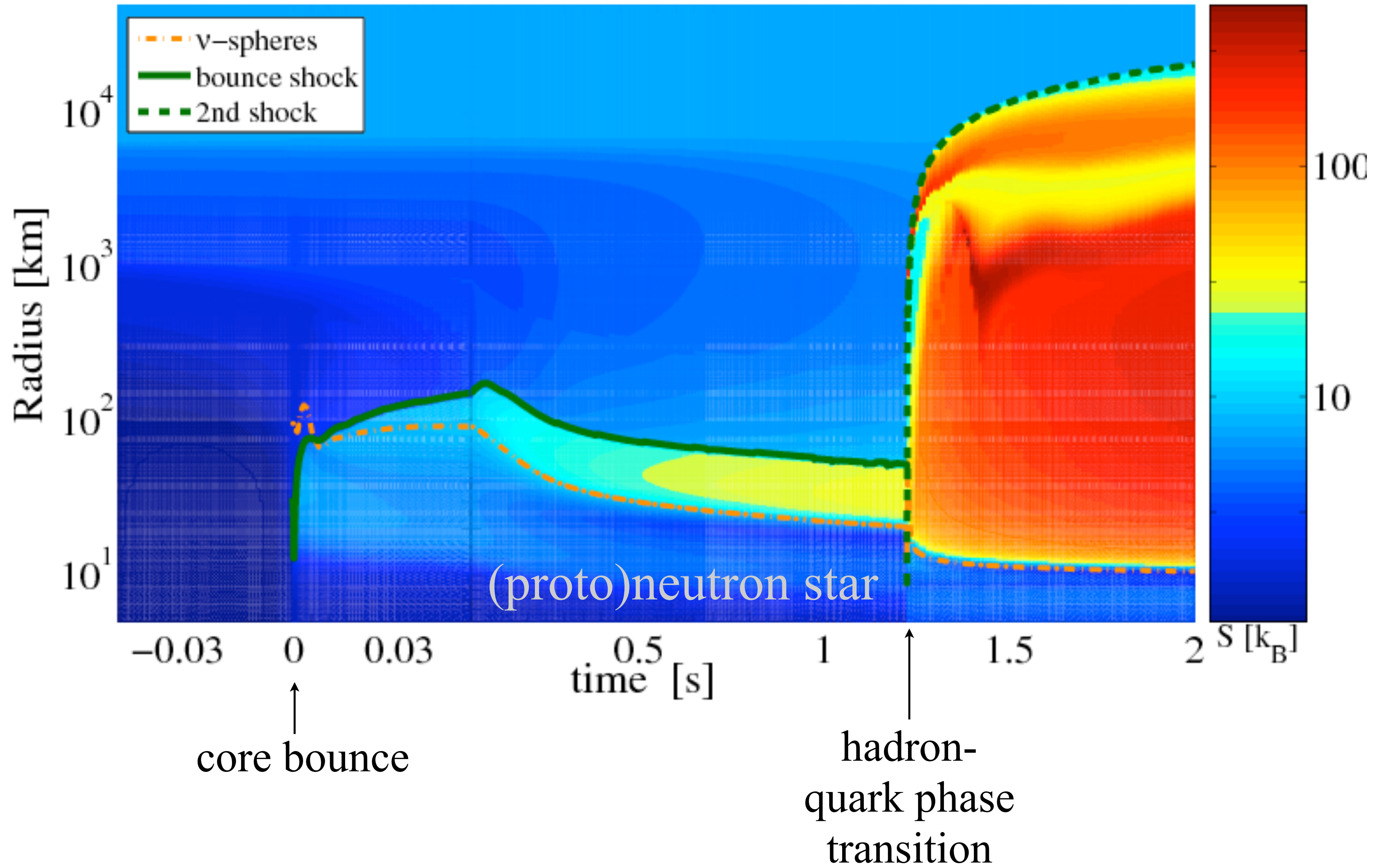


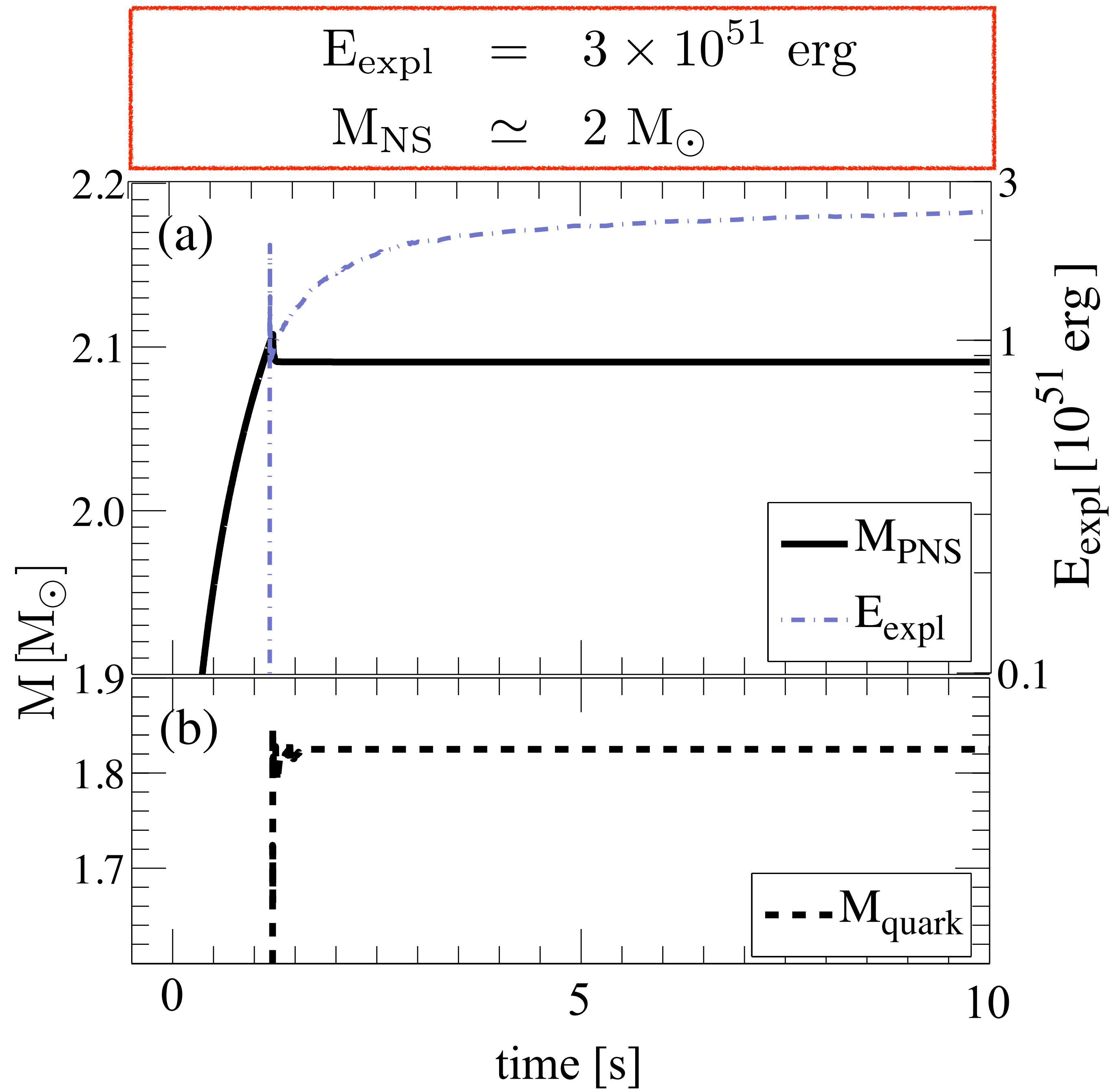
0.50165 s after bounce

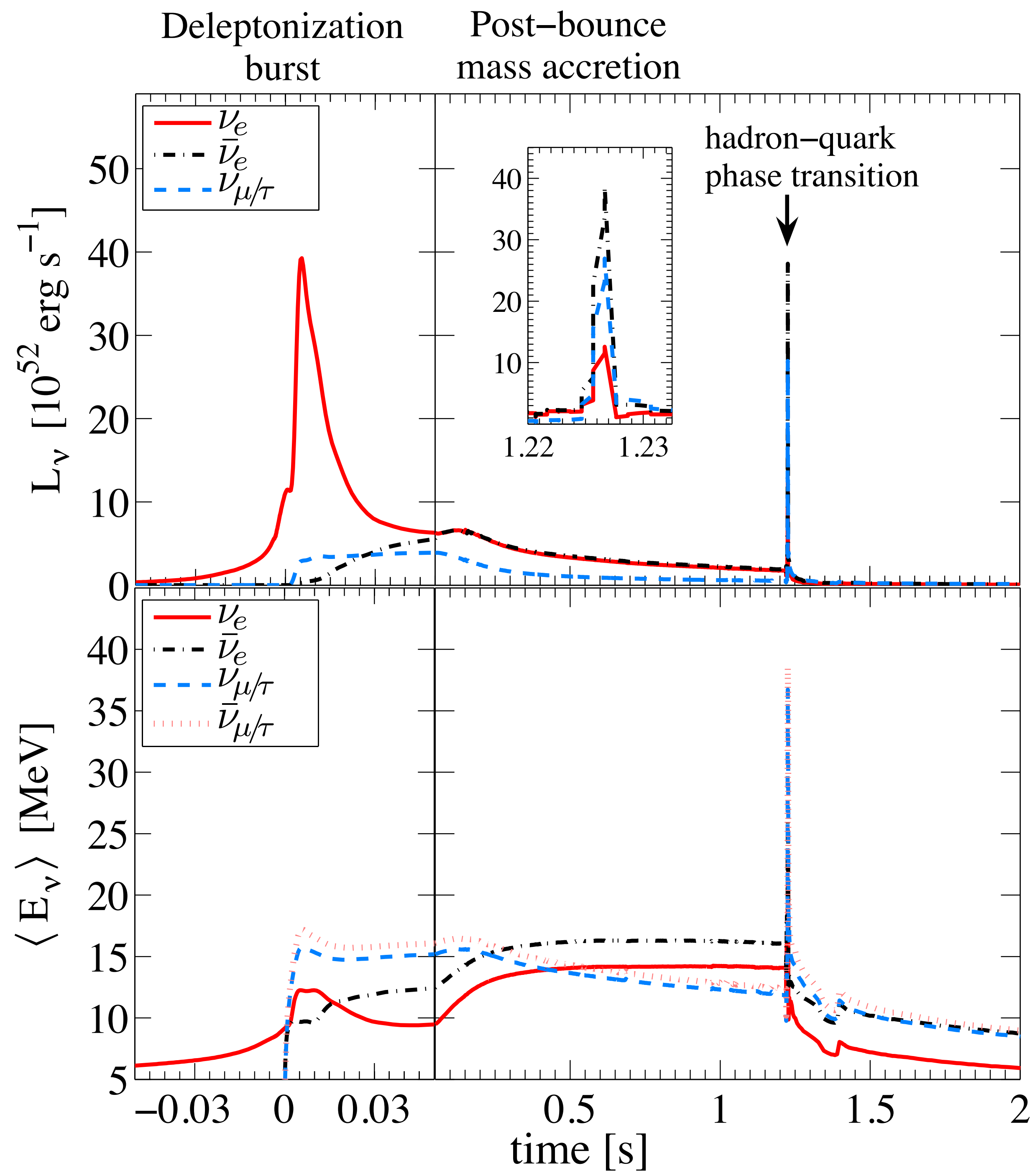




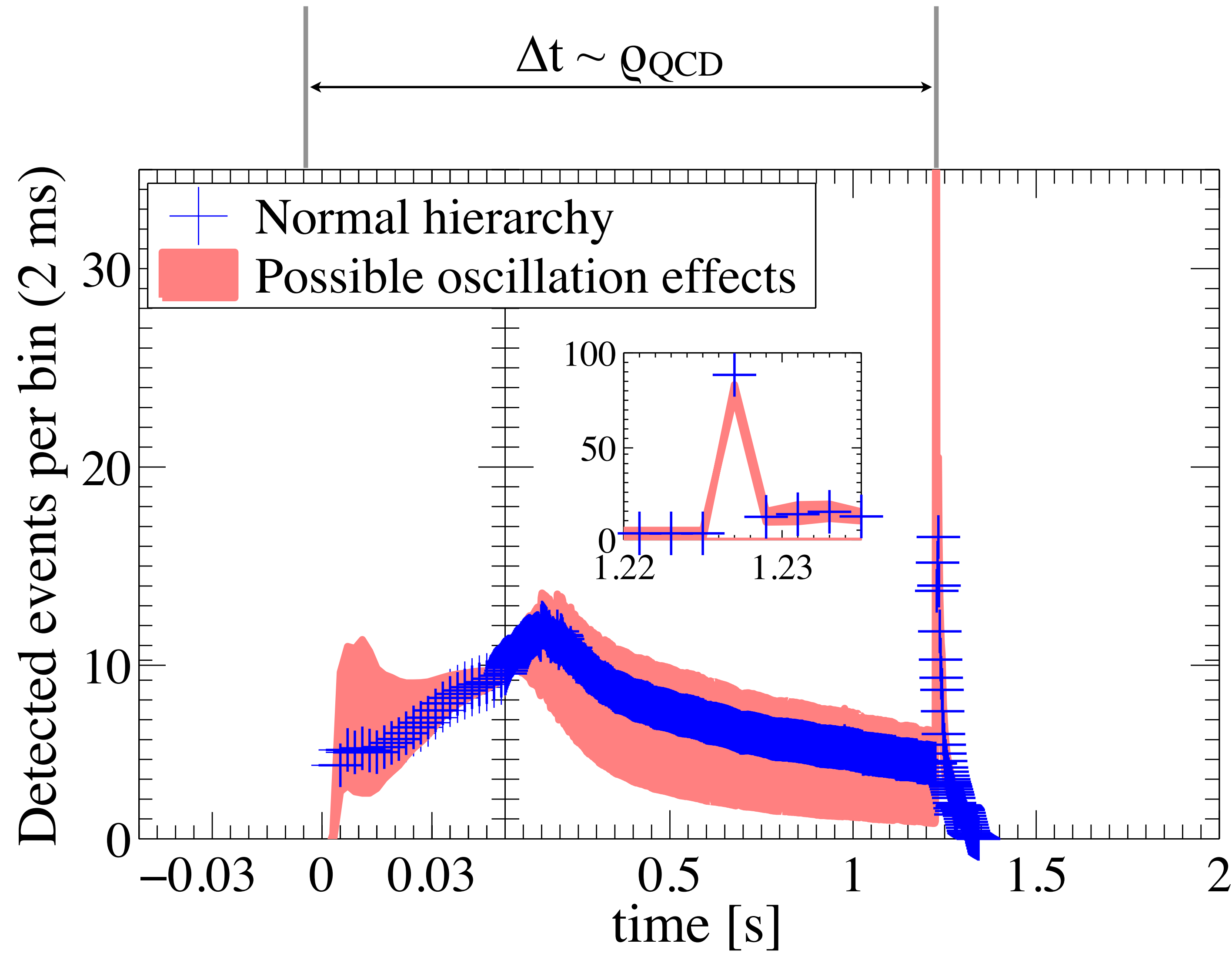
$$M_{\text{ZAMS}} = 50 M_{\odot}$$

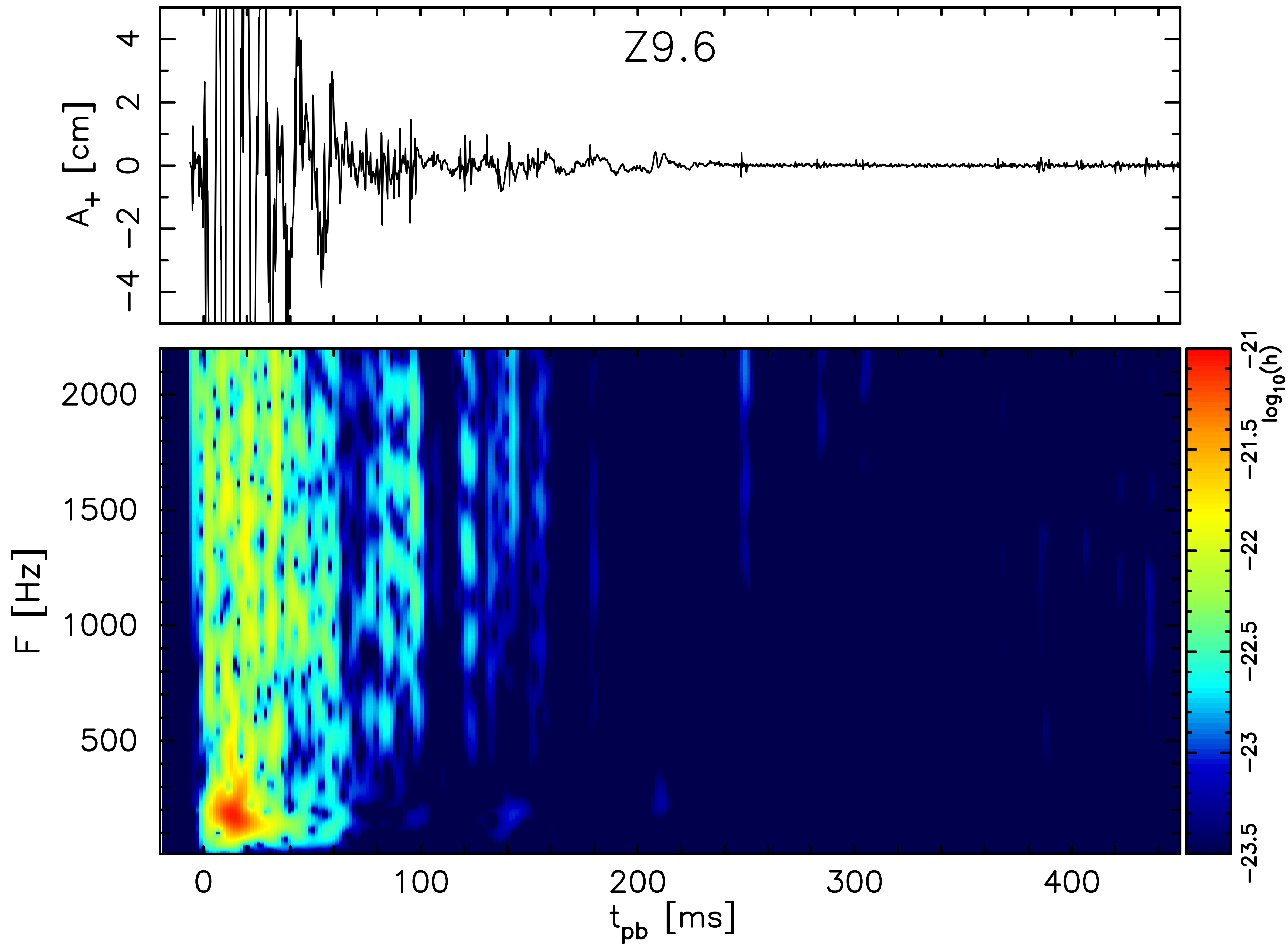


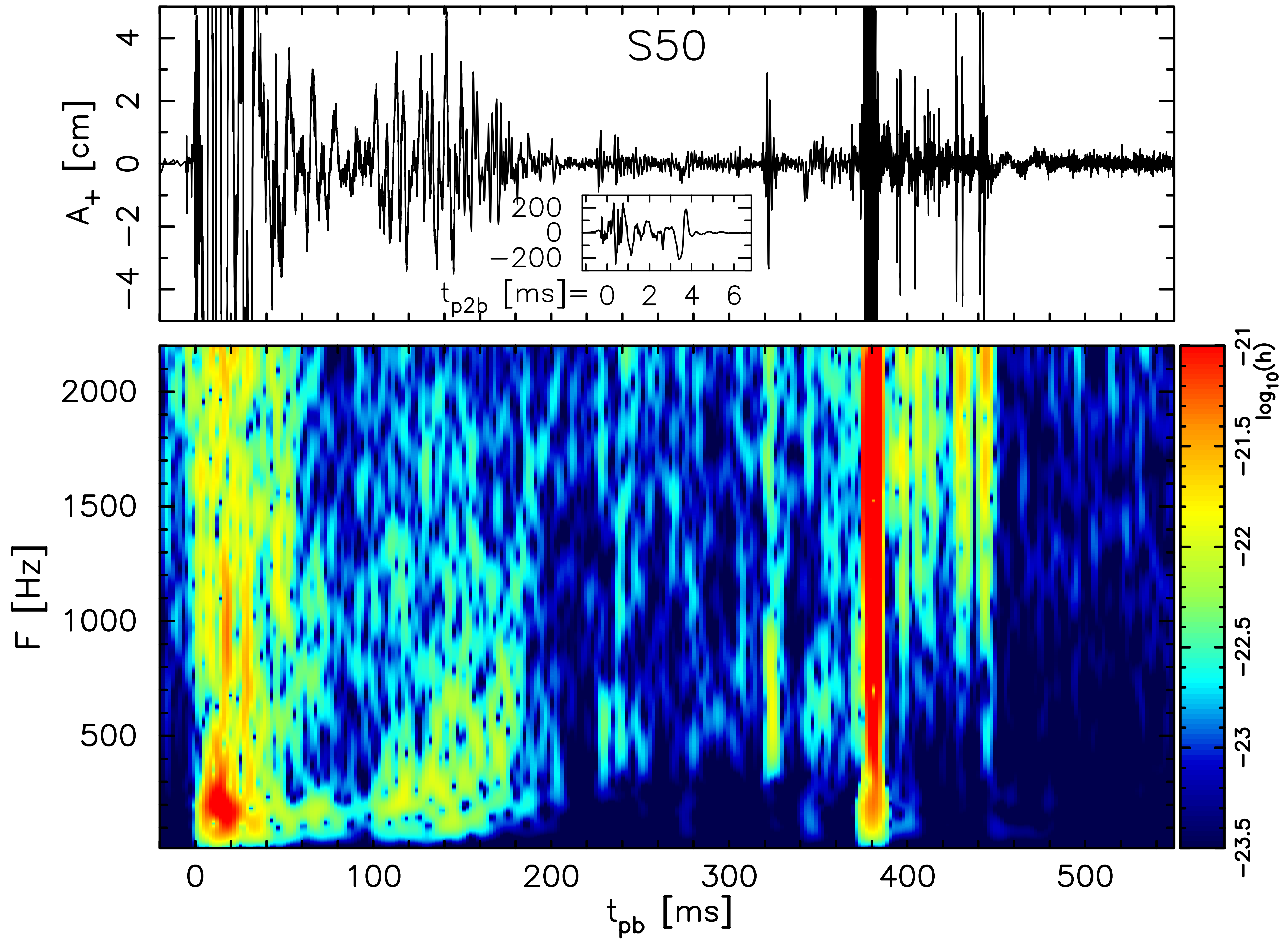


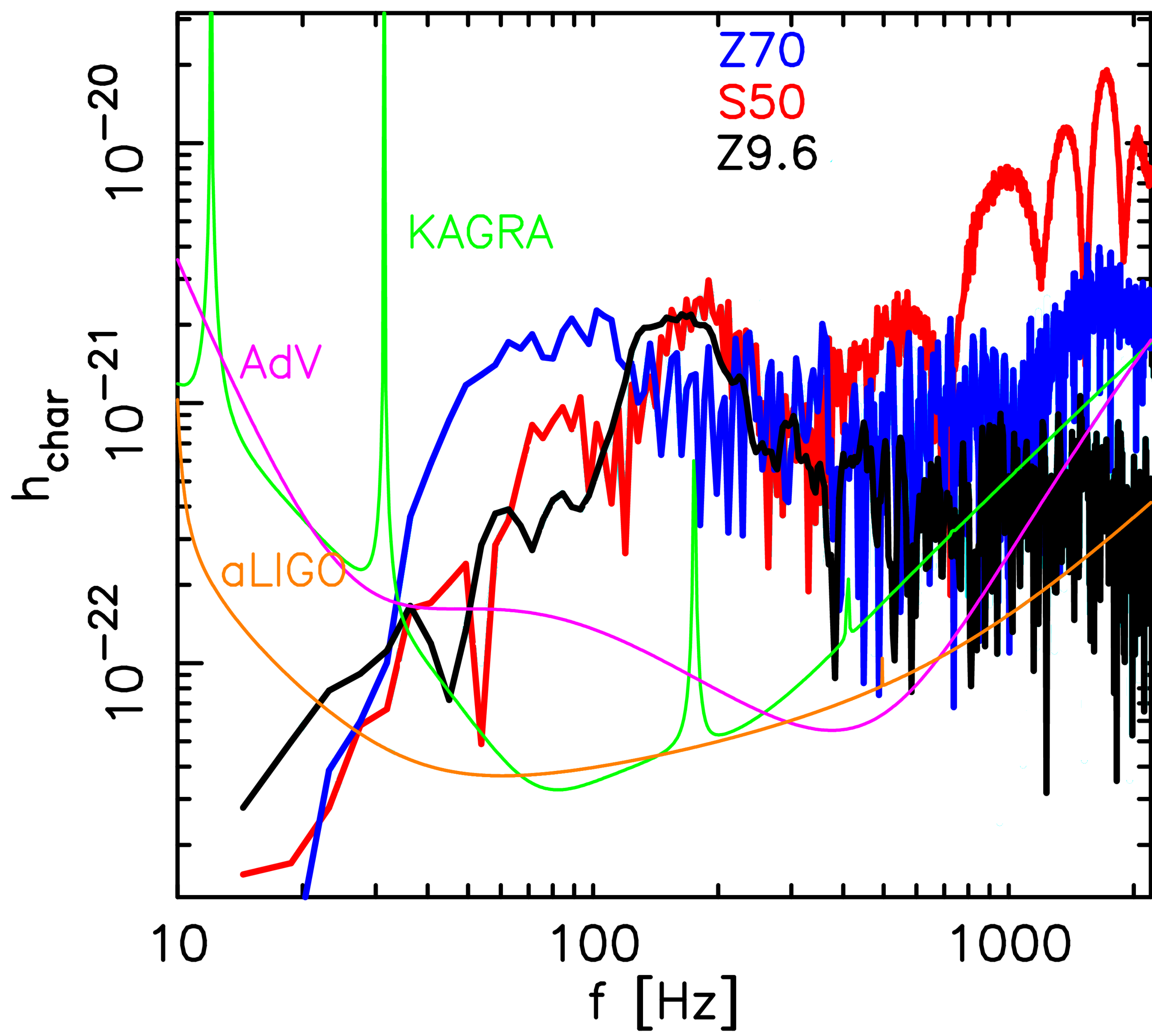


ν – signal @ Super-Kamiokande ($d \sim 10$ kpc)









Numerical tool for mode analysis : GREAT

$$Q_{ij} = \int \rho(r) (3r_i r_j - |\vec{r}|^2 \delta_{ij}) d^3 r \quad \rho \rightarrow \rho + \delta\rho, \quad \Delta\rho = \delta\rho + \xi^i \partial_i \rho$$

$$h_{ij} = \frac{2G}{c^4 r} \frac{d^2}{dt^2} Q_{ij} \quad \partial_t \xi^i = \delta v^{*i}$$

g-modes

$$\partial_r \eta_r = A \eta_r + B \eta_\perp,$$

$$\partial_r \eta_r + \left[\frac{2}{r} + 6 \frac{\partial_r \psi}{\psi} \right] \eta_r - \frac{l(l+1)}{r^2} \eta_\perp = 0,$$

$$\partial_r \eta_\perp = C \eta_r + D \eta_\perp.$$

$$\partial_r \eta_\perp - \left(1 - \frac{\mathcal{N}^2}{\sigma^2} \right) \eta_r + [\partial_r \ln q - G] \eta_\perp = 0.$$

$$\delta P = \delta \hat{P} Y_{lm} e^{-i\sigma t},$$

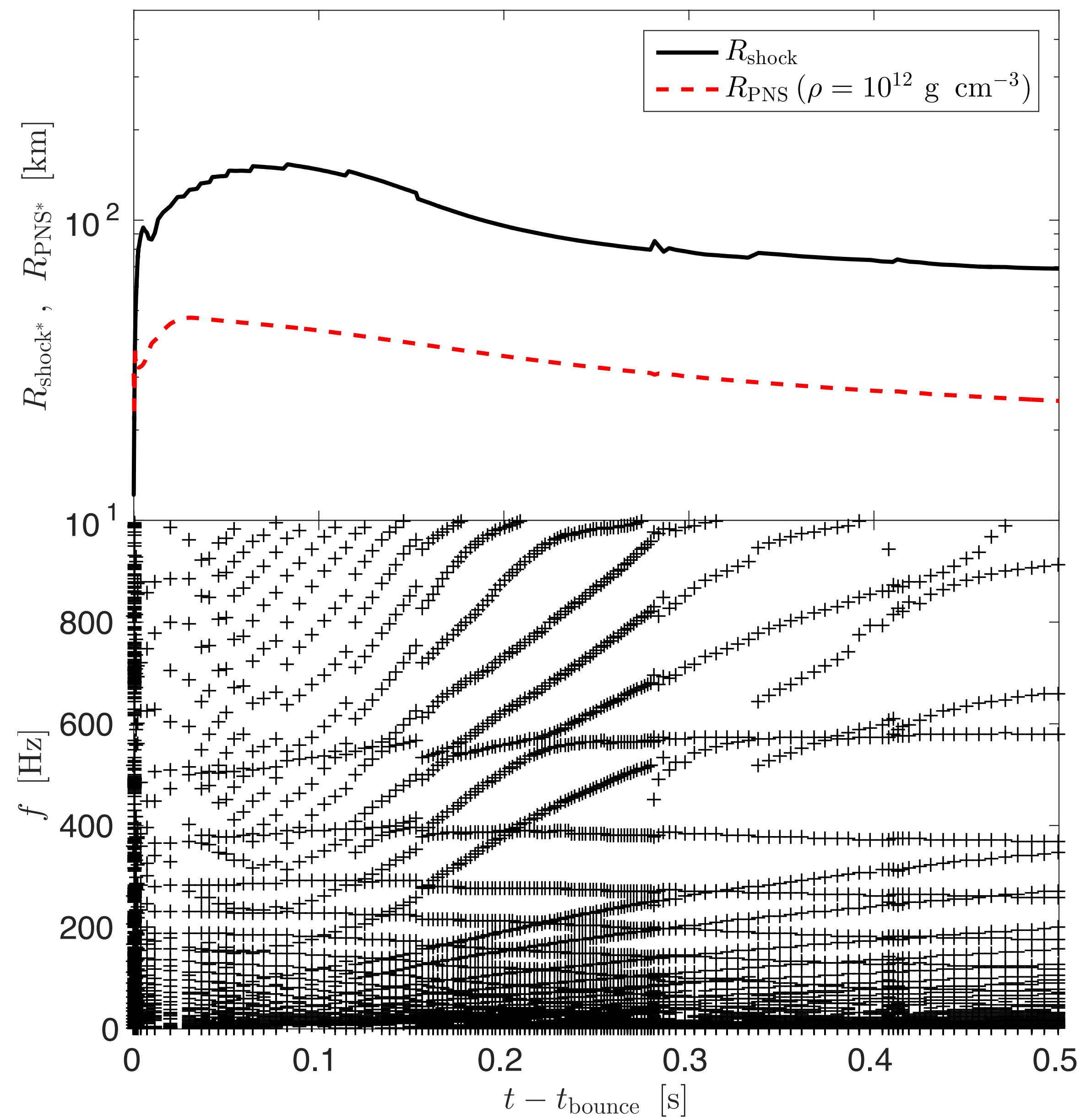
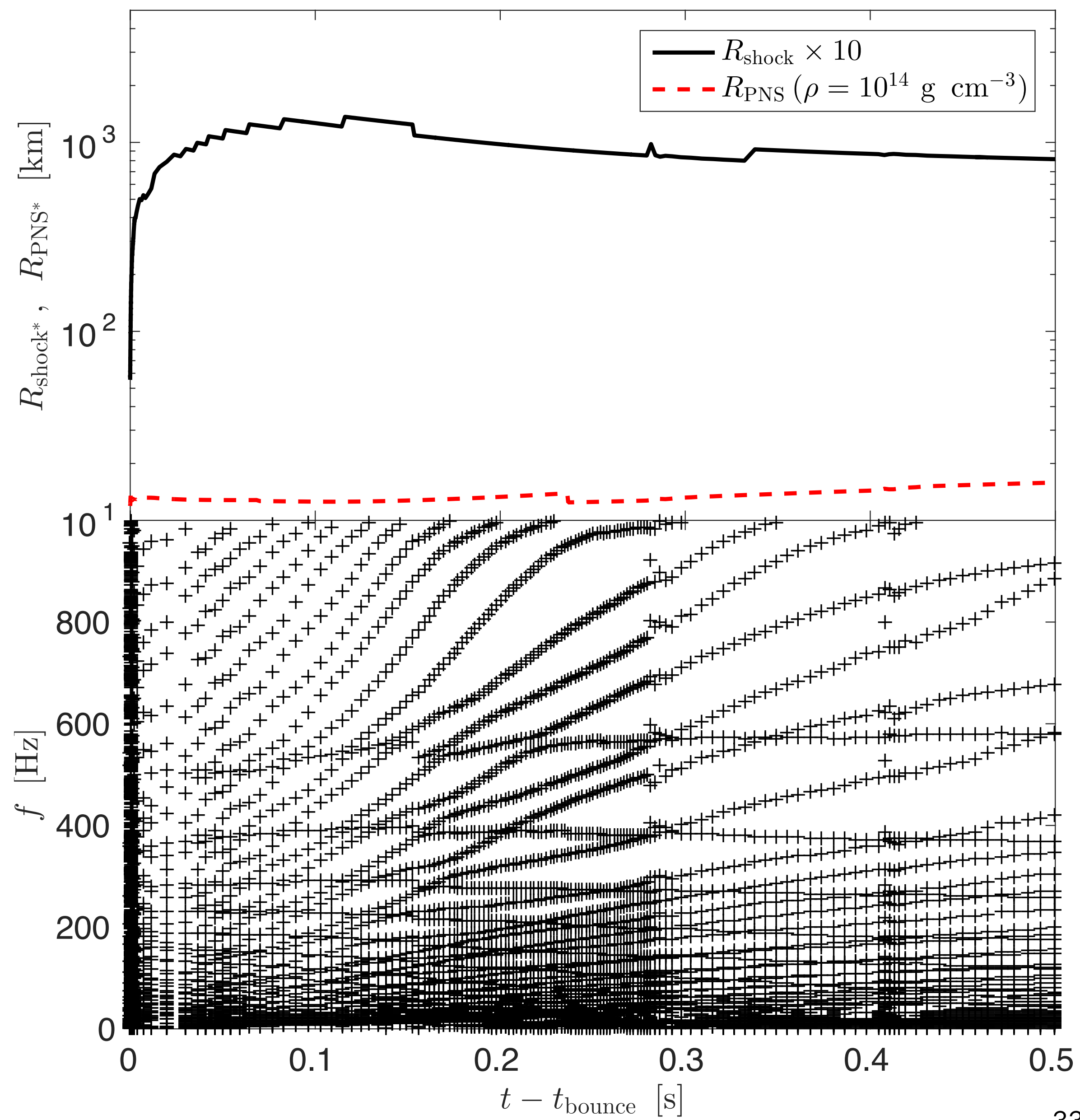
$$\xi^r = \eta_r Y_{lm} e^{-i\sigma t},$$

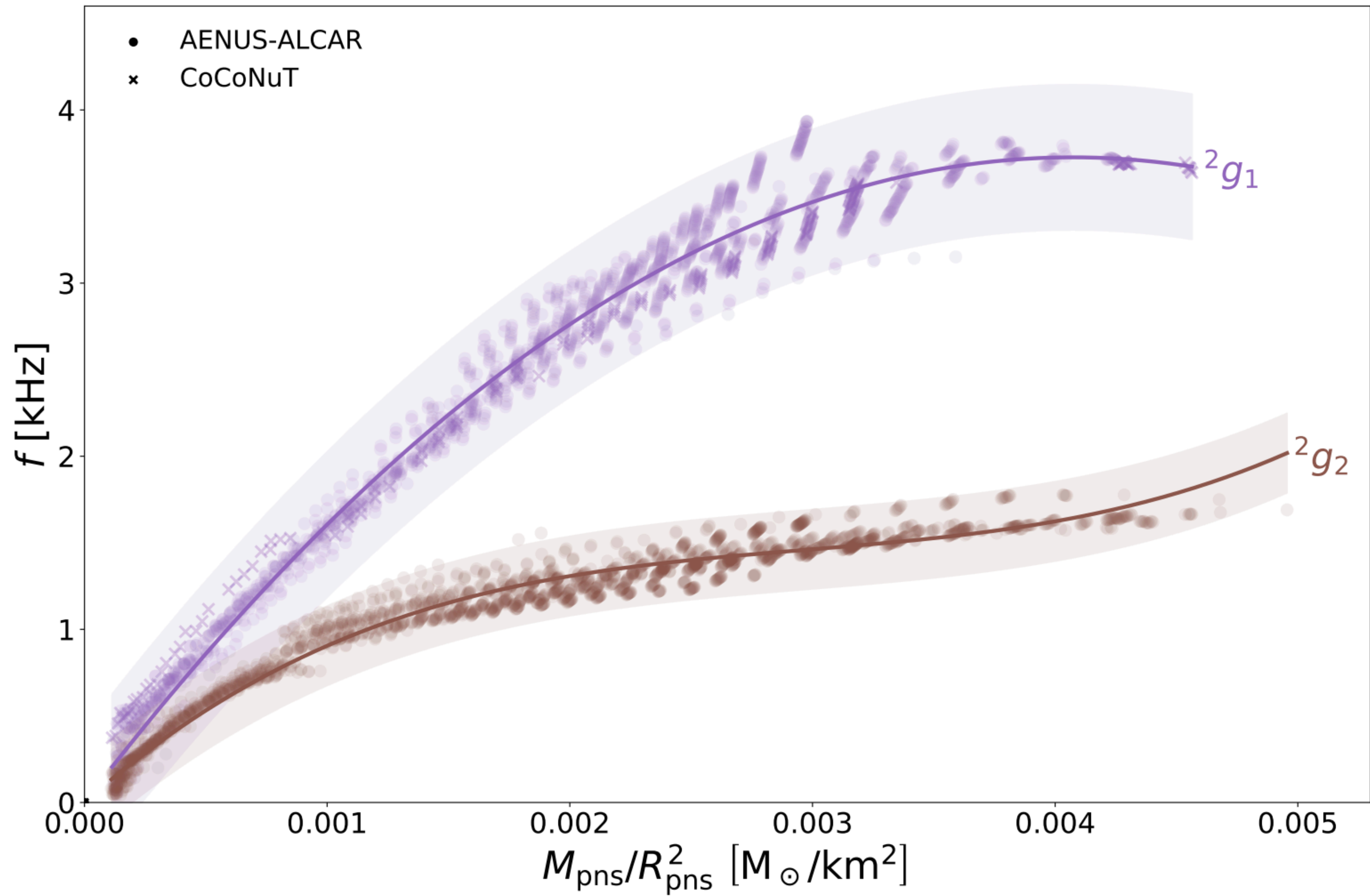
$$\xi^\theta = \eta_\perp \frac{1}{r^2} \partial_\theta Y_{lm} e^{-i\sigma t},$$

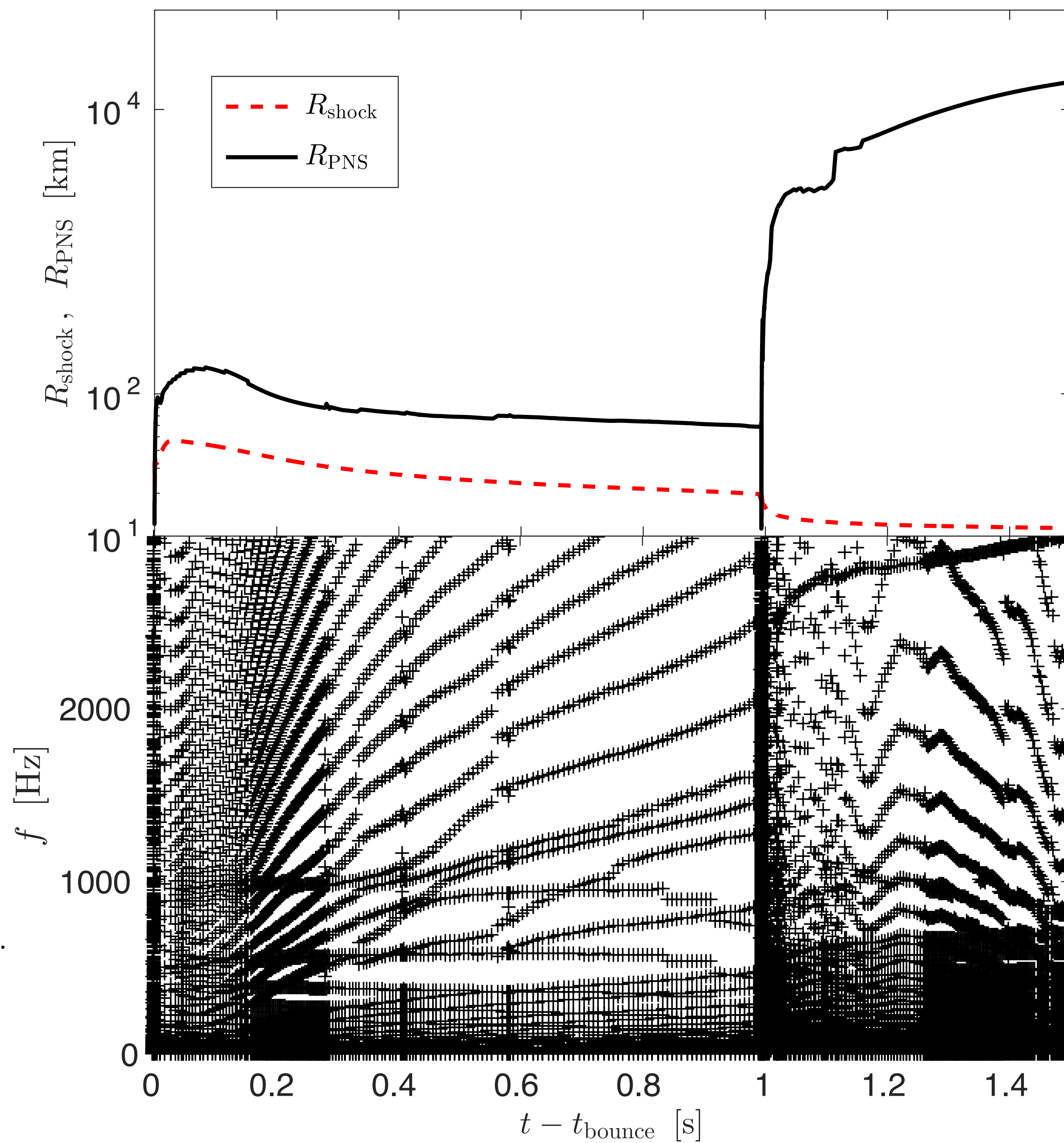
$$\xi^\varphi = \eta_\perp \frac{1}{r^2 \sin^2 \theta} \partial_\varphi Y_{lm} e^{-i\sigma t}$$

Torres-Forne et al. (2019), Mon. Not. Roy. Astron. Soc.

Torres-Forne et al. (2018), Mon. Not. Roy. Astron. Soc.







Model: s15.0
 Woosley et al. (2002), Rev. Mod. Phys.

EoS : DD2
 Typel S. et al., 2010, Phys. Rev. C

TERIMA KASIH
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MERCI
köszönöm
K SALAMAT
OBRIGADO
GRAZIE
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DZIĘKUJĘ
감사합니다
DANK U
DANKIE
TƏŞƏKKÜR
DĚKUJI
СПАСИБО
THANK YOU
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MULȚUMESC
TACK
FALEMNDERIT
PAKKA PÉR