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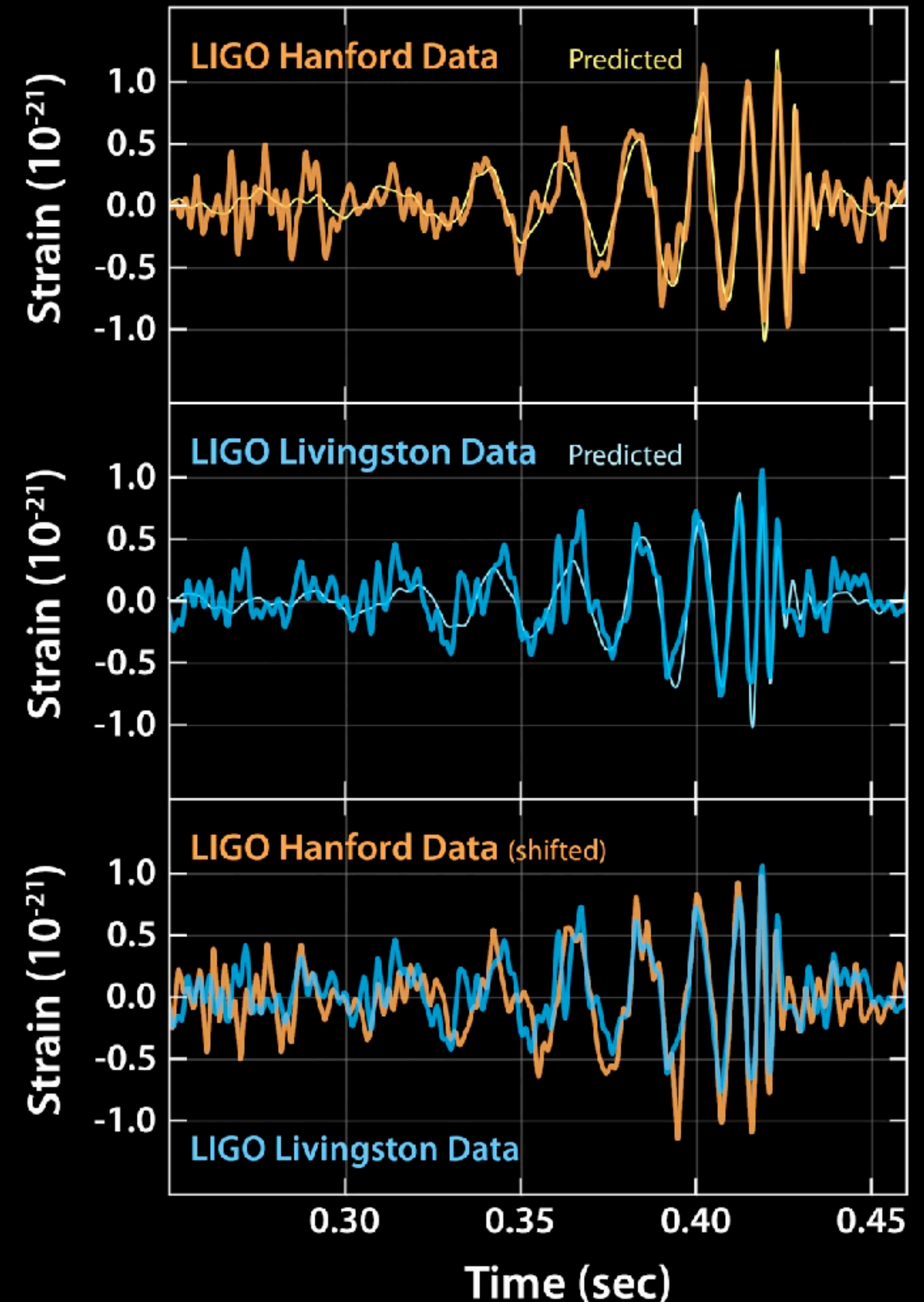


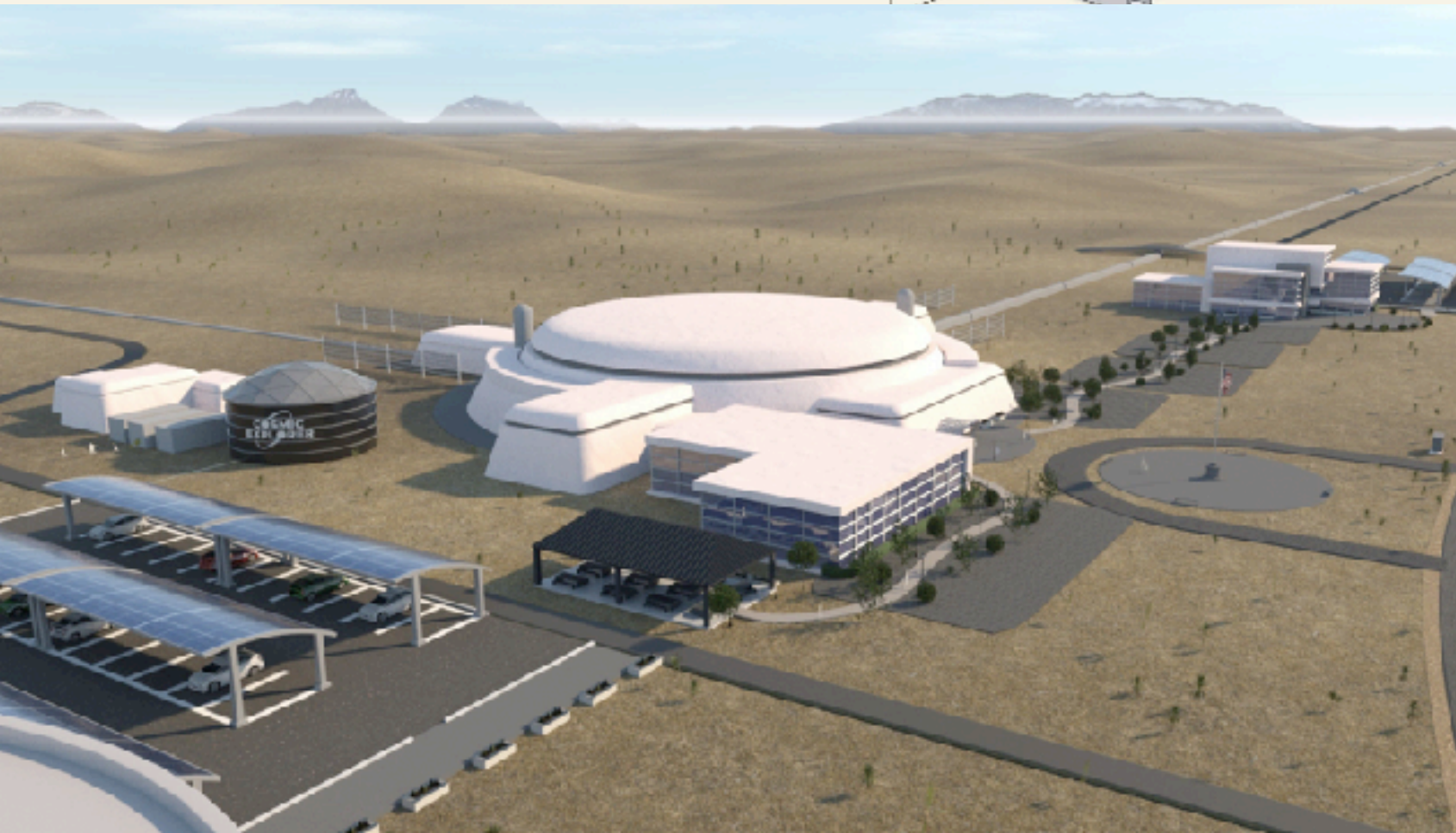
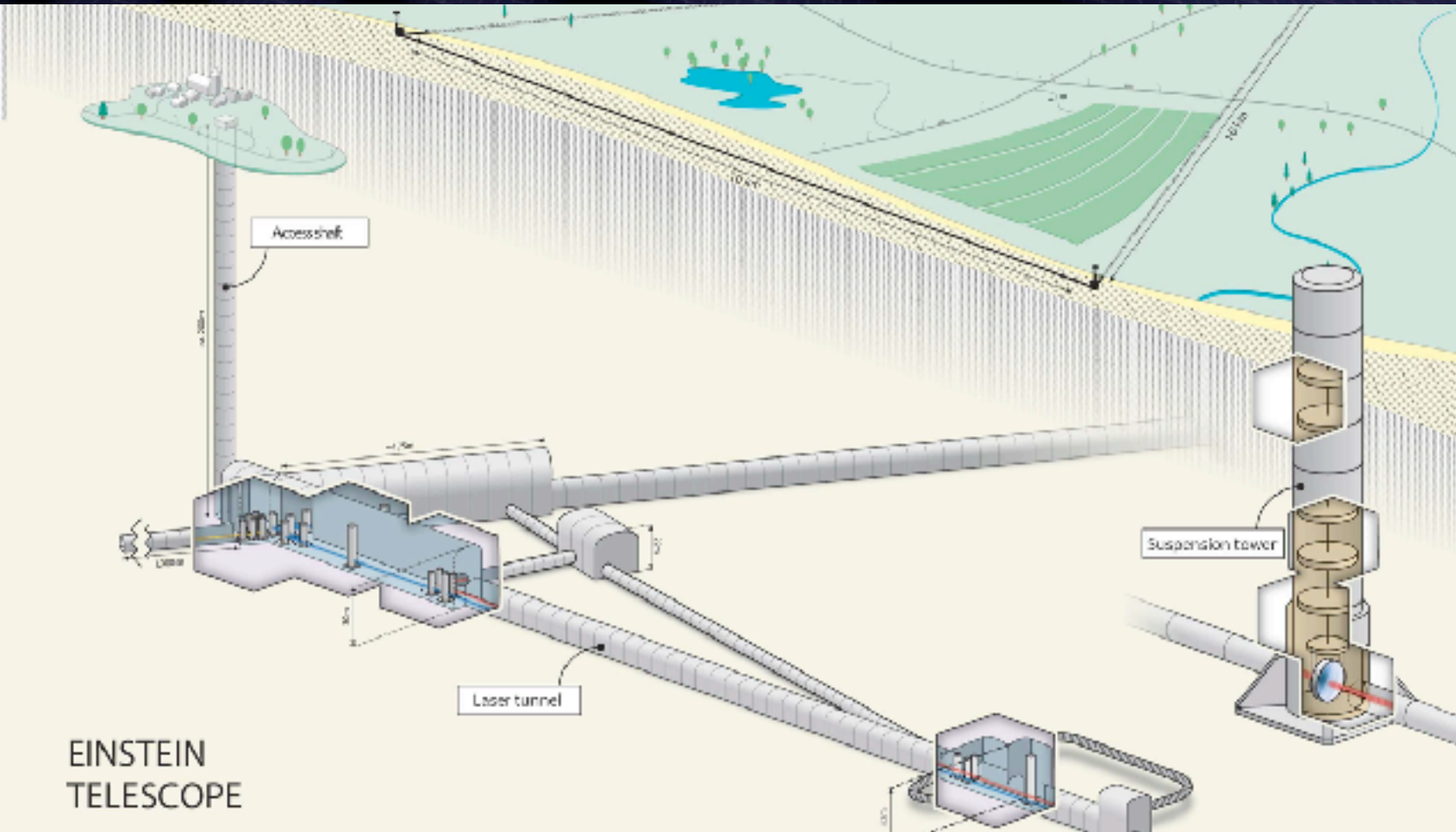
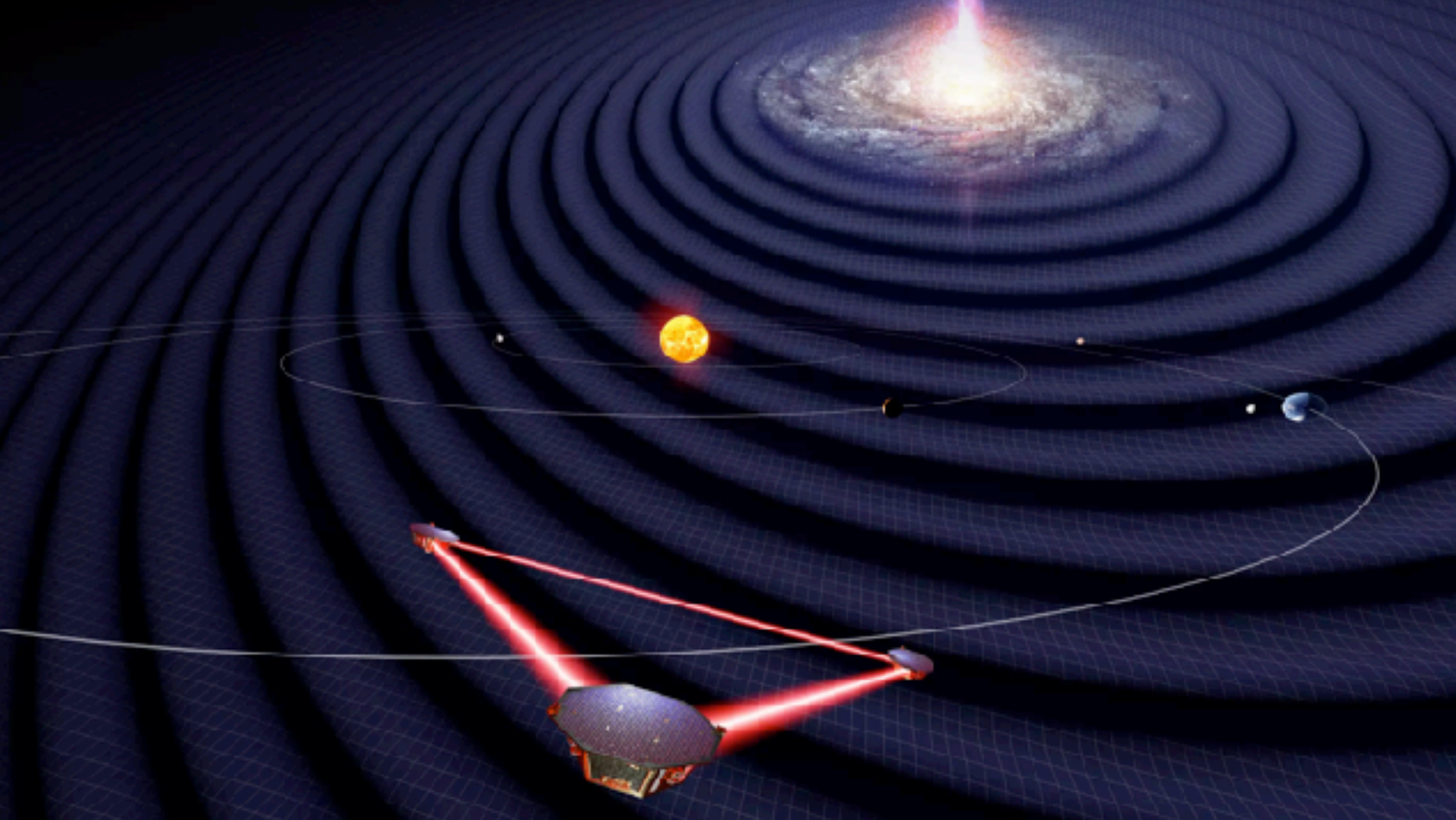
**Numerical Analysis of Mismatch Between
Eccentric Gravitational Wave Modeling
Numerical Codes
on a HTCondor cluster**

Balázs Kacs Kovics

Algorithms Used in GW Search

- In the LAL library (used by the mainstream): Numerical, EOB, Taylor waveforms
- Excellently modeling the currently detectable waveforms
- Problems:
 - ➔ Only short waveforms
 - ➔ Only specific waveforms
 - ➔ No eccentric waveforms
 - ➔ Mostly spin-aligned





New GW Detectors on the Horizon

eLISA, Einstein Telescope, Cosmic Explorer

- Targeting new sources like **NS-NS binaries**, merging galactic nuclei, supernovae, stochastic background
- Significantly longer observational times \Rightarrow longer waveforms (up to **6 months** — eLISA)
- Research of the inspiral phase
- **Eccentricity and spin effects** will be important in the orbital evolution of compact binaries
- **eLISA** got the **green light** this year

Gravitational Waves

Linearized theory

- Starting from the Einstein equation

$$R_{ab} - \frac{1}{2}g_{ab}R = \frac{8\pi G}{c^4}T_{ab}$$

- Take a small **perturbation** of the Einstein eq. **around a flat spacetime** (gauge symmetry of GR)

$$g_{ab} = \eta_{ab} + h_{ab}, \quad h_{ab} \ll 1$$

- The **Riemann-tensor** expressed in h_{ab} linear order

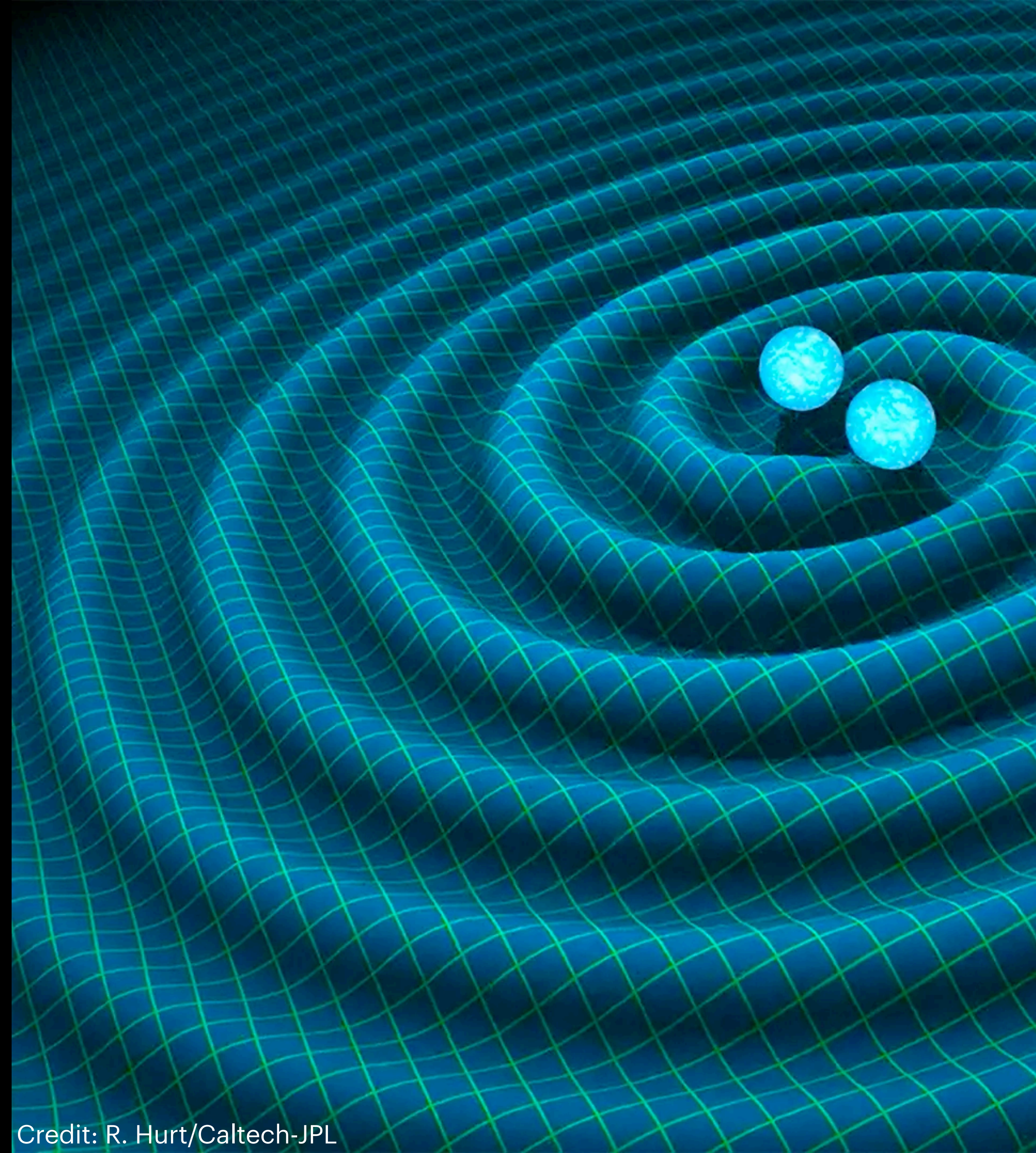
$$R_{abcd} = \frac{1}{2}(\partial_b\partial_c h_{ad} + \partial_a\partial_d h_{bc} - \partial_a\partial_c h_{bd} - \partial_b\partial_d h_{ac})$$

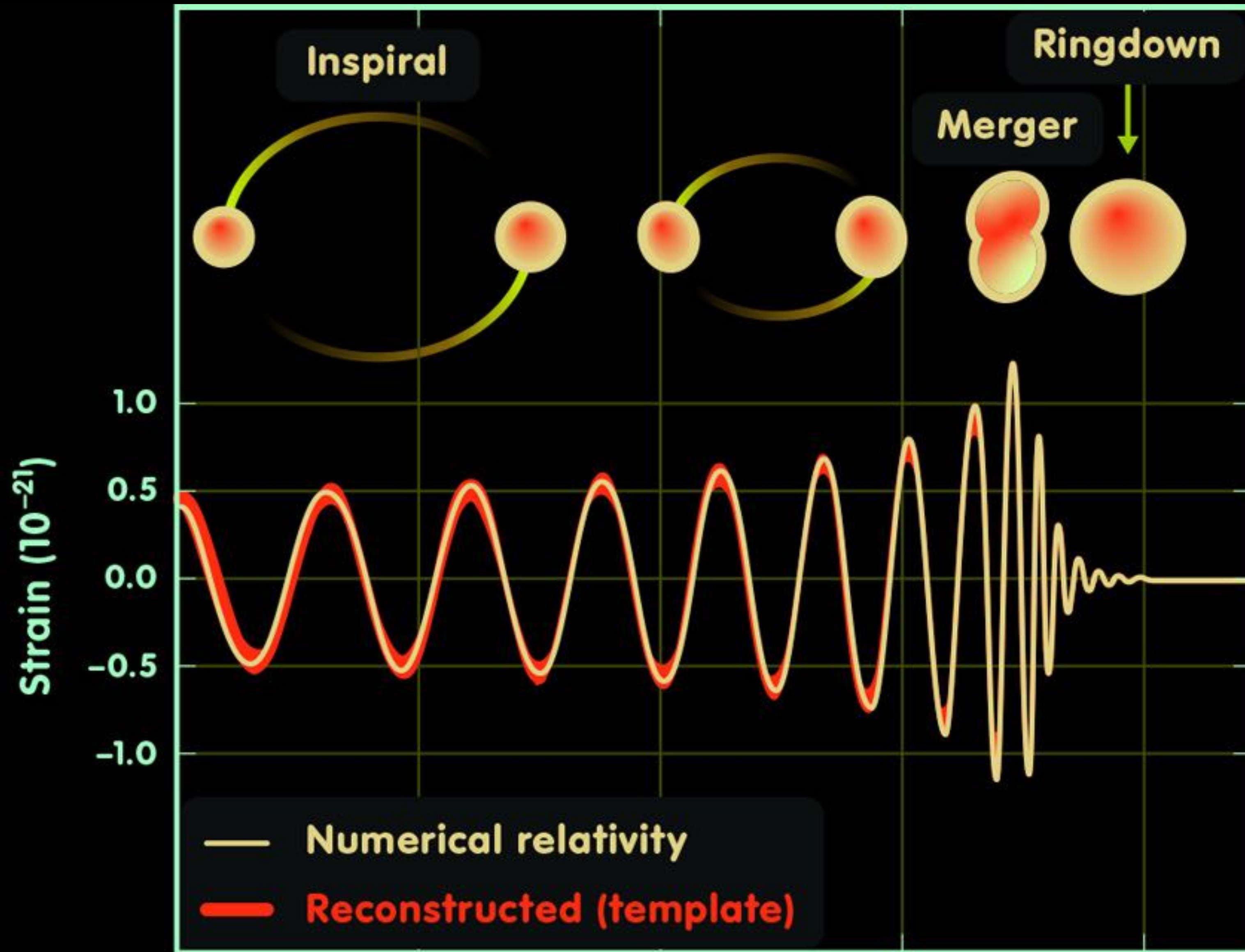
- The linearized Einstein equation

$$\square \bar{h}_{ab} + \eta_{ab}\partial^c\partial^d\bar{h}_{cd} - \partial^c\partial_b\bar{h}_{ac} - \partial^c\partial_a\bar{h}_{bc} = -\frac{16\pi G}{c^4}T_{ab}$$

- Using the gauge freedom of GR and choosing the **De Donder gauge**, $\partial^b\bar{h}_{ab} = 0$

$$\square \bar{h}_{ab} = -\frac{16\pi G}{c^4}T_{ab}$$

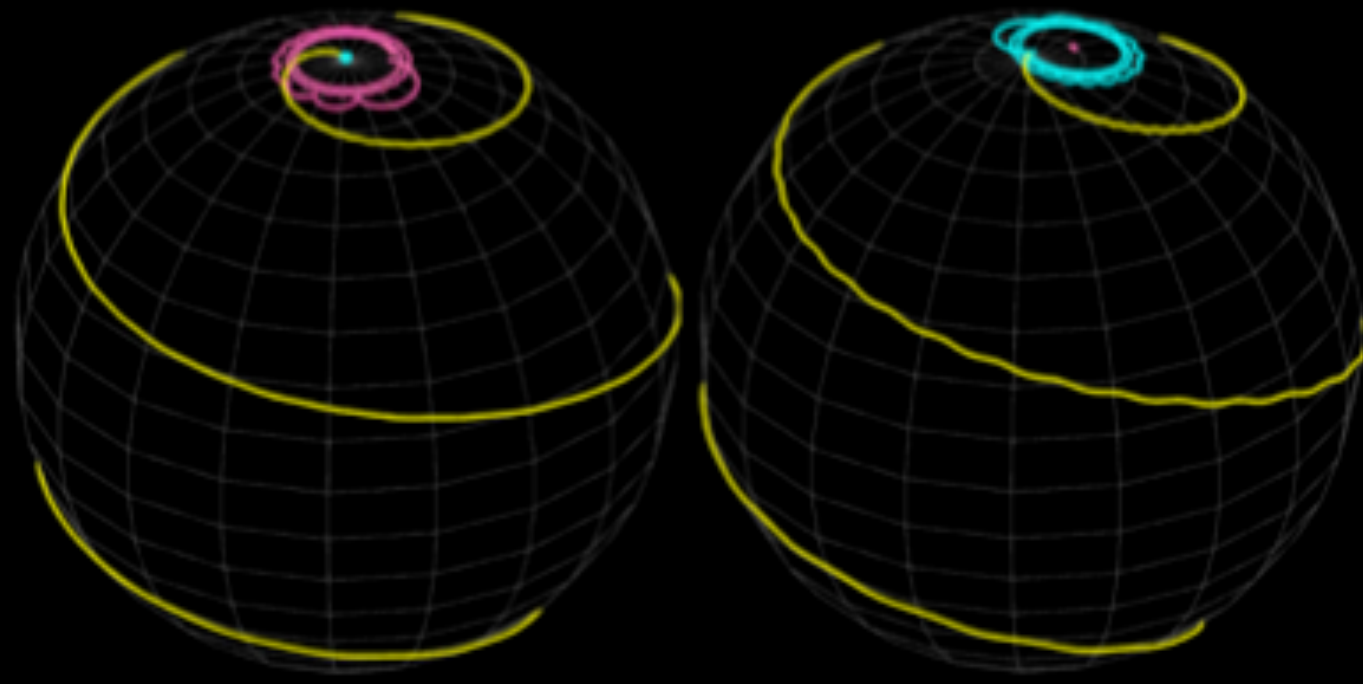




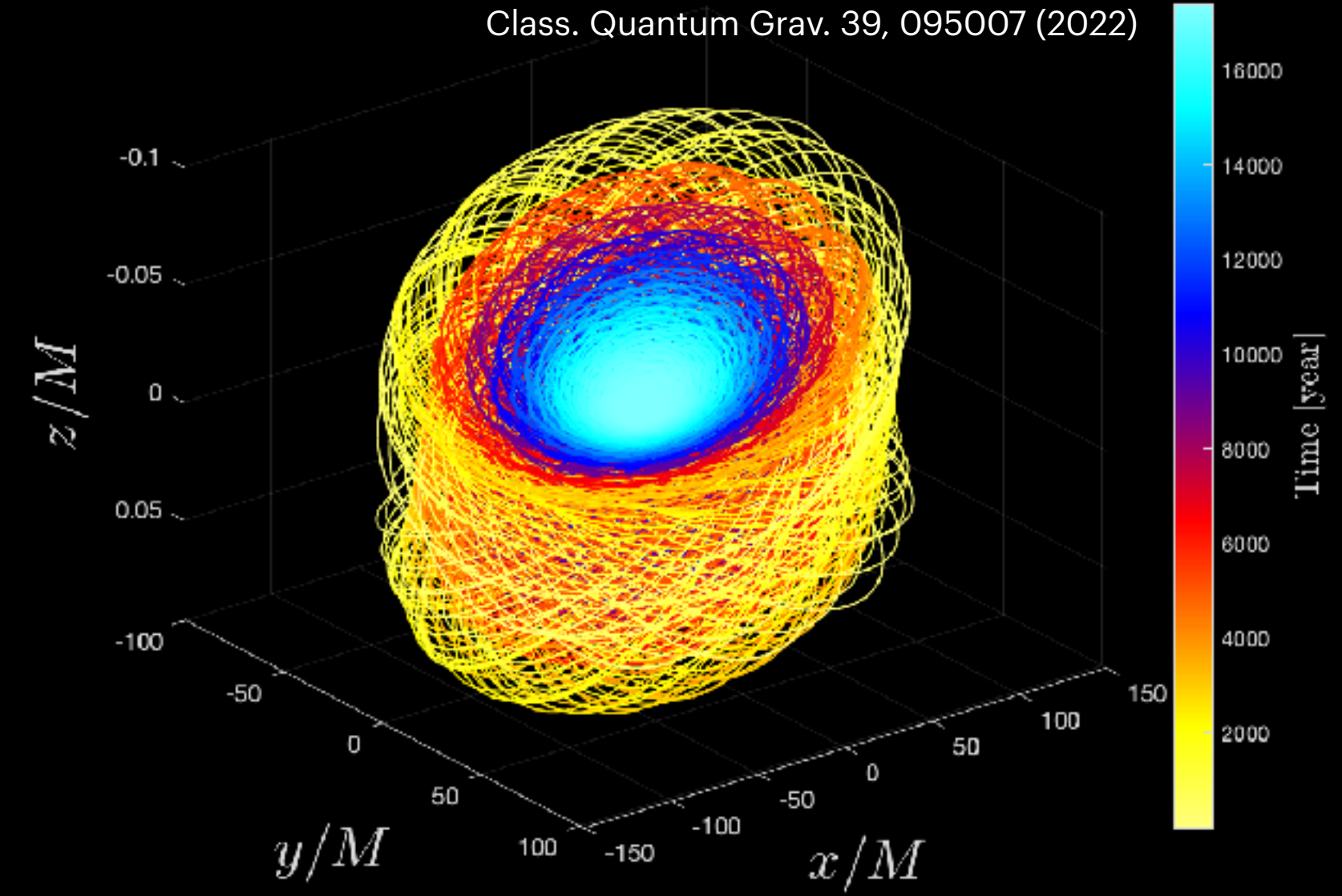
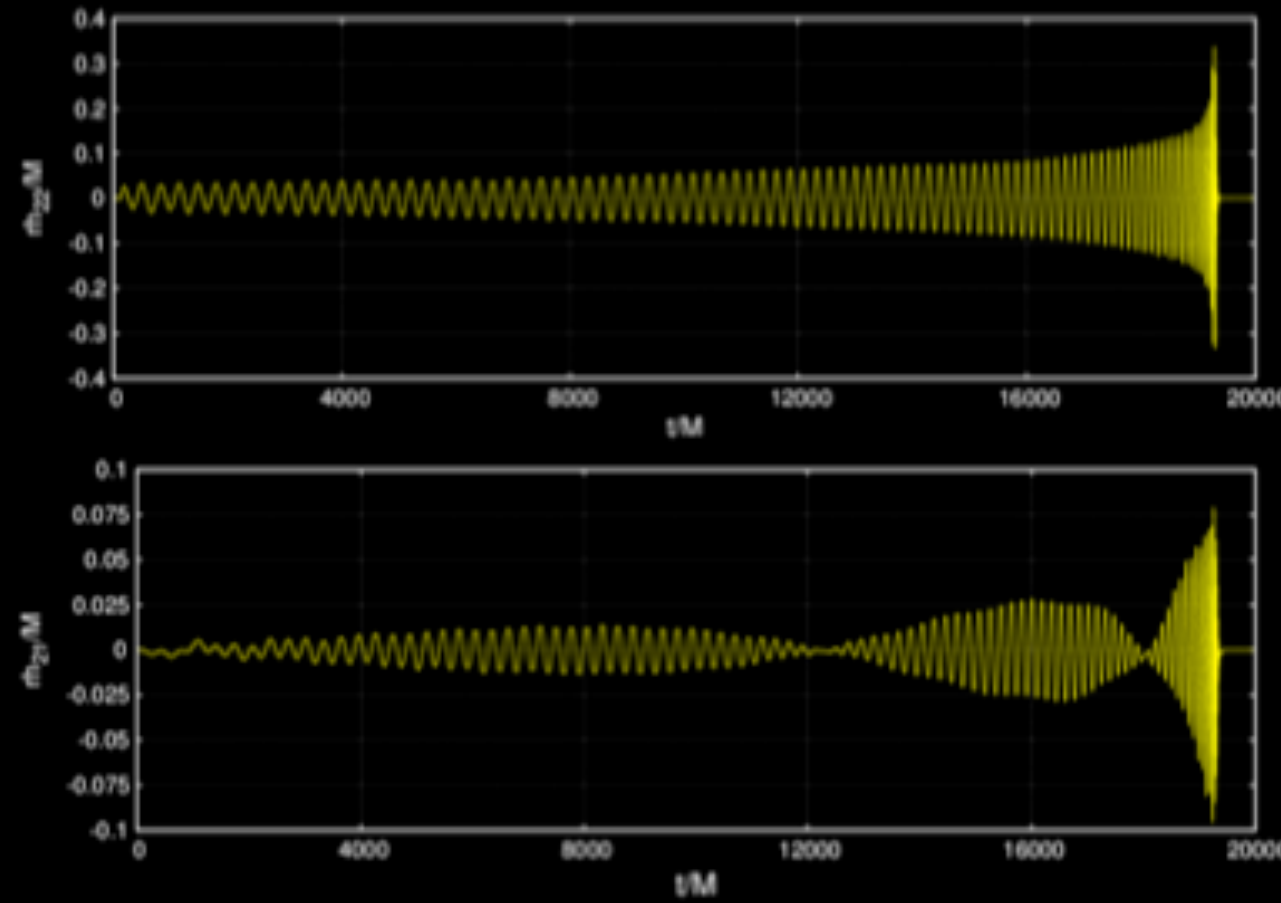
Motivation

Class. Quantum Grav. 39, 095007 (2022)

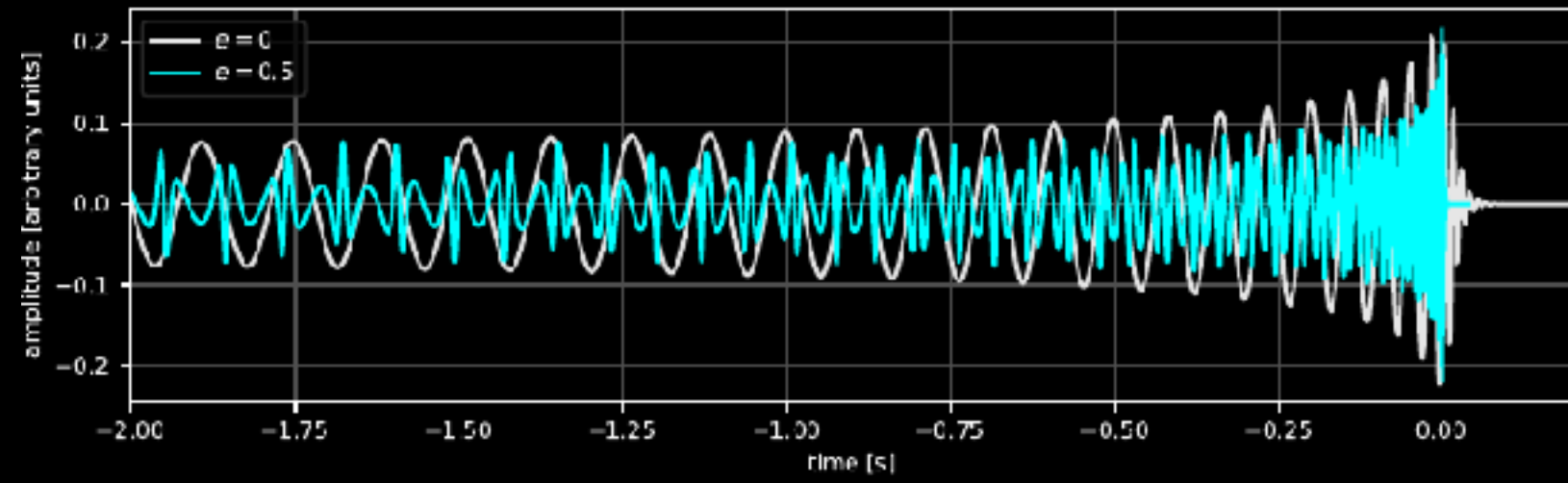
PRL 114, 141101 (2015)



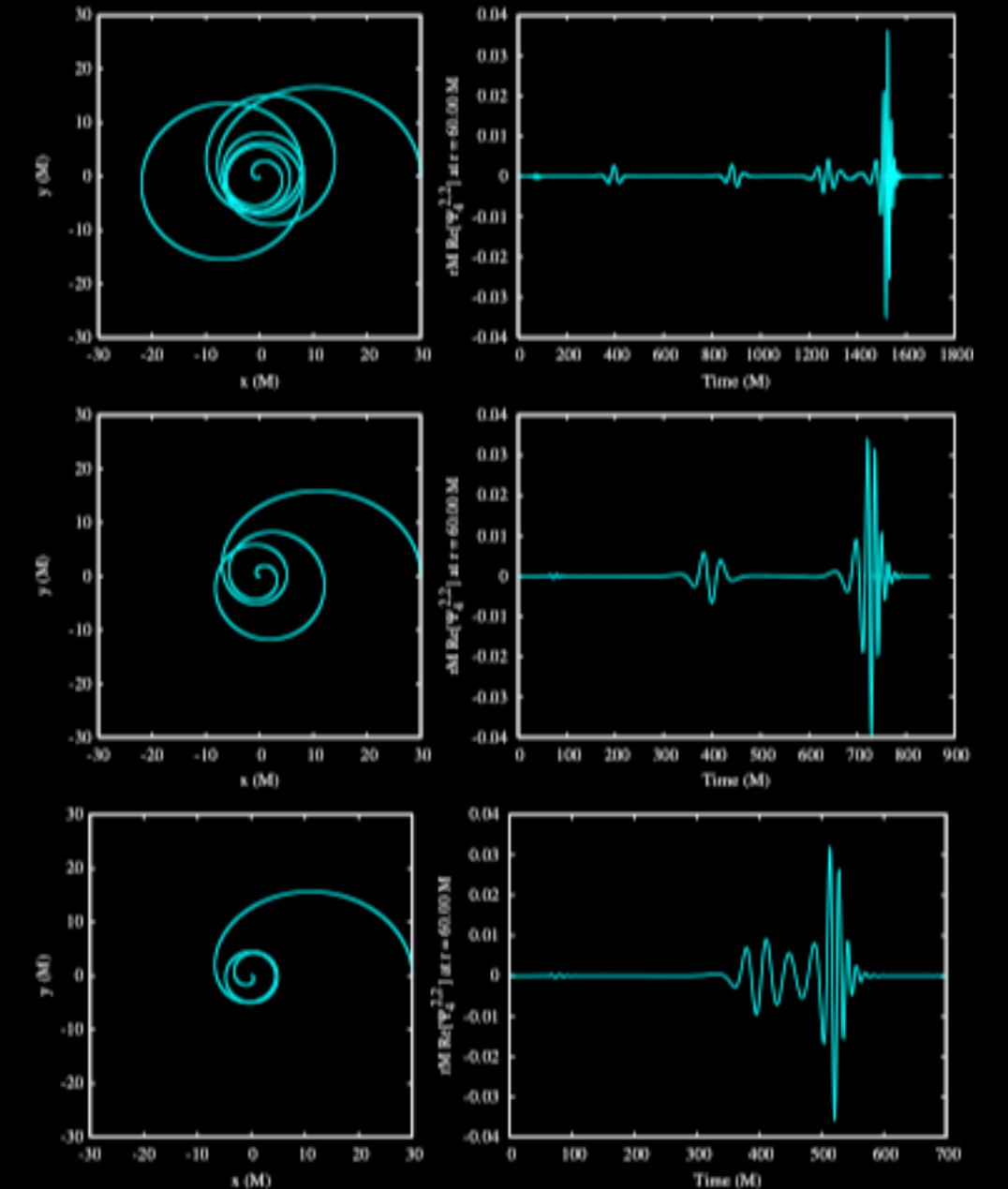
PRL 114, 141101 (2015)



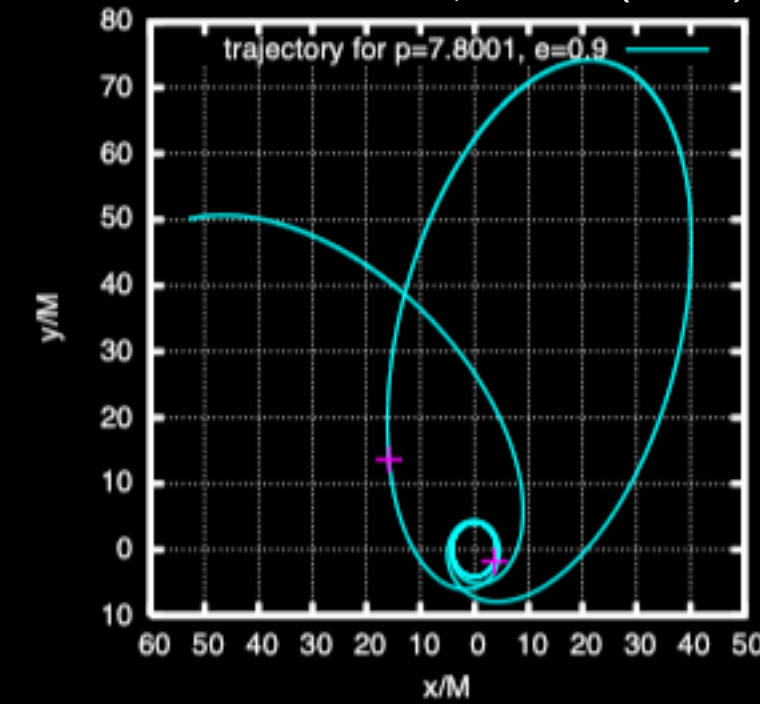
ApJ 883, 149 (2019)



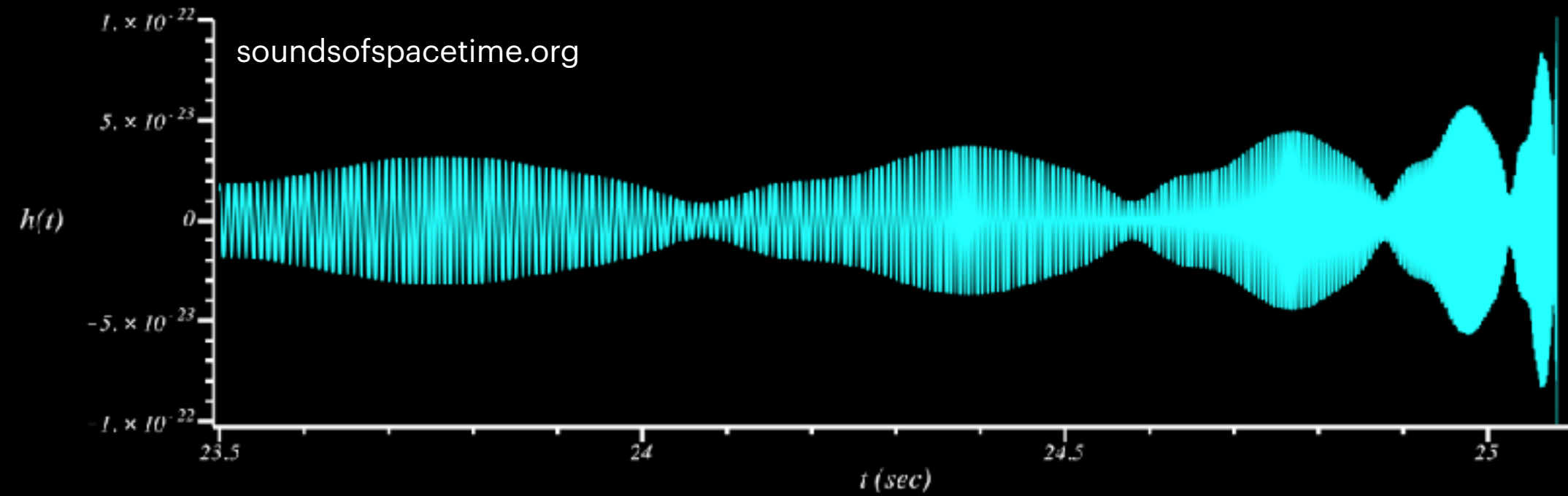
PRL 103, 131101 (2009)



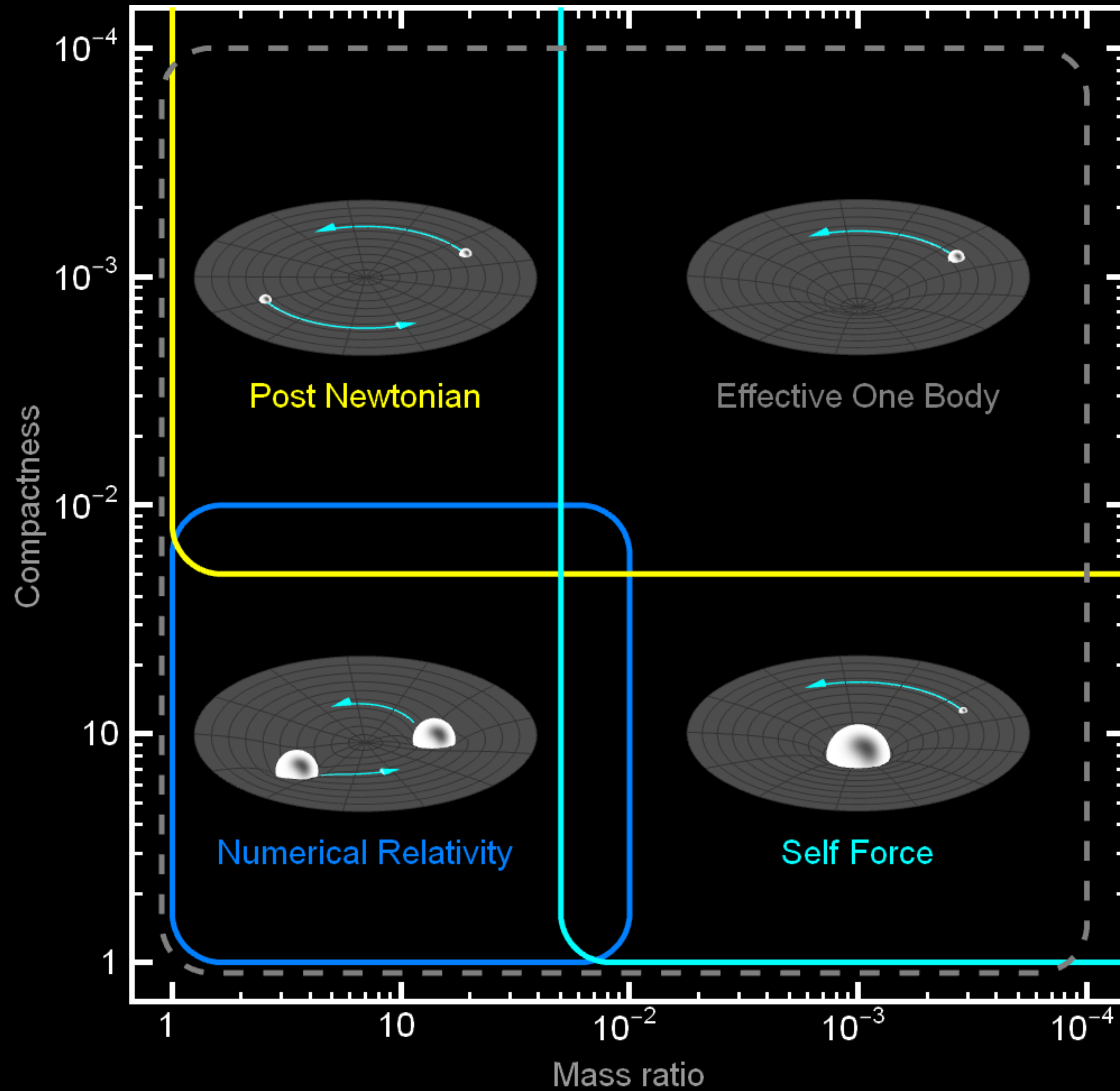
PRD 75, 124011 (2007)



sounds spacetime.org



Two-body problem of General Relativity



Post-Newtonian Expansion

- Built upon two assumptions:
 - gravity inside the source is weak like in the post-Minkowskian expansion
 - the motion of the components of the source is slow

- The equation of motion

$$\mathbf{a} = \mathbf{a}_N + \mathbf{a}_{PN} + \mathbf{a}_{2PN} + \mathbf{a}_{3PN} + \mathbf{a}_{4PN} + \mathbf{a}_{SO}^{1.5PN} + \mathbf{a}_{SS}^{2PN} + \mathbf{a}_{BT}^{RR, 2.5PN} + \mathbf{a}_{SO}^{2.5PN}$$

$$+ \mathbf{a}_{SO}^{3.5PN} + \mathbf{a}_{BT}^{RR, 3.5PN} + \mathbf{a}_{SS}^{RR, 3.5PN} + \mathbf{a}_{SO}^{RR, 3.5PN}$$

- The radiation field equation

$$h_{ij} = \frac{2G\mu}{c^4 D} [Q_{ij} + P^{0.5} Q_{ij} + P Q_{ij} + P^{1.5} Q_{ij} + P^2 Q_{ij} + P Q_{ij}^{SO}$$

$$+ P^{1.5} Q_{ij}^{SO} + P^2 Q_{ij}^{SO} + P Q_{ij}^{SS} + P^{1.5} Q_{ij}^{tail}]$$

Effective One-Body Approach

- reduce the conservative dynamics of the general relativistic two-body problem
- Mathisson–Papapetrou–Dixon equation is taken on a deformed Kerr black hole
- Hamiltonian of the Mathisson–Papapetrou–Dixon equations:

$$H_{\text{eff}} = M\eta \left(\beta^i p_i + \alpha \sqrt{1 + \gamma^{ij} p_i p_j + Q_4(p)} + H_S \right) + H_{SC}$$

$$H = M \sqrt{1 + 2\eta \left(\frac{H_{\text{eff}}}{M\eta} - 1 \right)}$$

- In the EOBNR framework, the quasicircular part of the radiation field is divided into two:
 - ❖ the inspiral-plunge
 - ❖ post-merger phase

$$h_{lm}^{(C)} = h_{lm}^{(N,\epsilon)} \hat{S}_{\text{eff}}^{(\epsilon)} T_{lm} e^{i\delta_{lm}} (\rho_{lm})^l N_{lm}$$
$$h_{lm}^{(N,\epsilon)} = \frac{M\eta}{D} n_{lm}^{(\epsilon)} c_{l+\epsilon} V_{\Phi}^l Y^{l-\epsilon, -m} \left(\frac{\pi}{2}, \Phi \right)$$

- For the eccentric part, in the radiation field terms up to the second post-Newtonian order are considered

Numerical Results

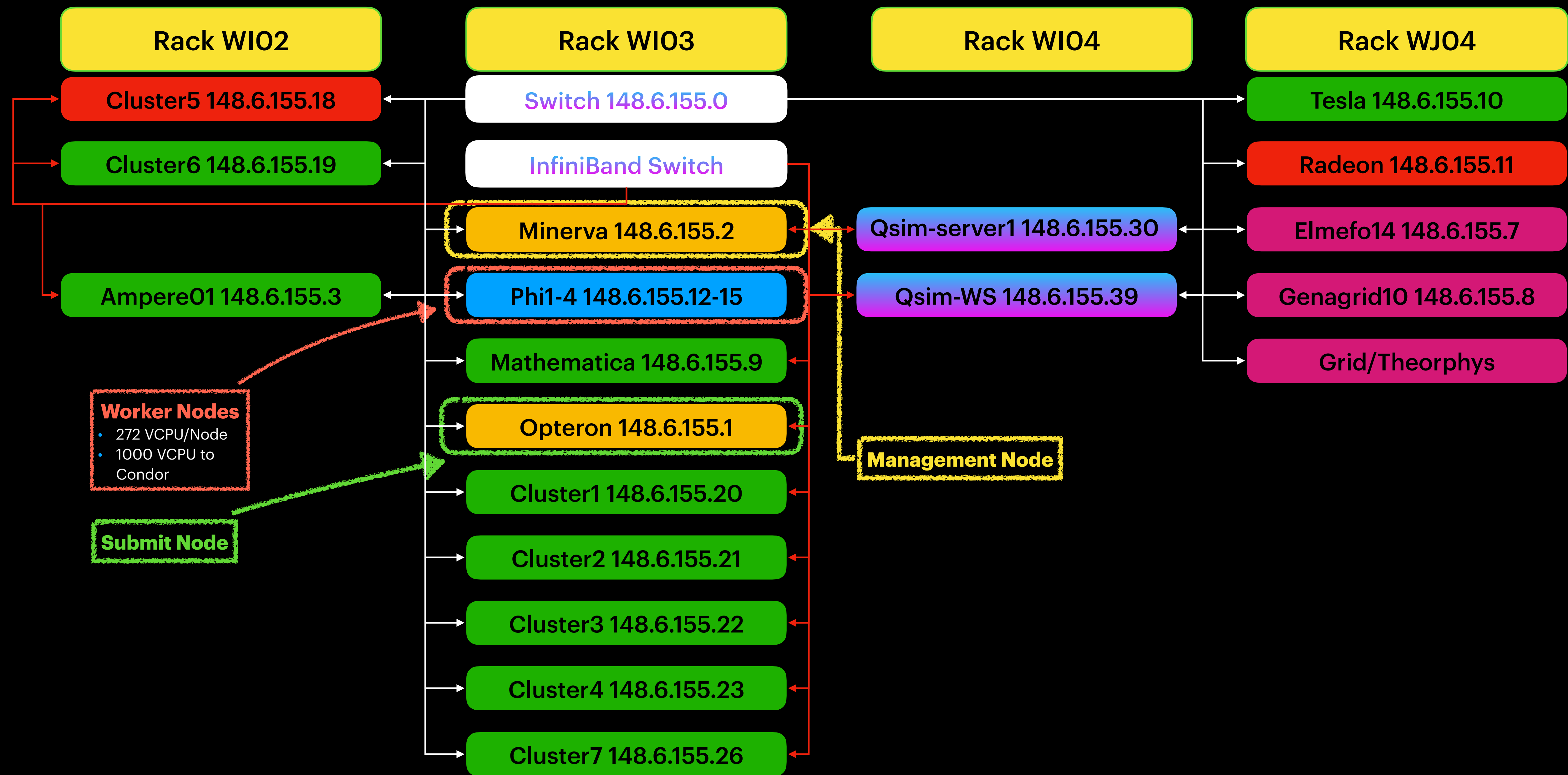
- 2 codes were used; one based on the PN, **CBwaves**; and one based on EOB, **SEOBNRE**
- both codes use a **4th-order Runge—Kutta** integrator
- on an identical initial parameter space

Initial Parameters

$m_1 [M_\odot]$	10 ... 100
$m_2 [M_\odot]$	10 ... 100
$R [M_{\text{tot}}]$	30
$R_{\text{min}} [M_{\text{tot}}]$	6
e_0	0.003
$dt [\text{sec}]$	1/4096

- **SEOBNRE** uses the initial orbital frequency:
$$f_{\text{init}} = \frac{c^3}{\pi G(m_1 + m_2)M_\odot \sqrt{r_0^3}}$$

Technical Background



Setting Up a Dynamic HTCondor Cluster

- 3 schemes for slot distribution
 1. Static
 2. Partionable
 3. Dynamic
- One **partionable** slot
- **Dynamic** slots are created from it
- **Dynamic** slots merge back to the **partionable** slot upon finishing
- The partionable slots split on: the **CPU, Memory, Disk**
- **How to configure:**

```
NUM_SLOTS = 1
NUM_SLOTS_TYPE_1 = 1
SLOT_TYPE_1 = cpus=100%
SLOT_TYPE_1_PARTITIONABLE = true
```
- Documentation to install HTCondor:

<https://htcondor.readthedocs.io/en/latest/admin-manual/index.html>

Edit it to change the number of CPUs used

Name	OpSys	Arch	State	Activity	LoadAv	Mem	ActvtyTime		
slot1@ph1.gpu.wigner.hun-ren.hu	LINUX	X86_64	Unclaimed	Idle	0.000	64244	54+17:27:35		
slot1@phi2.gpu.wigner.hun-ren.hu	LINUX	X86_64	Unclaimed	Idle	0.000	64226	186+18:10:59		
slot1@phi3.gpu.wigner.hun-ren.hu	LINUX	X86_64	Unclaimed	Idle	0.000	64226	186+18:08:56		
slot1@phi4.gpu.wigner.hun-ren.hu	LINUX	X86_64	Unclaimed	Idle	0.000	64226	186+18:09:48		
	Total	Owner	Claimed	Unclaimed	Matched	Preempting	Drain	Backfill	BkIdle
X86_64/LINUX	4	0	0	4	0	0	0	0	0
Total	4	0	0	4	0	0	0	0	0

Name	OpSys	Arch	State	Activity	LoadAv	Mem	ActvtyTime		
slot1@ph1.wigner.hun-ren.hu	LINUX	X86_64	Unclaimed	Idle	0.000	63220	54+18:10:11		
slot1_1@ph1.wigner.hun-ren.hu	LINUX	X86_64	Claimed	Busy	0.000	128	0+00:00:01		
slot1_2@ph1.wigner.hun-ren.hu	LINUX	X86_64	Claimed	Busy	0.000	128	0+00:00:01		
slot1_3@ph1.wigner.hun-ren.hu	LINUX	X86_64	Claimed	Busy	0.000	128	0+00:00:00		
slot1_4@ph1.wigner.hun-ren.hu	LINUX	X86_64	Claimed	Busy	0.000	128	0+00:00:00		
slot1_5@ph1.wigner.hun-ren.hu	LINUX	X86_64	Claimed	Busy	0.000	128	0+00:00:01		
slot1_6@ph1.wigner.hun-ren.hu	LINUX	X86_64	Claimed	Busy	0.000	128	0+00:00:00		
slot1_7@ph1.wigner.hun-ren.hu	LINUX	X86_64	Claimed	Busy	0.000	128	0+00:00:01		
slot1_8@ph1.wigner.hun-ren.hu	LINUX	X86_64	Claimed	Busy	0.000	128	0+00:00:00		
slot1@phi2.gpu.wigner.hun-ren.hu	LINUX	X86_64	Unclaimed	Idle	0.000	63202	186+18:53:00		
slot1_1@phi2.gpu.wigner.hun-ren.hu	LINUX	X86_64	Claimed	Busy	0.000	128	0+00:00:01		
slot1_2@phi2.gpu.wigner.hun-ren.hu	LINUX	X86_64	Claimed	Busy	0.000	128	0+00:00:01		
slot1_3@phi2.gpu.wigner.hun-ren.hu	LINUX	X86_64	Claimed	Busy	0.000	128	0+00:00:00		
slot1_4@phi2.gpu.wigner.hun-ren.hu	LINUX	X86_64	Claimed	Busy	0.000	128	0+00:00:00		
slot1_5@phi2.gpu.wigner.hun-ren.hu	LINUX	X86_64	Claimed	Busy	0.000	128	0+00:00:00		
slot1_6@phi2.gpu.wigner.hun-ren.hu	LINUX	X86_64	Claimed	Busy	0.000	128	0+00:00:00		
slot1_7@phi2.gpu.wigner.hun-ren.hu	LINUX	X86_64	Claimed	Busy	0.000	128	0+00:00:01		
slot1_8@phi2.gpu.wigner.hun-ren.hu	LINUX	X86_64	Claimed	Busy	0.000	128	0+00:00:00		
slot1@phi3.gpu.wigner.hun-ren.hu	LINUX	X86_64	Unclaimed	Idle	0.000	63202	186+18:53:13		
slot1_1@phi3.gpu.wigner.hun-ren.hu	LINUX	X86_64	Claimed	Busy	0.000	128	0+00:00:01		
slot1_2@phi3.gpu.wigner.hun-ren.hu	LINUX	X86_64	Claimed	Busy	0.000	128	0+00:00:01		
slot1_3@phi3.gpu.wigner.hun-ren.hu	LINUX	X86_64	Claimed	Busy	0.000	128	0+00:00:00		
slot1_4@phi3.gpu.wigner.hun-ren.hu	LINUX	X86_64	Claimed	Busy	0.000	128	0+00:00:00		
slot1_5@phi3.gpu.wigner.hun-ren.hu	LINUX	X86_64	Claimed	Busy	0.000	128	0+00:00:00		
slot1_6@phi3.gpu.wigner.hun-ren.hu	LINUX	X86_64	Claimed	Busy	0.000	128	0+00:00:00		
slot1_7@phi3.gpu.wigner.hun-ren.hu	LINUX	X86_64	Claimed	Busy	0.000	128	0+00:00:01		
slot1_8@phi3.gpu.wigner.hun-ren.hu	LINUX	X86_64	Claimed	Busy	0.000	128	0+00:00:00		
slot1@phi4.gpu.wigner.hun-ren.hu	LINUX	X86_64	Unclaimed	Idle	0.000	63202	186+18:53:08		
slot1_1@phi4.gpu.wigner.hun-ren.hu	LINUX	X86_64	Claimed	Busy	0.000	128	0+00:00:01		
slot1_2@phi4.gpu.wigner.hun-ren.hu	LINUX	X86_64	Claimed	Busy	0.000	128	0+00:00:01		
slot1_3@phi4.gpu.wigner.hun-ren.hu	LINUX	X86_64	Claimed	Busy	0.000	128	0+00:00:00		
slot1_4@phi4.gpu.wigner.hun-ren.hu	LINUX	X86_64	Claimed	Busy	0.000	128	0+00:00:01		
slot1_5@phi4.gpu.wigner.hun-ren.hu	LINUX	X86_64	Claimed	Busy	0.000	128	0+00:00:00		
slot1_6@phi4.gpu.wigner.hun-ren.hu	LINUX	X86_64	Claimed	Busy	0.000	128	0+00:00:00		
slot1_7@phi4.gpu.wigner.hun-ren.hu	LINUX	X86_64	Claimed	Busy	0.000	128	0+00:00:01		
slot1_8@phi4.gpu.wigner.hun-ren.hu	LINUX	X86_64	Claimed	Busy	0.000	128	0+00:00:00		
	Total	Owner	Claimed	Unclaimed	Matched	Preempting	Drain	Backfill	BkIdle
X86_64/LINUX	36	0	32	4	0	0	0	0	0
Total	36	0	32	4	0	0	0	0	0

Submitting Jobs With Python

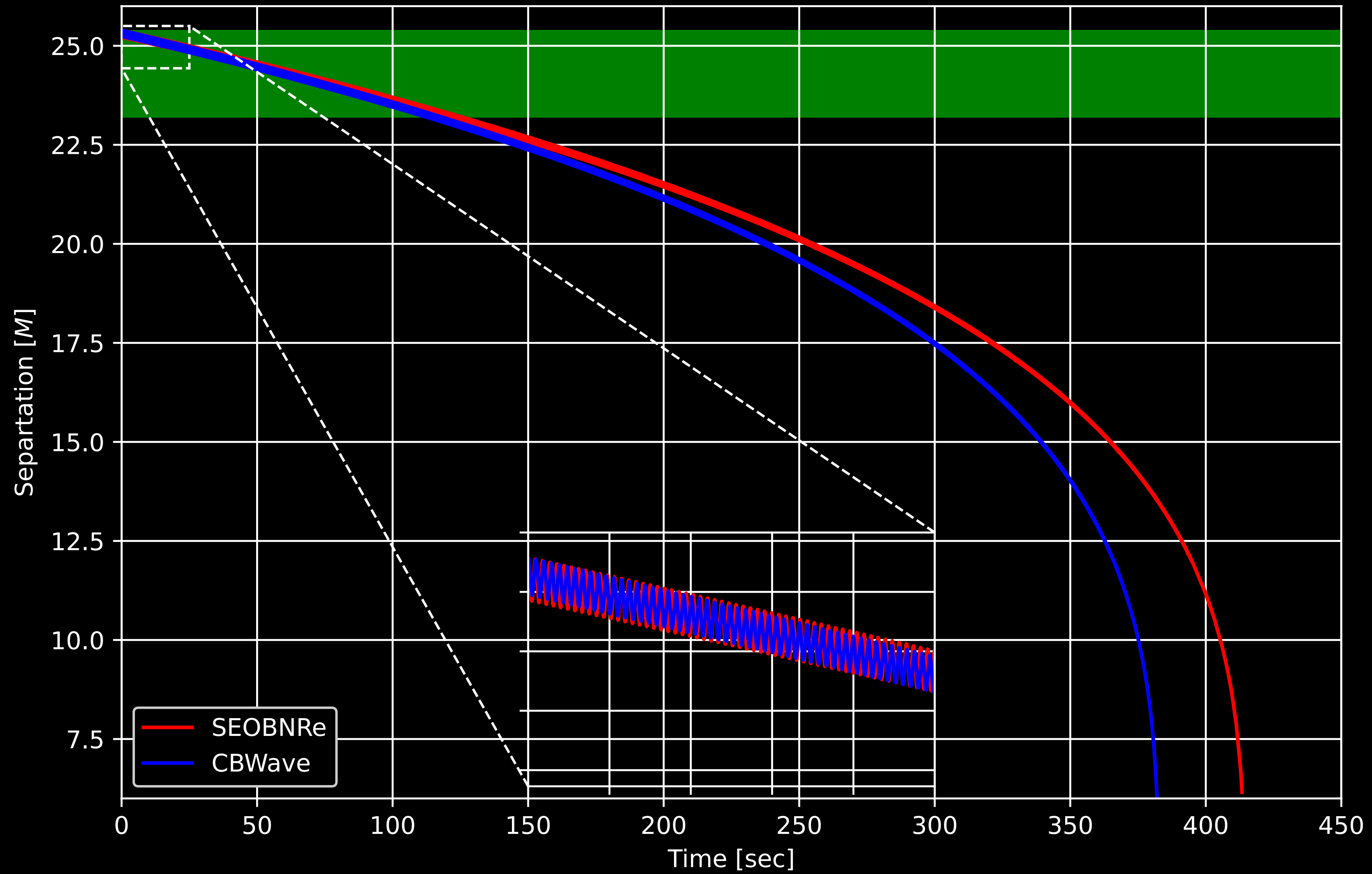
```
for j in range(m):
    for j in range(m):
        schedd = htc.Schedd()

        prefix = str(int(M1*np.power(10,acc))) + "_" + str(int(e0[j]*np.power(10,acc)))
        condor_cbw_out = "out/cbwaves_" + prefix + ".out"
        condor_seo_out = "out/seobnre_" + prefix + ".out"
        condor_cbw_err = "error/cbwaves_" + prefix + ".err"
        condor_seo_err = "error/seobnre_" + prefix + ".err"

        *** HTCondor submit for SEOBNRE ***
        cmd_seo = path + "/seobnre.py --m1 " + str(M1) + " --m2 " + str(M2) + " --e0 " + str(e0[j]) + " --sampleRate " + str(int(dt)) + " --spin1x " + str(S1x) + " --spin1y " + str(S1y) + " --spin1z " + str(S1z)
        + " --spin2x " + str(S2x) + " --spin2y " + str(S2y) + " --spin2z " + str(S2z) + " --fmin " + str(Fmin) + " --printstep " + str(printstep) + " --fnameprefix " + prefix + " --dir " + dirpath + " --seobin " +
        str(seobin)

        seobnre_job = htc.Submit({
            "initialdir": inidir,          # sets the location of the initial working directory
            "executable": "/usr/bin/python3.10",  # the program to run on the execute node
            "arguments": cmd_seo,          # the arguments of the program, that will run on the execute node
            "getenv": "True",
            "output": condor_seo_out,      # anything the job prints to standard output will end up in this file
            "error": condor_seo_err,       # anything the job prints to standard error will end up in this file
            "log": "seobnre.log",          # this file will contain a record of what happened to the job
            "request_cpus": "1",           # how many CPU cores we want
        })
        submit_seo = schedd.submit(seobnre_job)  # submit the job
```


Evolution of the orbital separation with 5 Hz initial orbital frequency at $q = 1/100$



Mismatch/Unfaithfulness

- To calculate the mismatch, one first has to calculate the **Overlap**:

$$\mathcal{O} = \frac{\langle h_1, h_2 \rangle}{\sqrt{\langle h_1, h_1 \rangle \langle h_2, h_2 \rangle}}$$

where

$$\langle h_1, h_2 \rangle = 4\Re \int_{f_{\max}}^{f_{\min}} \frac{\tilde{h}_1 \tilde{h}_2^*}{S_n(f)} df$$

- The **mismatch** (or unfaithfulness) is the **marginalized overlap** over some quantities

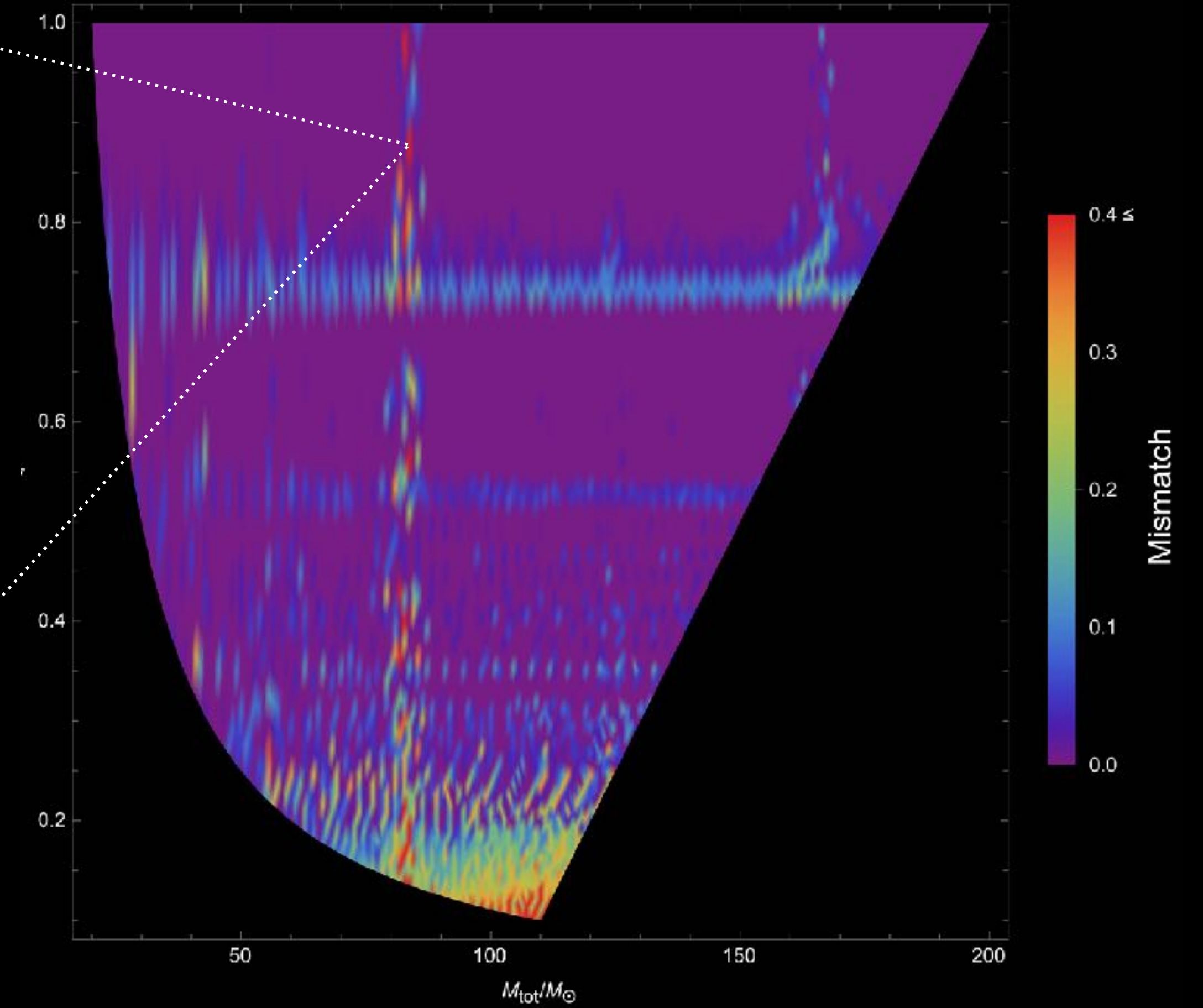
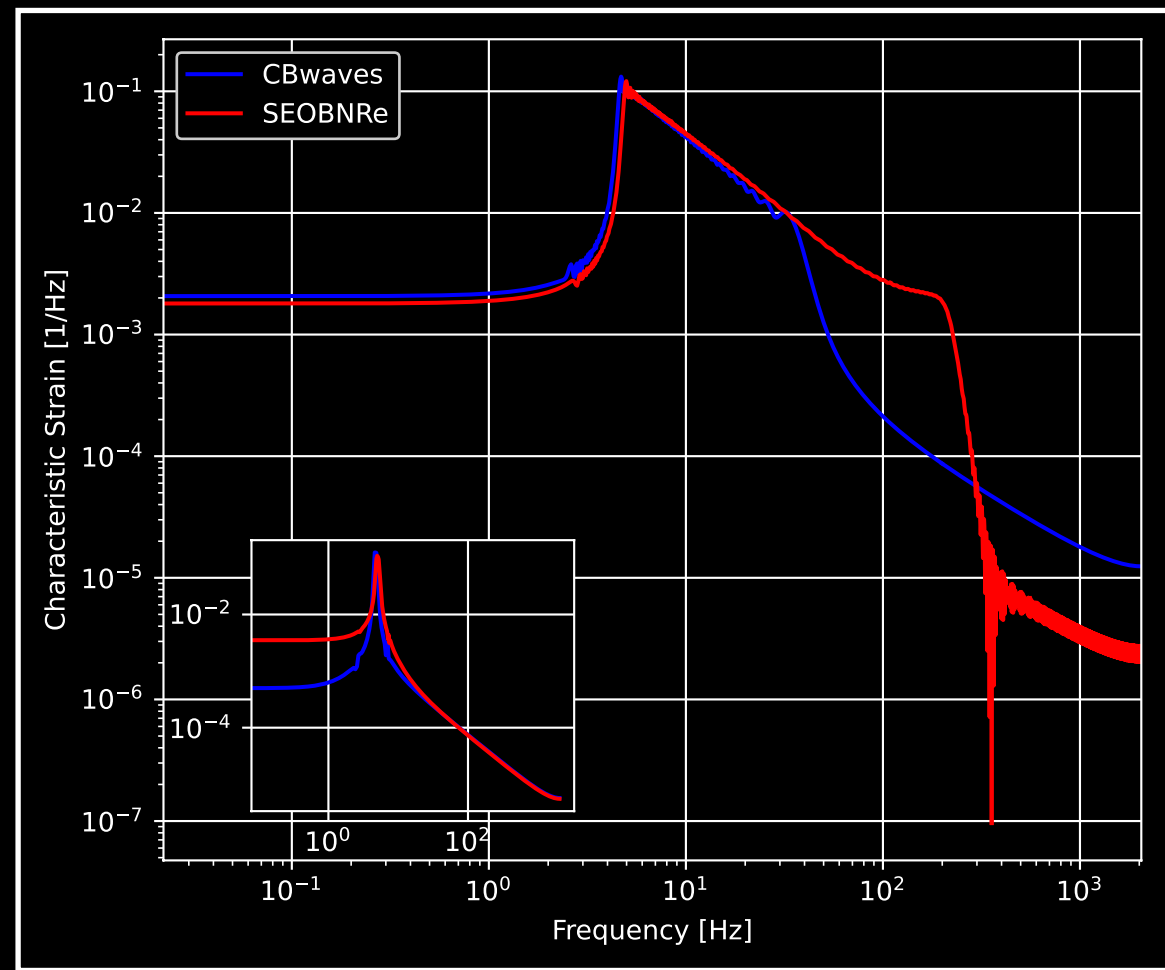
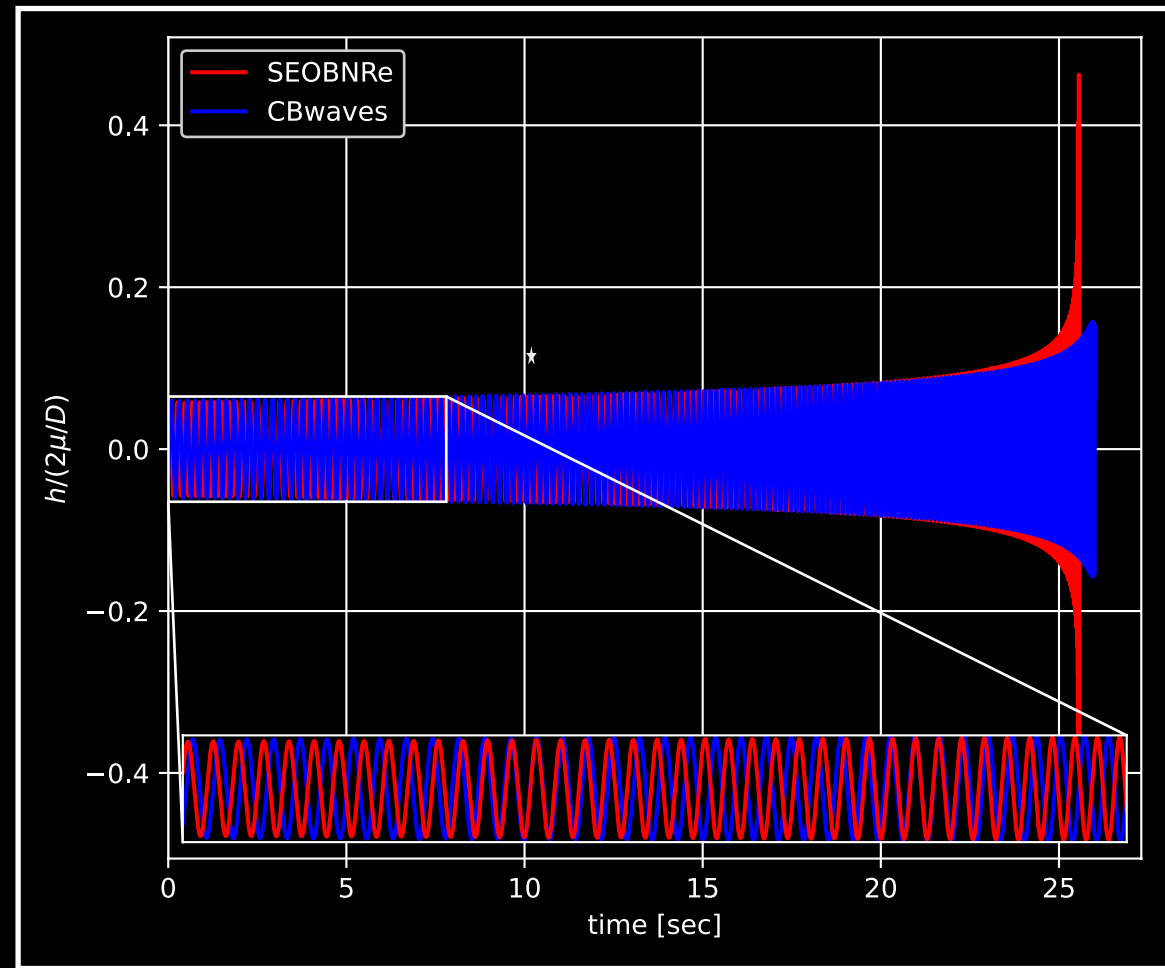
$$\mathcal{M} = \max_{t, \phi, \psi} \mathcal{O}(h_1, h_2)$$

where the **max** was taken over timeshifts, polarization angles, and phase

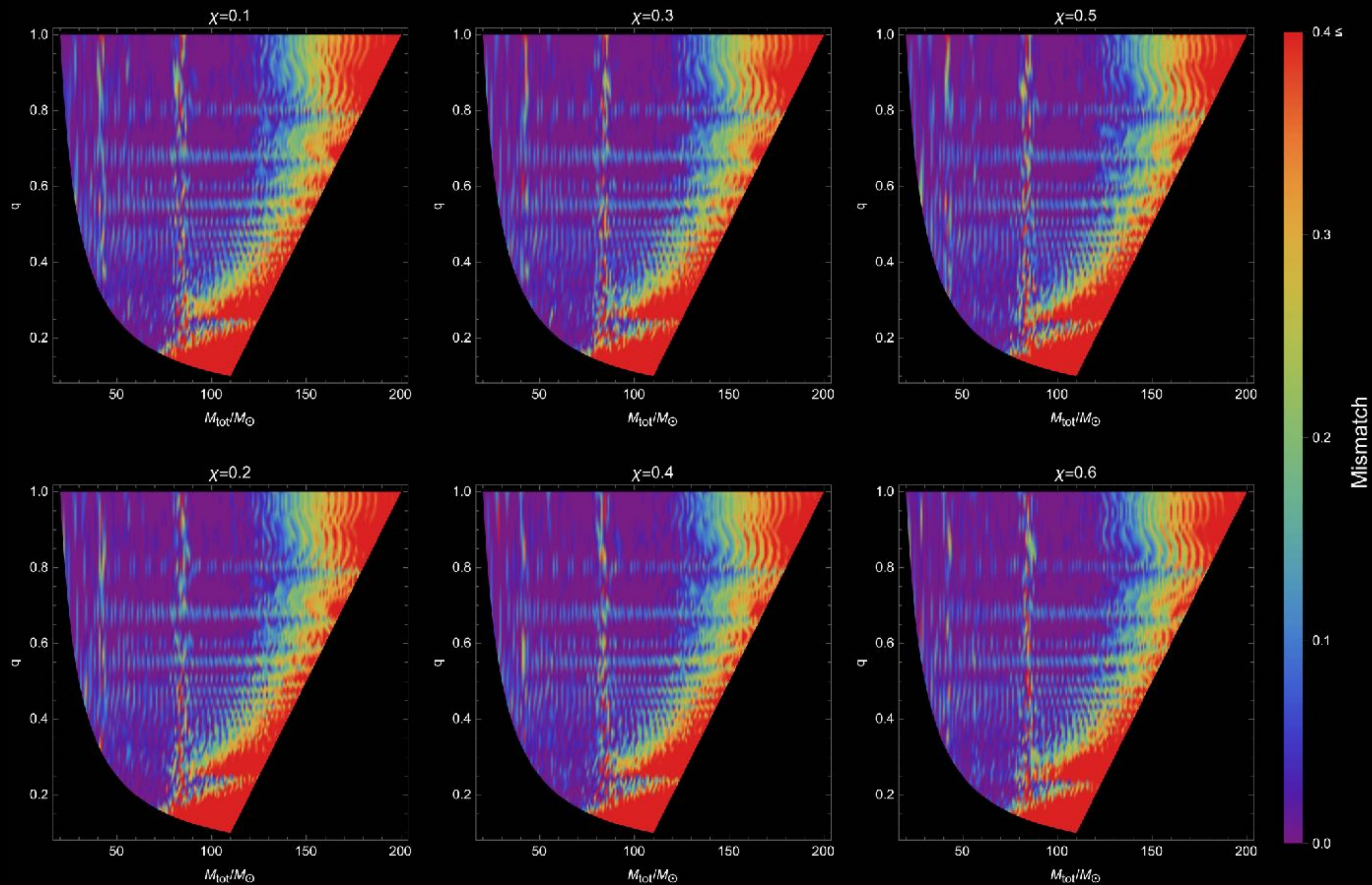
- The **kuibit** Python library was used.

Mismatch map for the not-spinning configurations

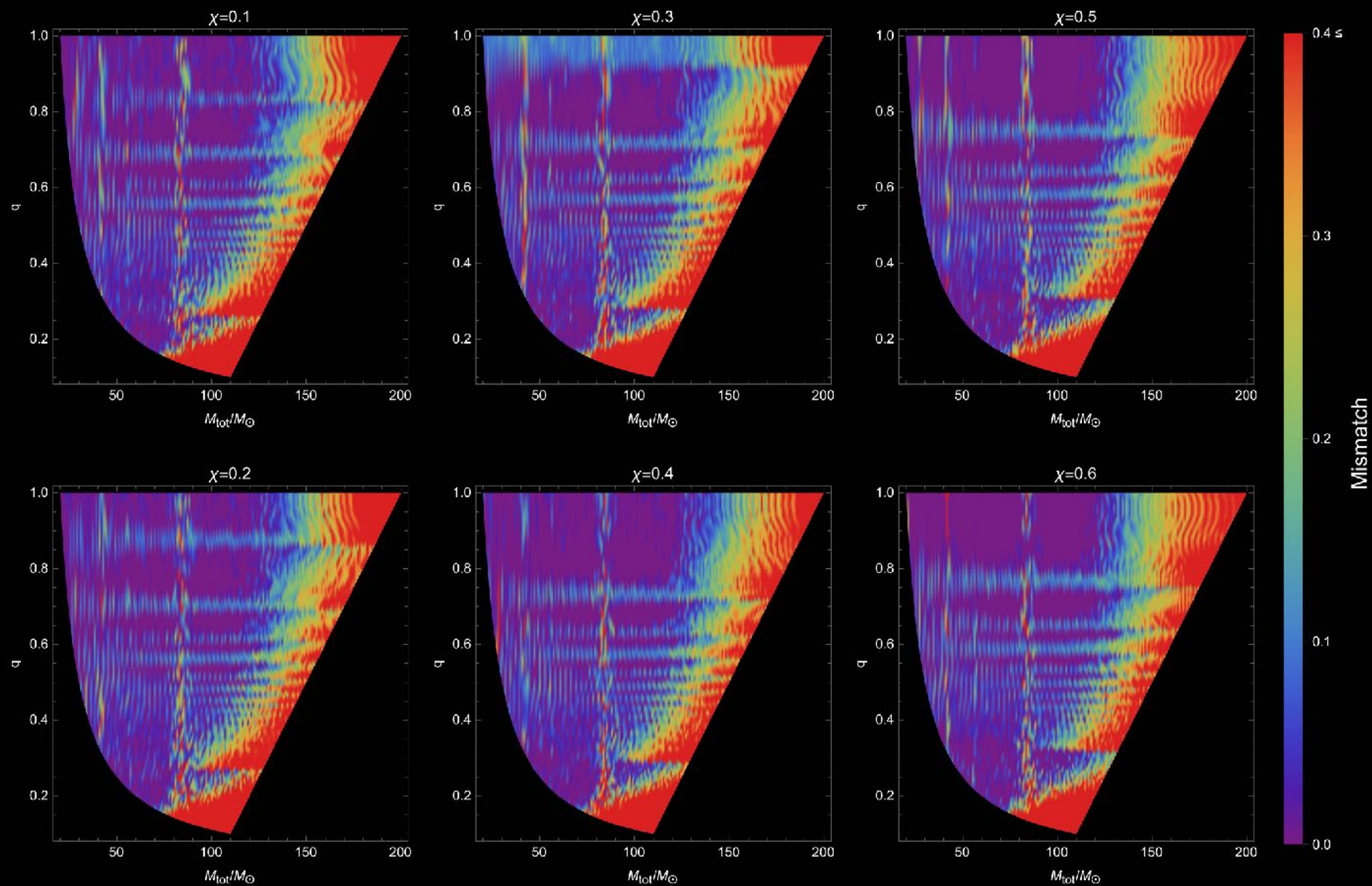
$m_1 = 44.5455 M_\odot$, $m_2 = 39.0909 M_\odot$ and $q = 0.8775499209$



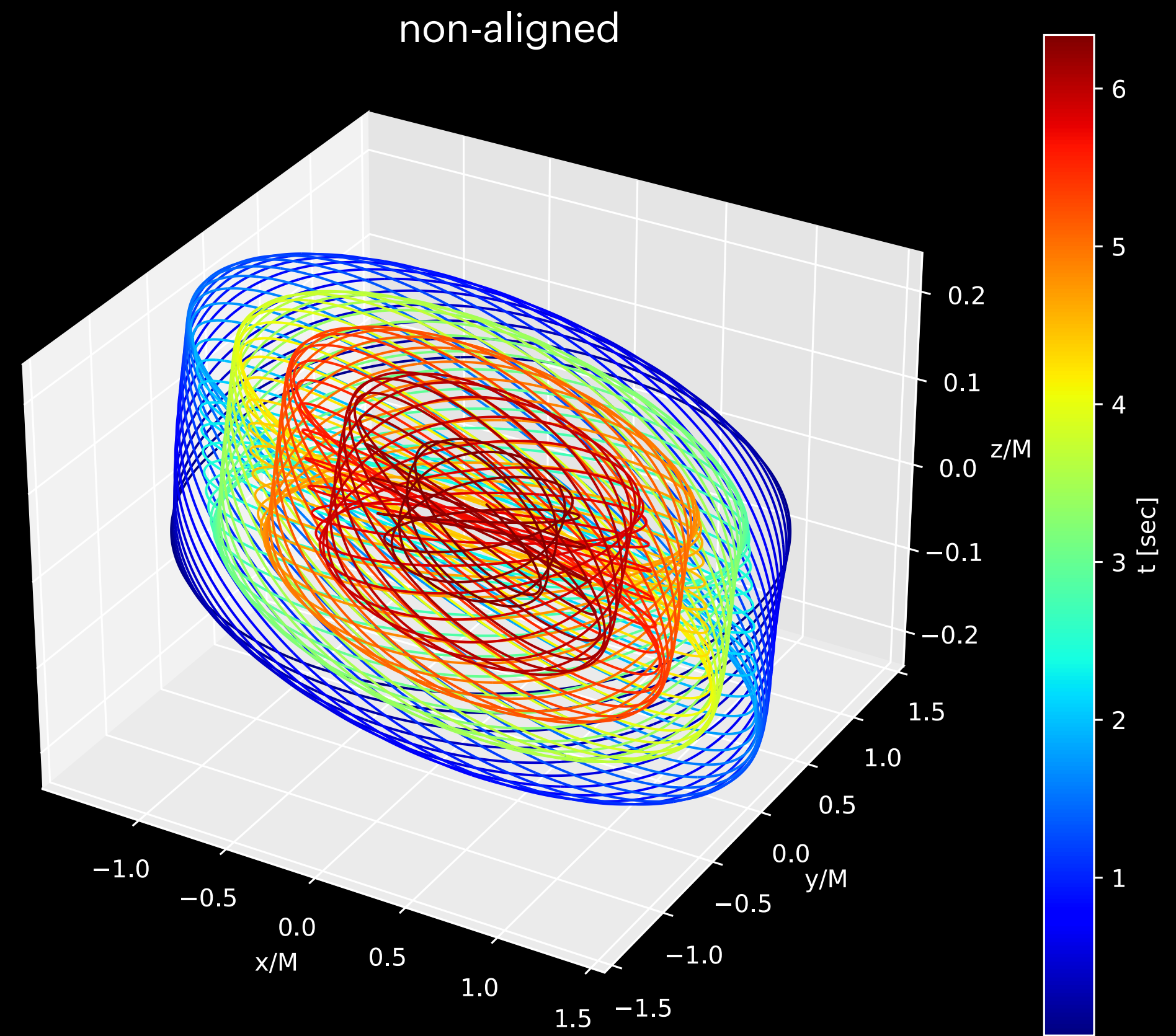
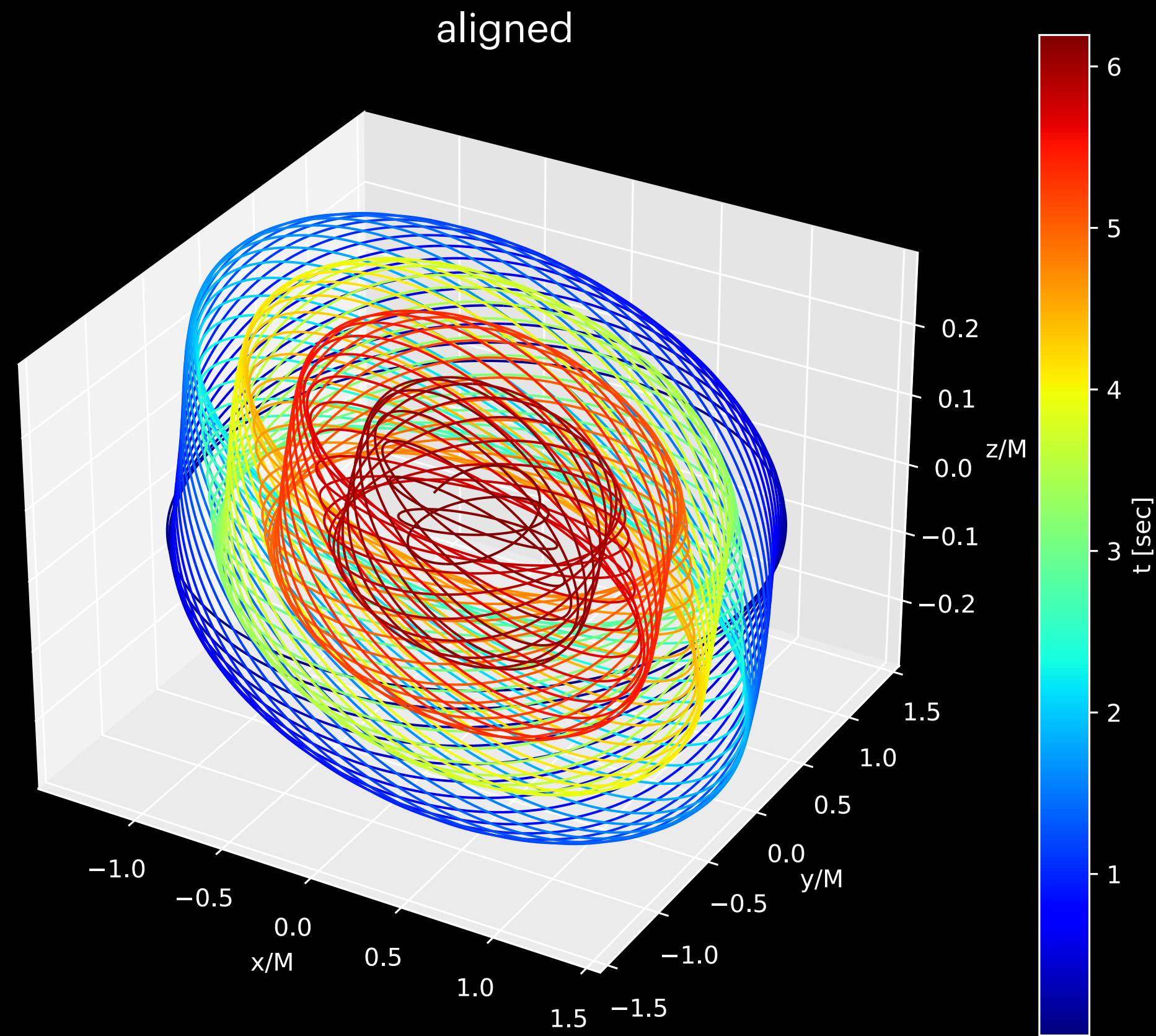
Mismatch map for the spin-aligned configurations



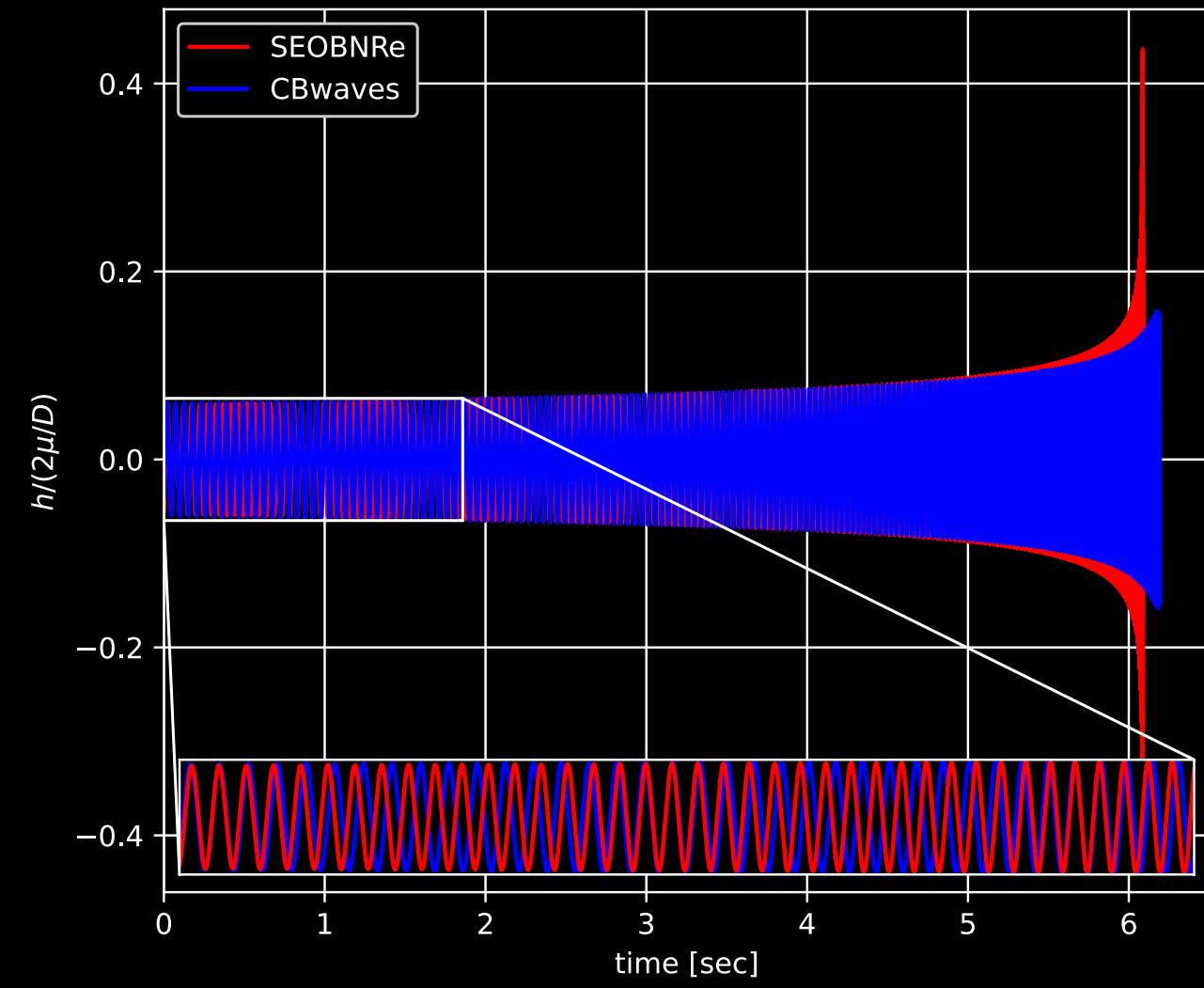
Mismatch map for the non-aligned spin configurations



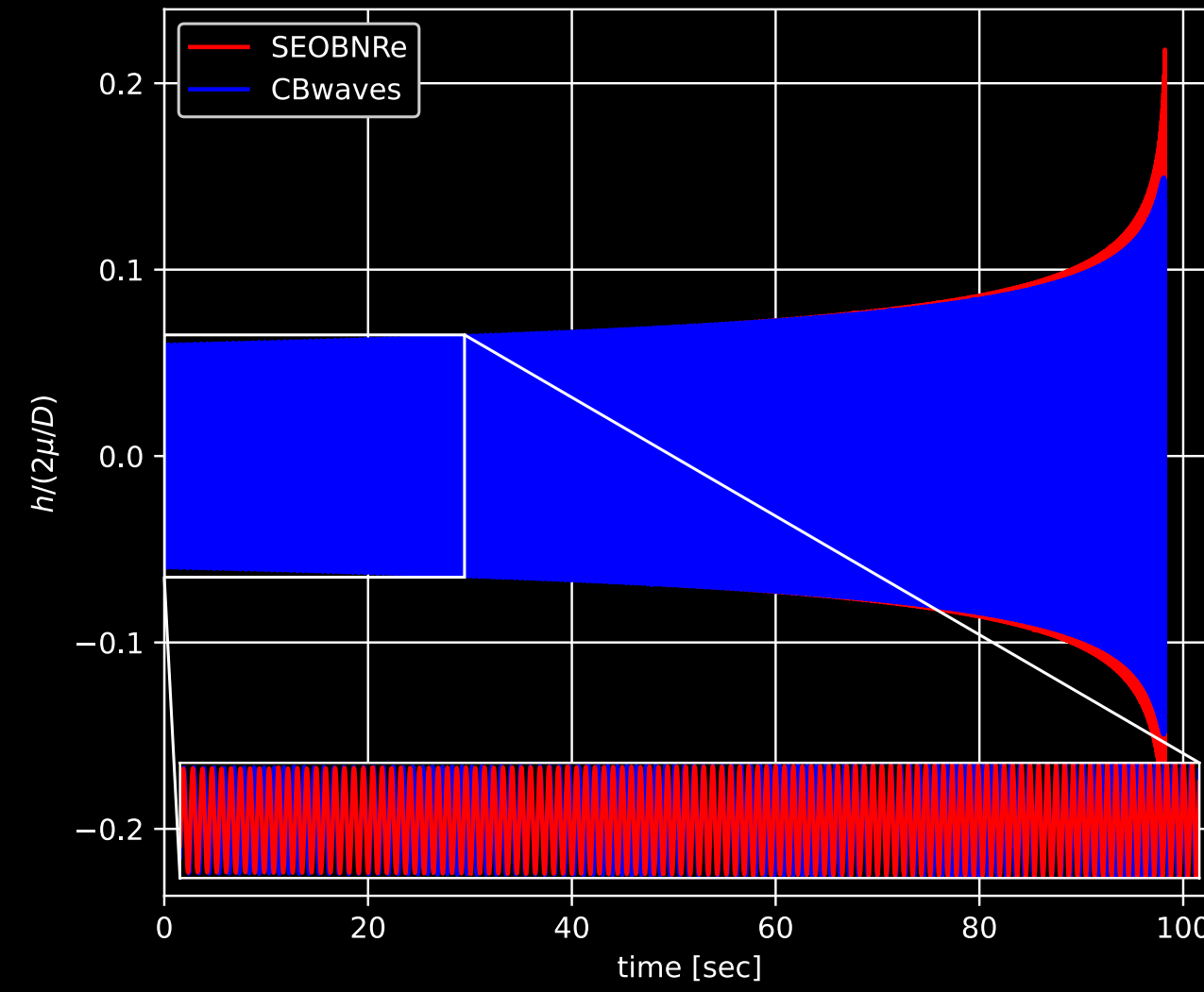
Orbital evolution of the binary at $\chi_1 = 0.6$, $m_1 = m_2 = 10 M_\odot$



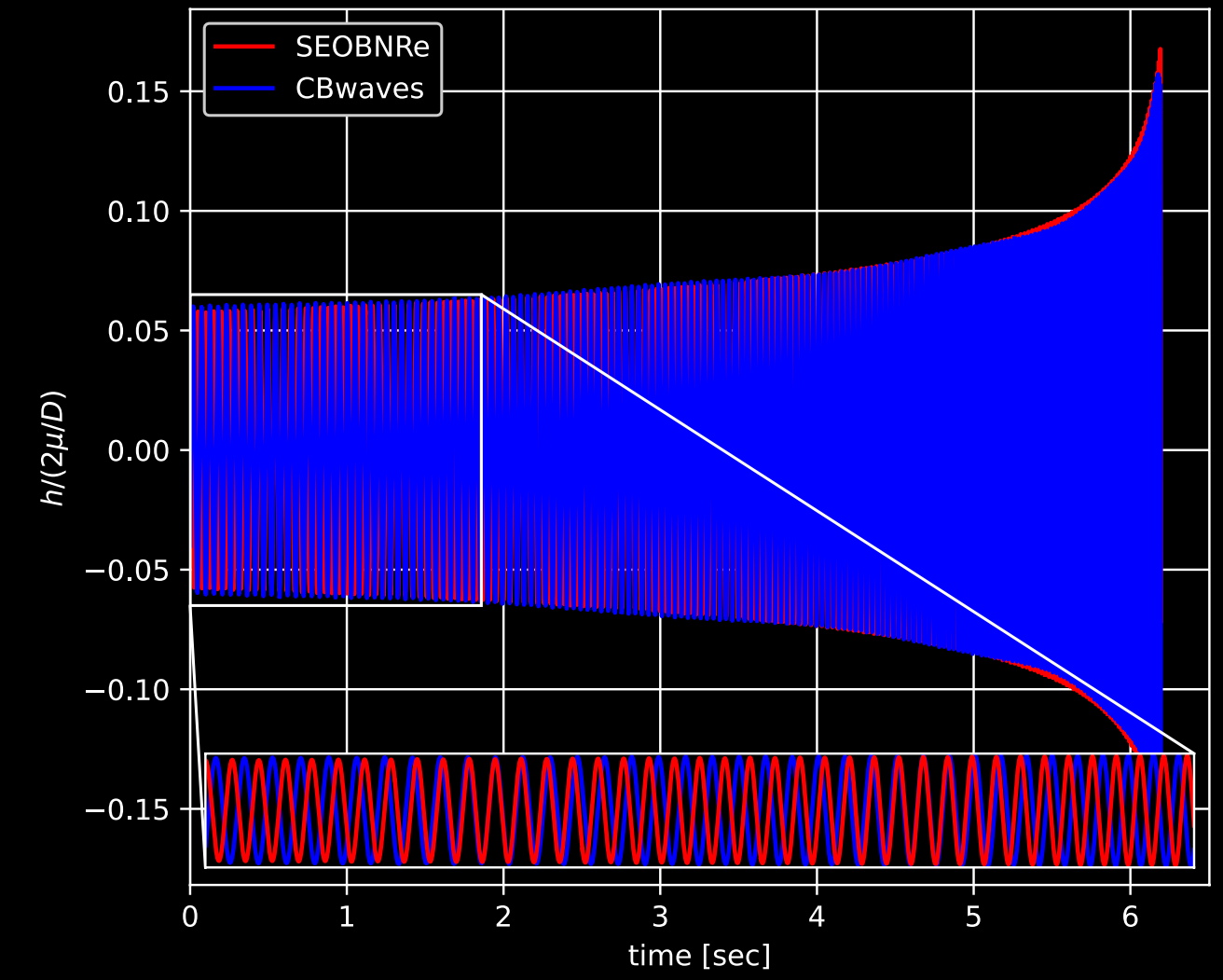
not-spinning, $m_2 = 10 M_\odot$, $m_2 = 10 M_\odot$



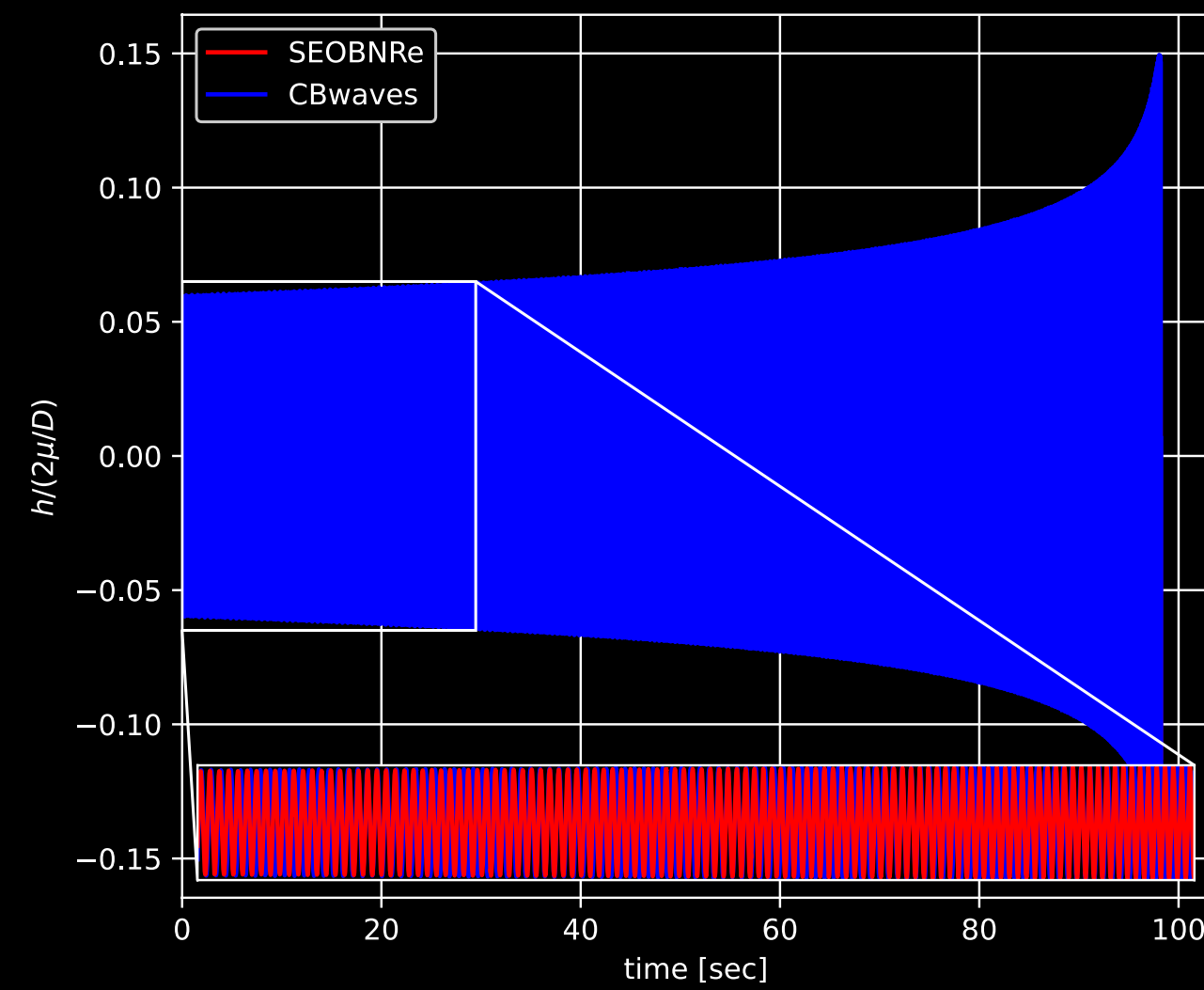
not-spinning, $m_2 = 100 M_\odot$, $m_2 = 10 M_\odot$



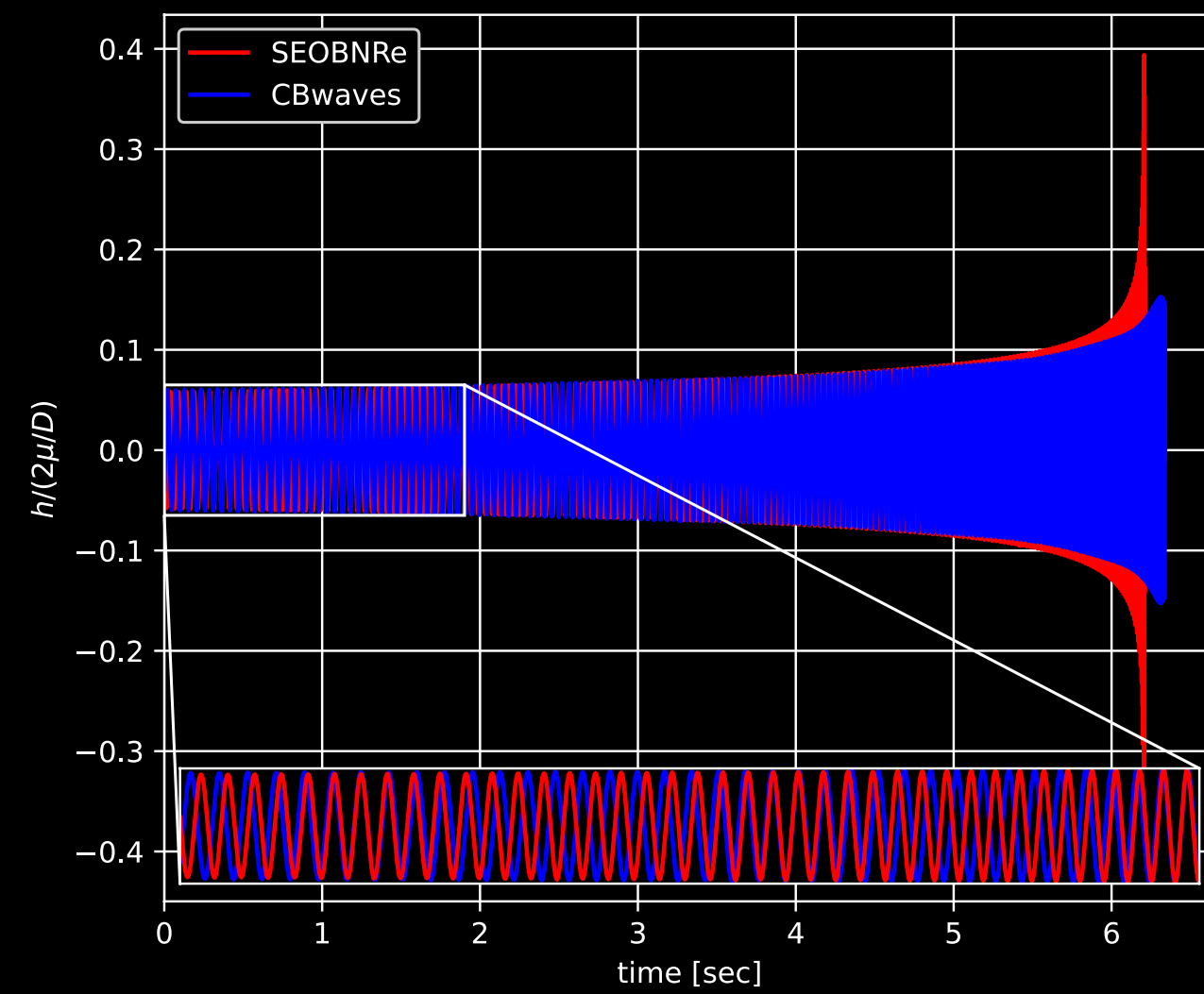
$\chi_1 = 0.6$, aligned, $m_2 = 10 M_\odot$, $m_2 = 10 M_\odot$



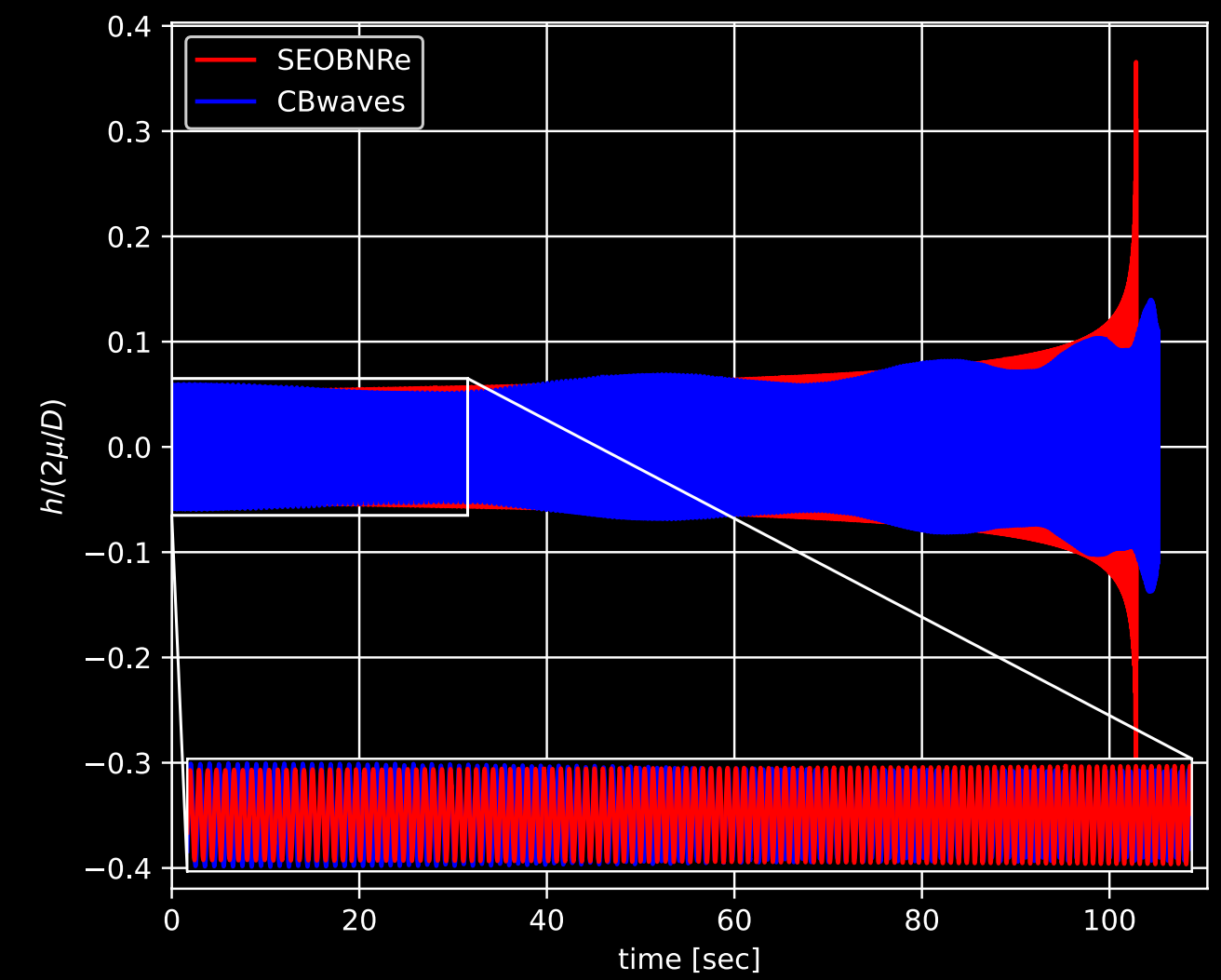
$\chi_1 = 0.6$, aligned, $m_2 = 100 M_\odot$, $m_2 = 10 M_\odot$



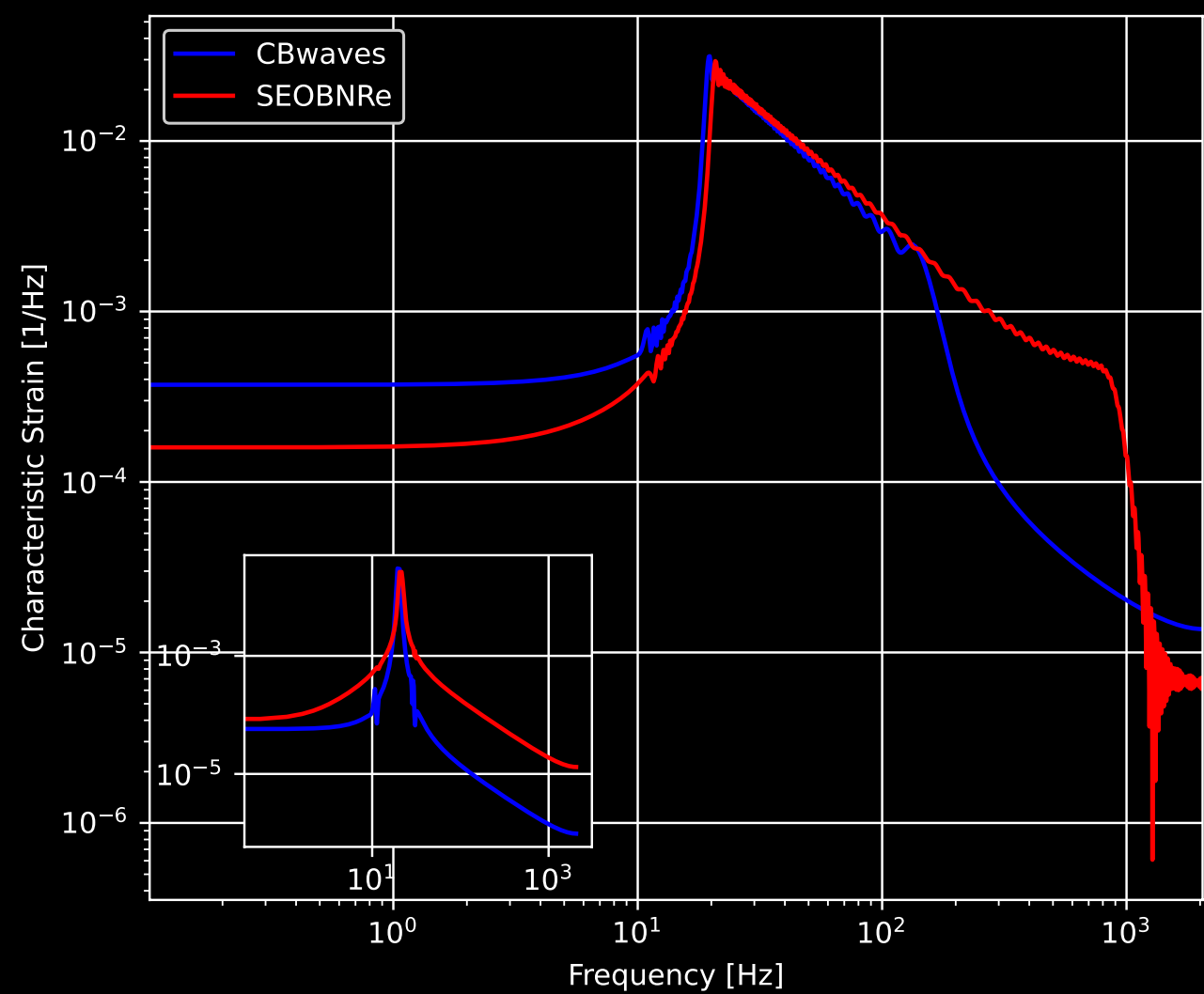
$\chi_1 = 0.6$, aligned, $m_2 = 10 M_\odot$, $m_2 = 10 M_\odot$



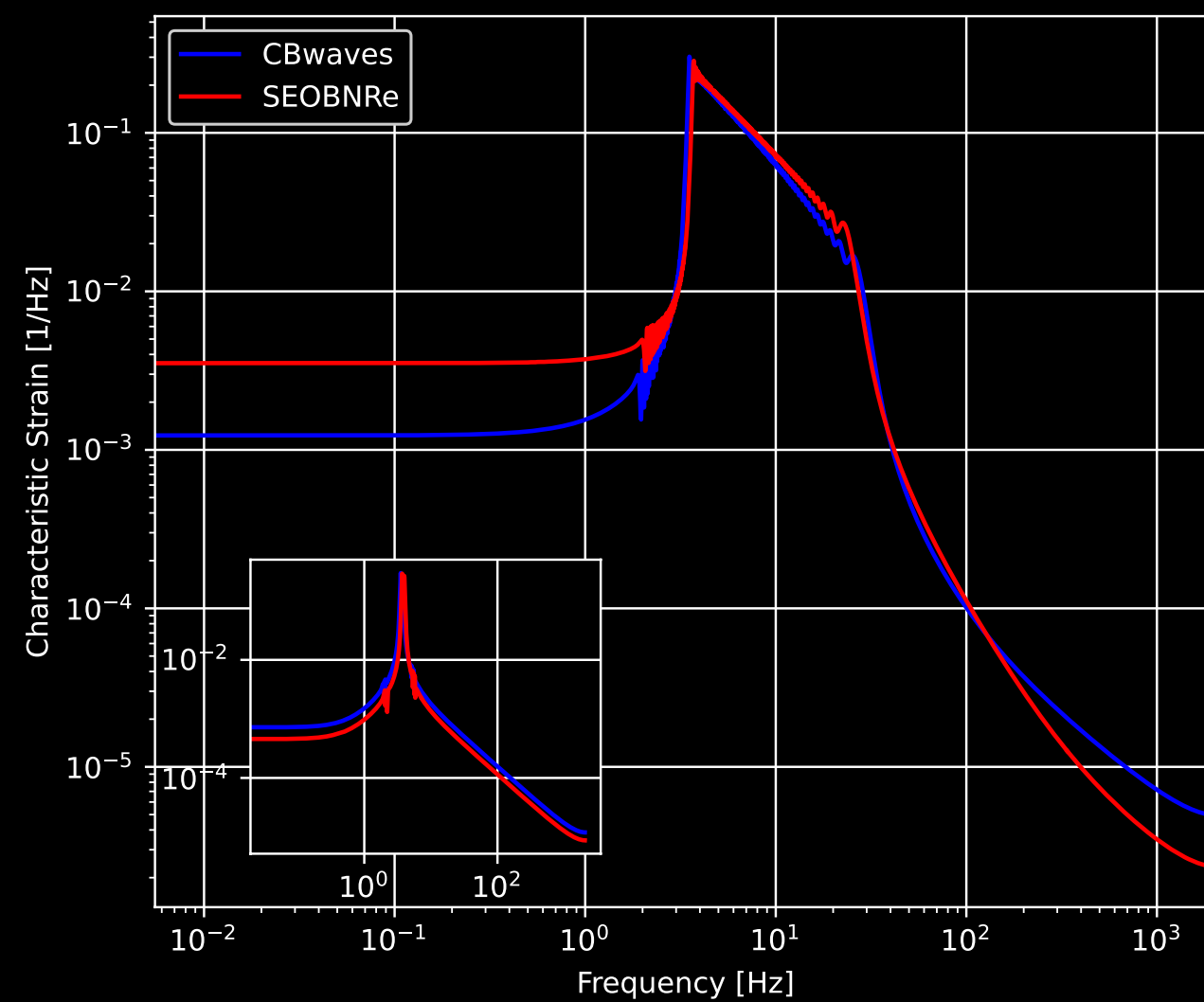
$\chi_1 = 0.6$, aligned, $m_2 = 100 M_\odot$, $m_2 = 10 M_\odot$



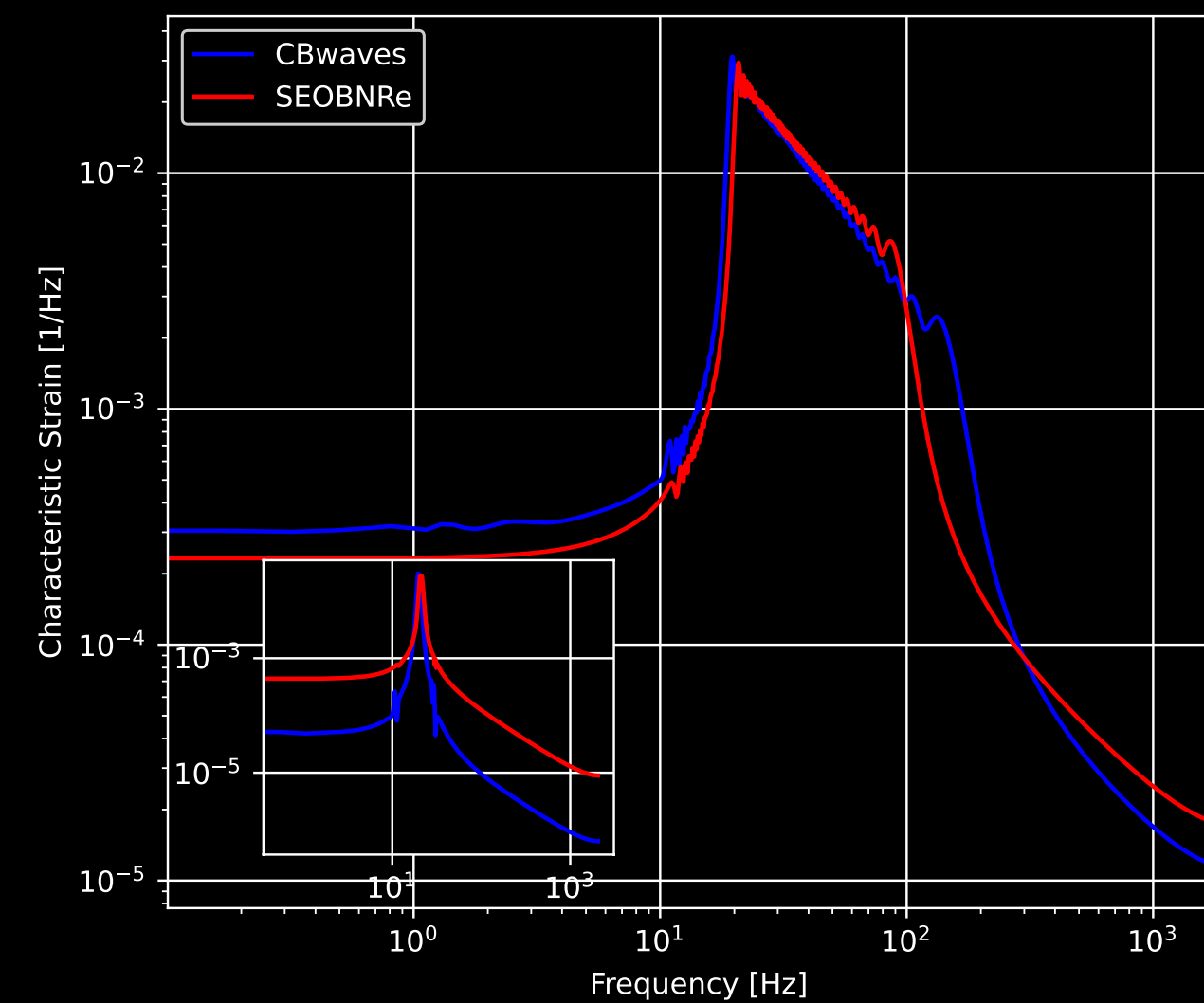
not-spinning, $m_2 = 10 M_\odot$, $m_2 = 10 M_\odot$



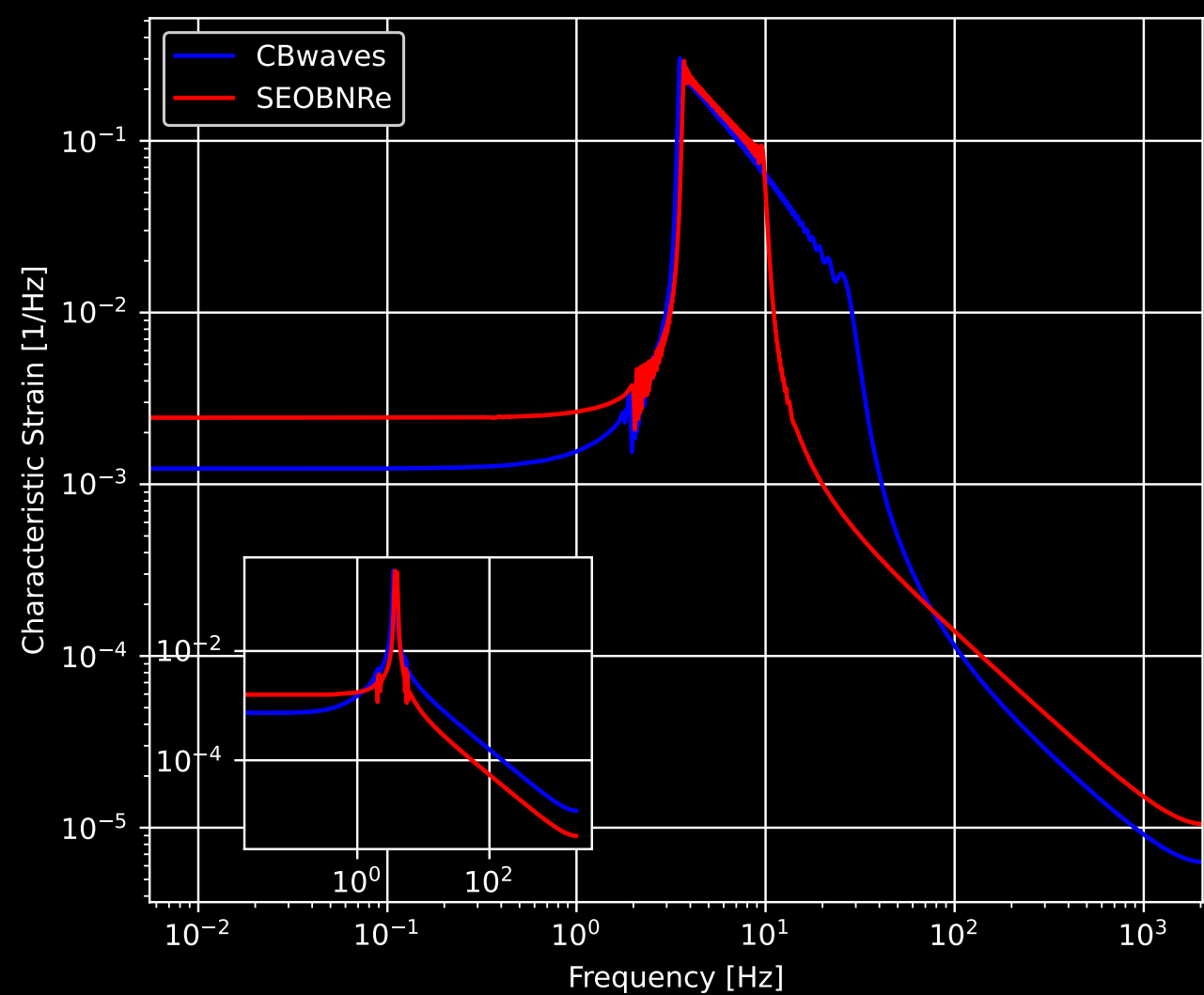
not-spinning, $m_2 = 100 M_\odot$, $m_2 = 10 M_\odot$



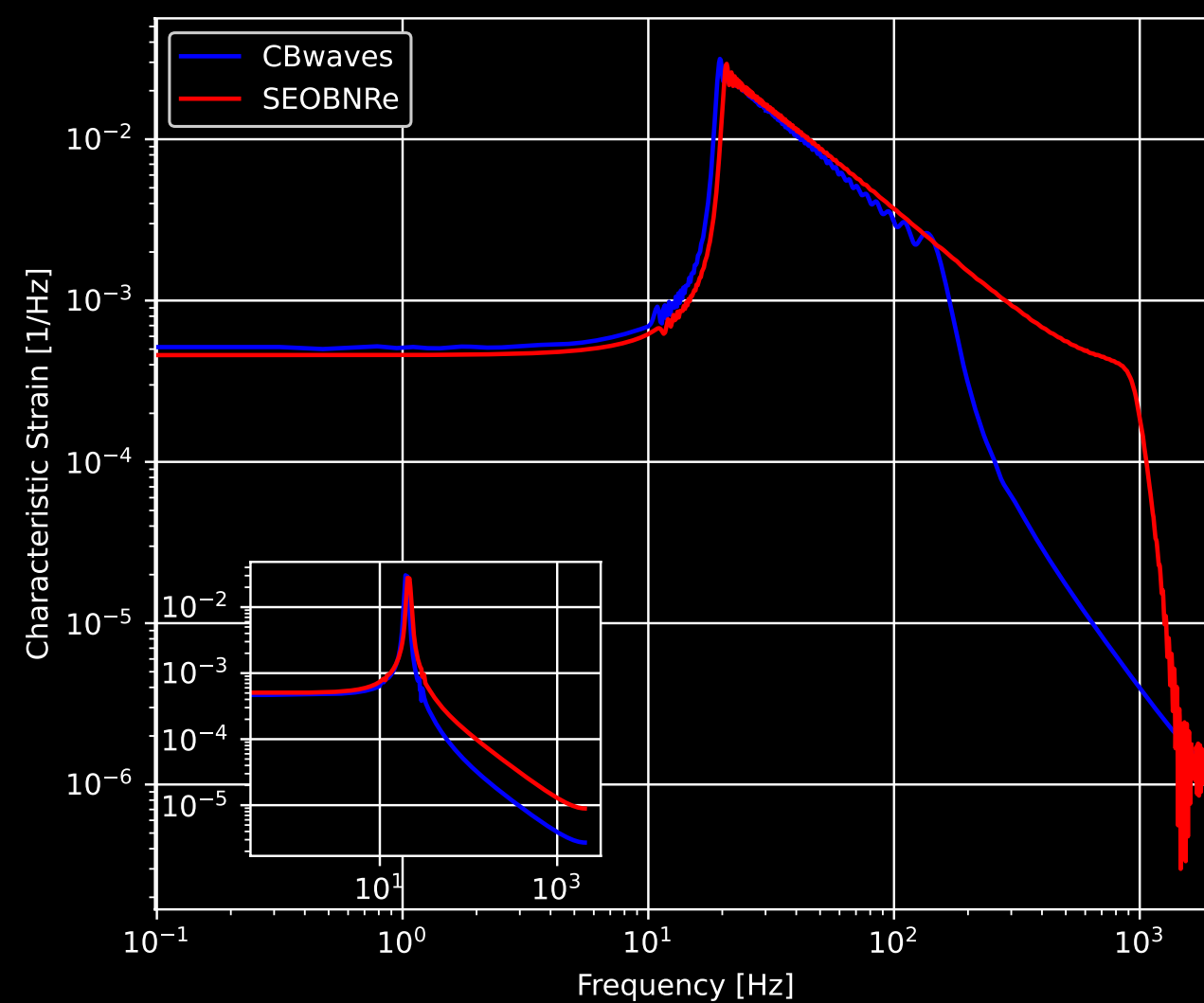
$\chi_1 = 0.6$, aligned, $m_2 = 10 M_\odot$, $m_2 = 10 M_\odot$



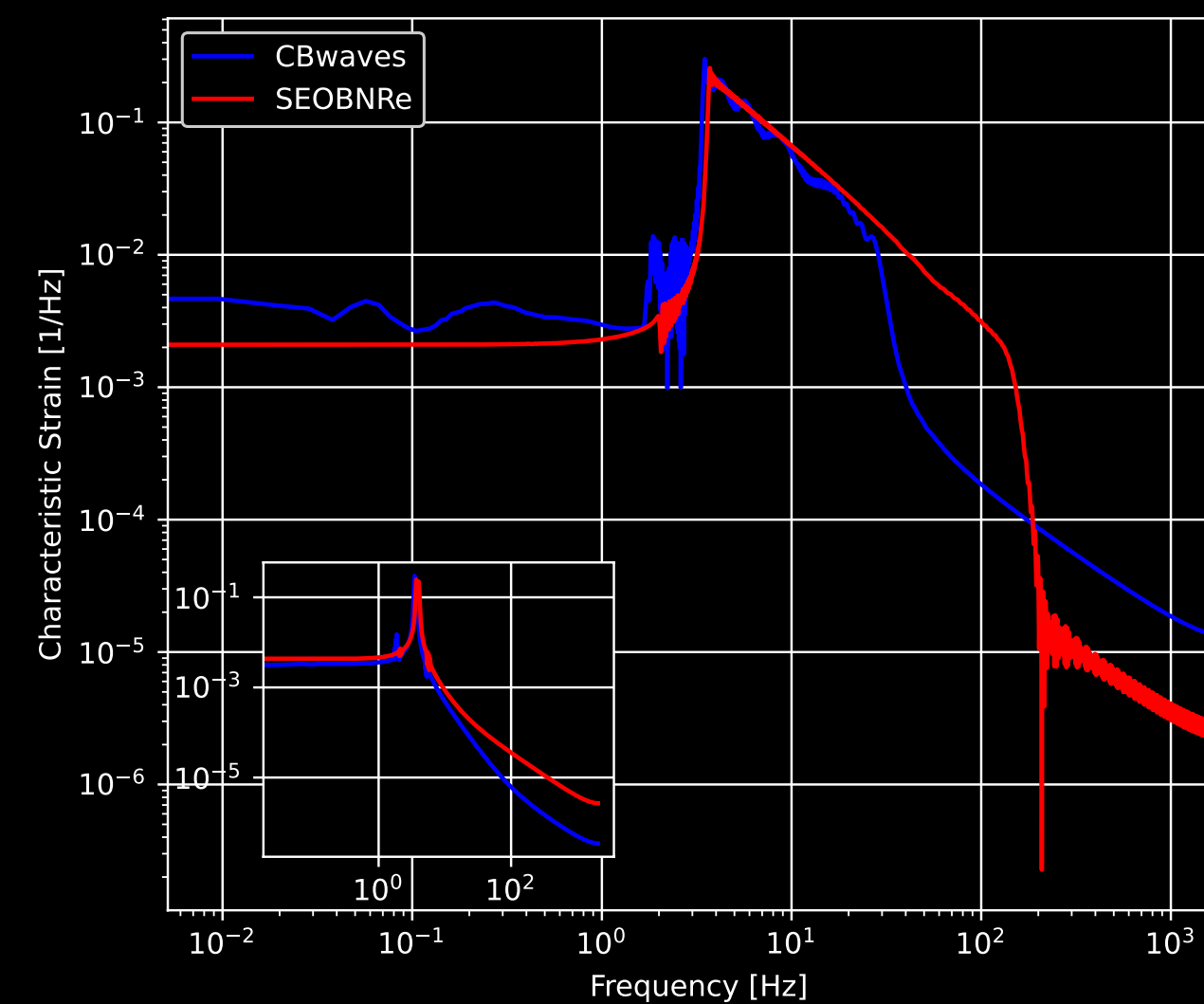
$\chi_1 = 0.6$, aligned, $m_2 = 100 M_\odot$, $m_2 = 10 M_\odot$



$\chi_1 = 0.6$, aligned, $m_2 = 10 M_\odot$, $m_2 = 10 M_\odot$



$\chi_1 = 0.6$, aligned, $m_2 = 100 M_\odot$, $m_2 = 10 M_\odot$



Discussion and Conclusion

- At $1 : 20$ mass-ratio, the separation computed by both codes is in close agreement
- For configurations with $q < 1/20$ the $6 M$ limit reached earlier by SEBNRE
- For configurations with $q > 1/20$ the $6 M$ limit reached earlier by CBWaves
- Made detailed contour maps for the mismatch (or unfaithfulness) of various spin configurations
- As the mass-ratio is closing $1 : 10$, the mismatch between the two models grows larger
- A similar behavior is exhibited toward larger total masses with spins, but irrespective of the spin alignment
- the spins did not retain the initial alignment set in CBwaves
- However, the effects of the spin are unnoticeable on the aligned waveforms

Acknowledgment

- ▶ Deeply grateful to **WSCLab.** for computational resources provided
- ▶ This work was made in collaboration with **Dániel Barta**
- ▶ Thanks for the insight and pieces of advice given by **László Á. Gergely**



The background features a stylized mountain range with several peaks of varying heights and widths. The mountains are rendered in a gradient of colors, from a light, almost white purple at the top to a dark, deep blue at the bottom. The overall effect is a soft, atmospheric landscape.

Thanks for your attention!