

Ab-initio theory of nuclear spin flip processes within NV center of diamond *via orbital degrees of freedom*

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experimental *theoretical model*

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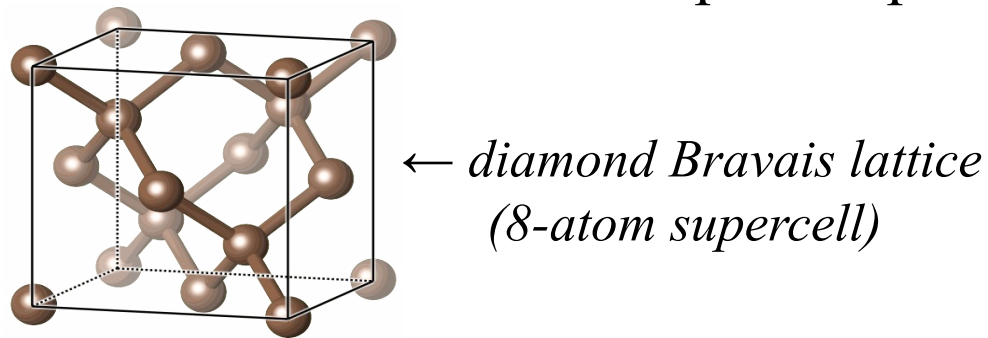
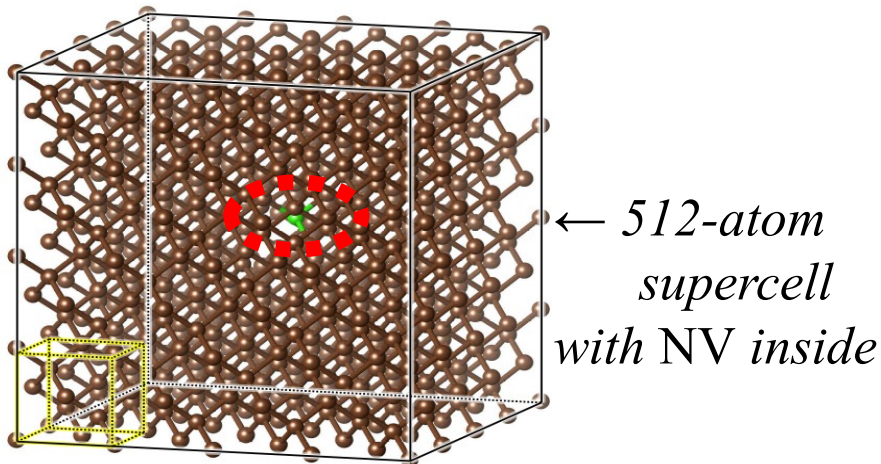
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DFT is conventionally used for: electronic structure

simulate defects in ~100-1000-atom supercells to get:

- 👉 **Formation energies**
- 👉 **Optical excitations** (*~0.1 eV precision – HSE06 hybrid functional*)
- 👉 **Electron-phonon coupling** (*vibronic sideband for optical centers*)

DFT: *density functional theory* – simulate the electronic structure on HPC supercomputers



DFT is conventionally used for: electronic structure

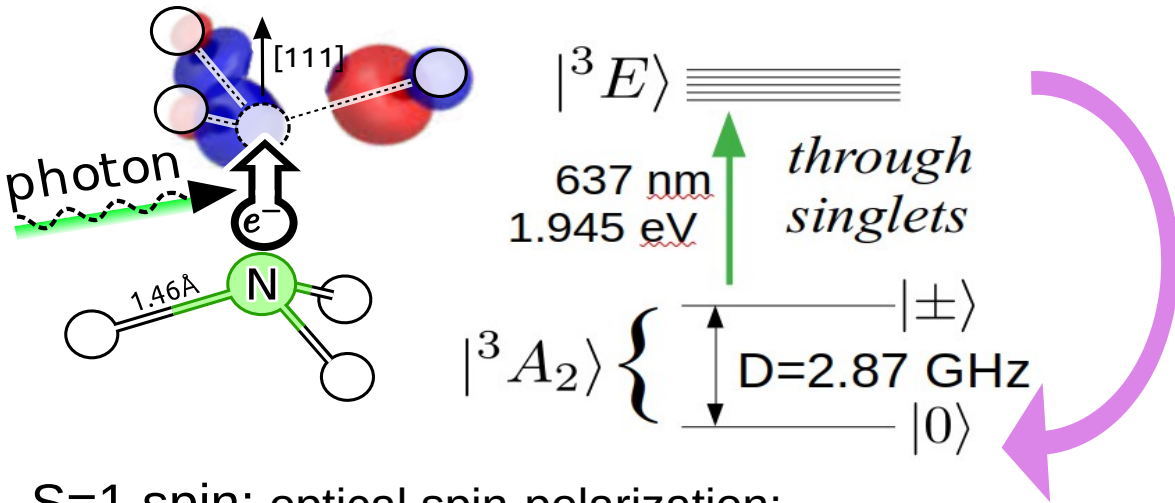
simulate defects in ~100-1000-atom supercells to get: “hidden” properties of qubits

- 👉 **Formation energies**
- 👉 **Optical excitations** (~ 0.1 eV precision – HSE06 hybrid functional)
- 👉 **Electron-phonon coupling** (*vibronic sideband for optical centers*)
- 👉 **Spin-phonon relaxation**
- 👉 **Spin-orbit matrix elements:** λLS (λ up to $\sim 20\%$ precision)
- 👉 **Spin-spin interaction:** **ZFS** (*Zero field splitting*) *magnetic dipole-dipole: SDS*
- 👉 **Hyperfine interaction:** *electronic + nuclear spin dipole-dipole: SAI*
- 👉 **Nuclear quadrupolar interaction: IQI**

aim: determine unconventional (spin) parameters inaccessible by experiments

This talk: **Predict** motion of ^{14}N nuclear spin during optical cycles

Introduction: *qubit* $|0\rangle$ initialization for NV – *electronic spin*



S=1 spin: optical spin-polarization:
ms=0 is preferentially populated
over **ms=±1**
upon exposure to green light

G. Thiering, A. Gali, Phys. Rev. B **96**, 081115(R) (2017)

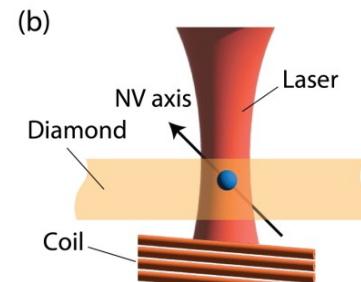
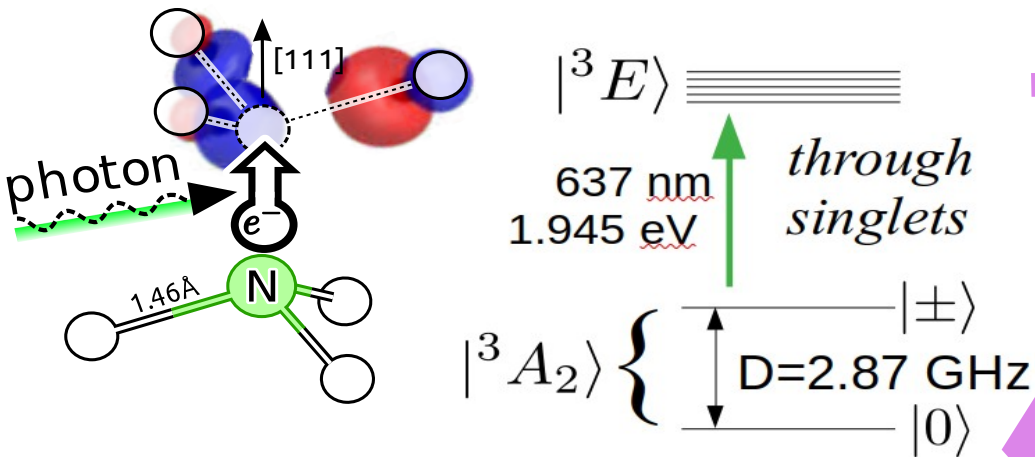
G. Thiering, A. Gali, Phys. Rev. B **98**, 085207 (2018)

M. L. Goldman, ... Phys. Rev. Lett. **114**, 145502 (2015)

... and many other studies

Introduction: *qubit* $|0\rangle$ initialization for $NV - {}^{14}\text{N}$ nuclear spin

This talk: **Predict** relaxation of ${}^{14}\text{N}$ nuclear spin during optical cycles

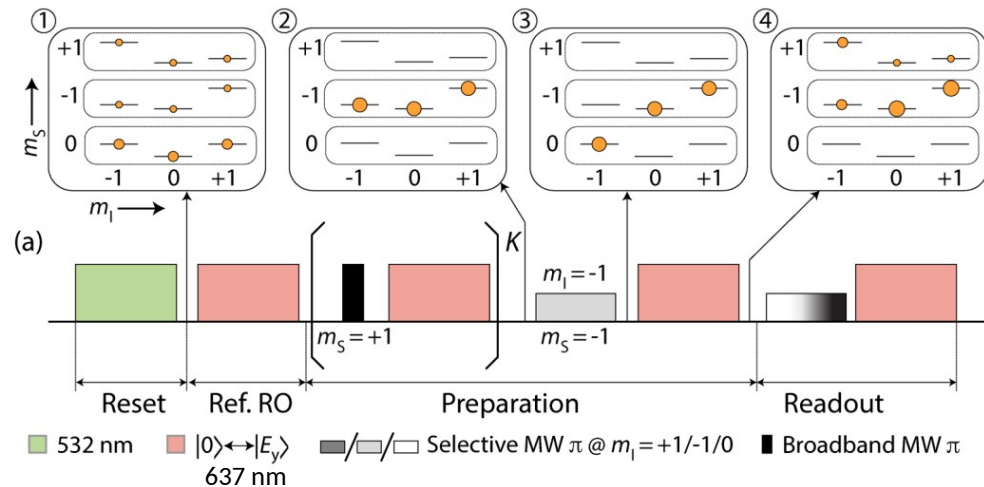


${}^{14}\text{N}$ nuclear spin can be **initialized** too: hyperpolarization

quantum memory
long T_1 time: seconds, hours
 T_2 time \sim ms

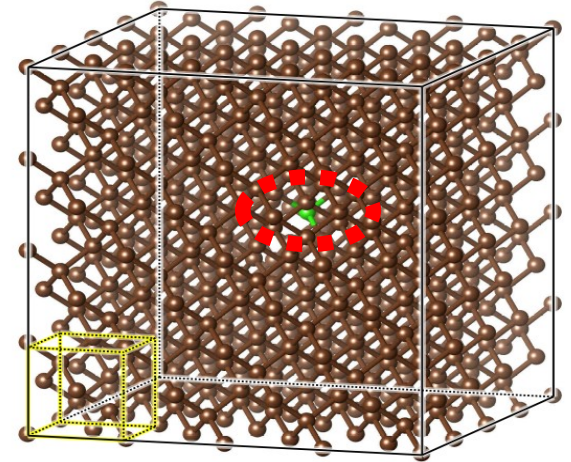
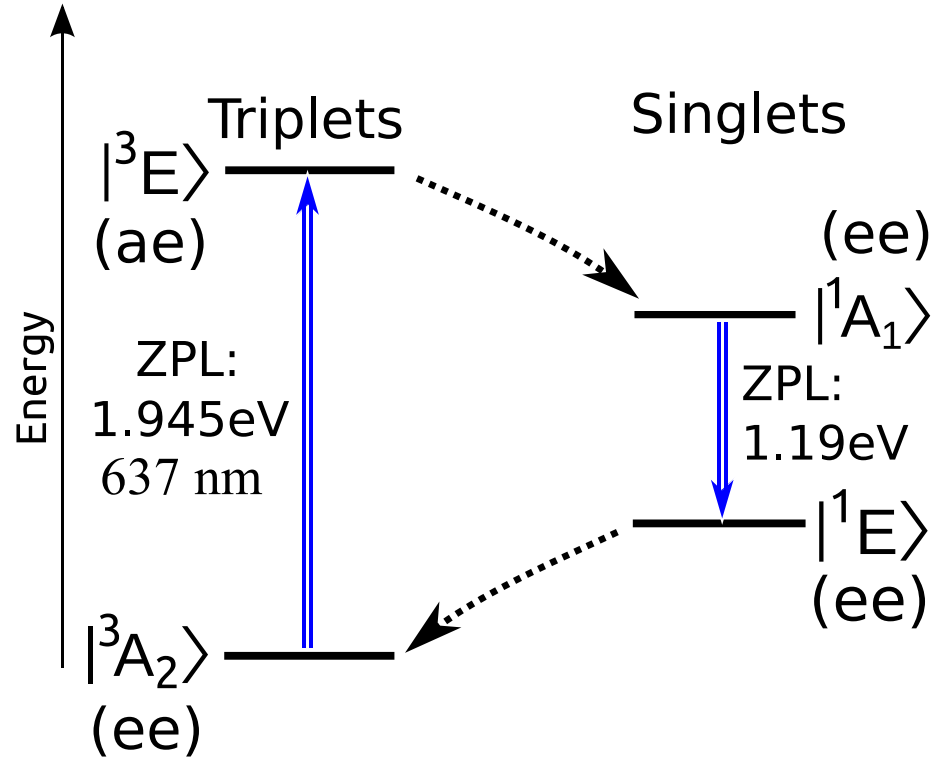
Preparation and readout of ${}^{14}\text{N}$ spins:

S=1 spin: optical spin-polarization:
 $m_s=0$ is preferentially populated
over $m_s=\pm 1$
upon exposure to green light



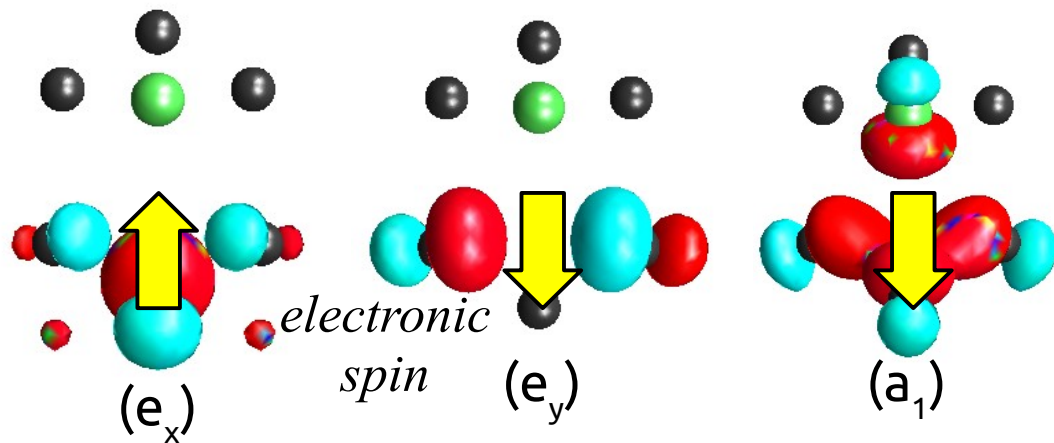
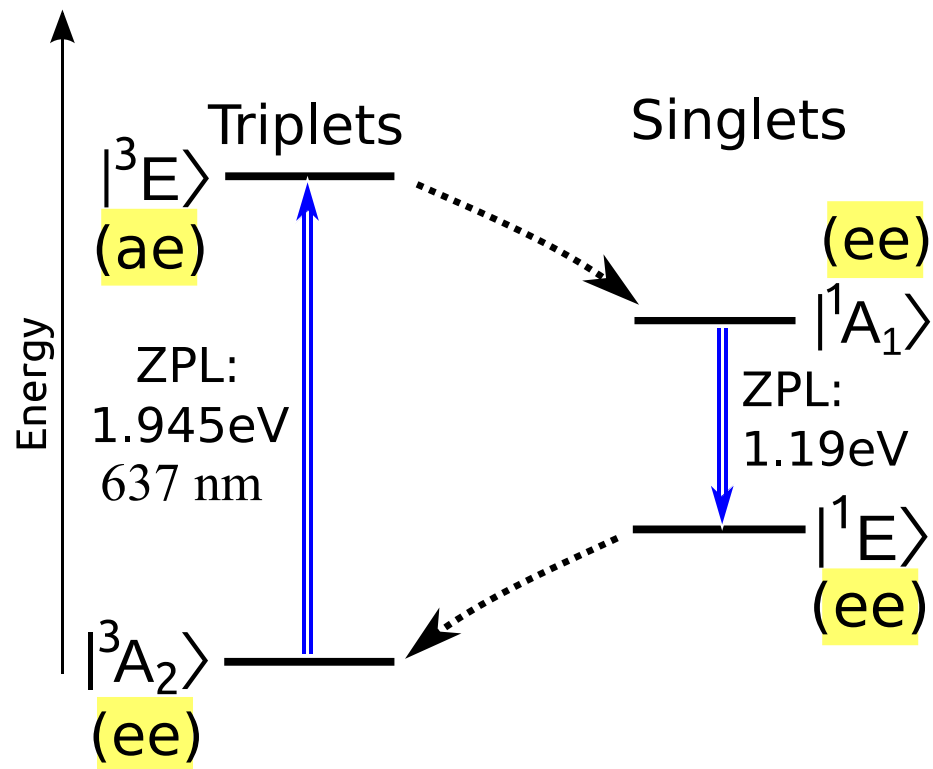
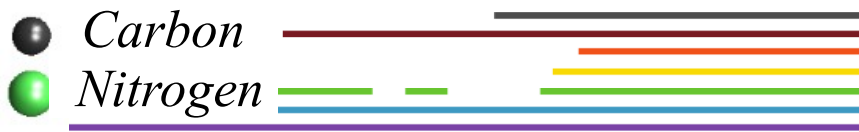
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- ... and many other studies

Electronic structure of NV



← 512-atom diamond supercell with NV inside

Electronic structure of NV



Electronic structure in short:

pick two orbitals above

for

two particle wavefunctions

also arrange their electronic spin

$|^3E\rangle$:

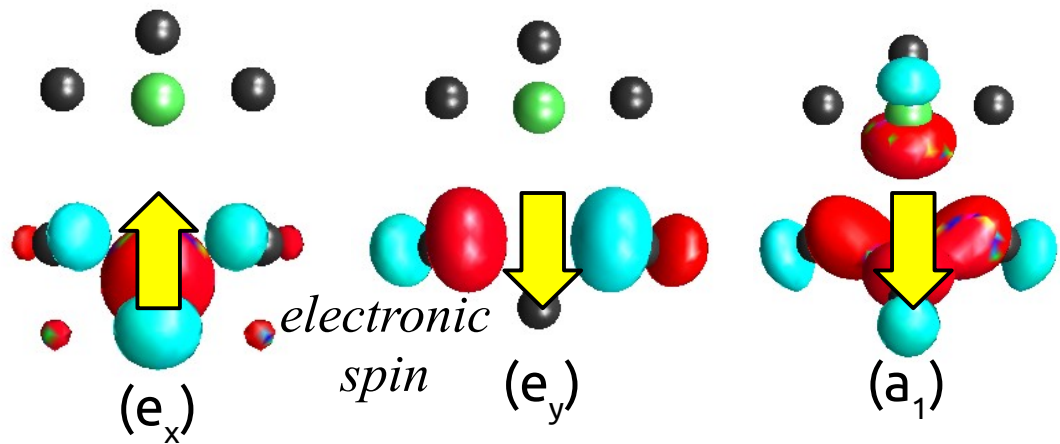
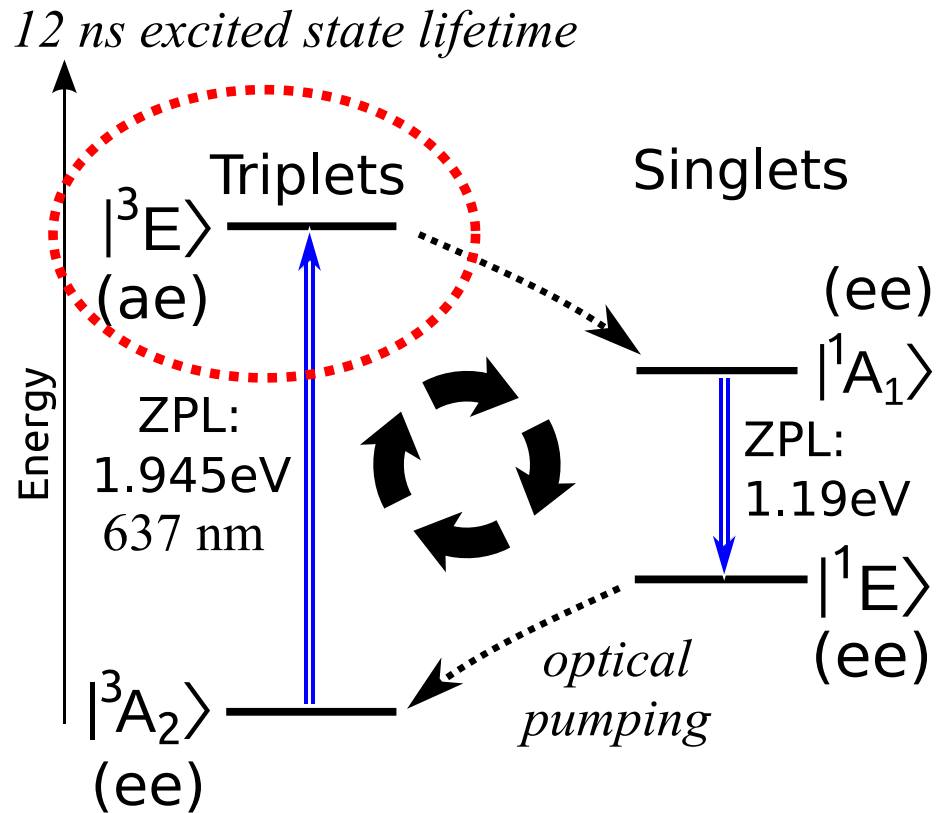
m_L - orbital quantum number: +1, -1 or e_x, e_y

m_S - electronic spin: +1, 0, -1

$|^3A_2\rangle$:

m_S - electronic spin: +1, 0, -1

Electronic structure of NV



Electronic structure in short:
 pick two orbitals above
 for
 two particle wavefunctions
 also arrange their electronic spin

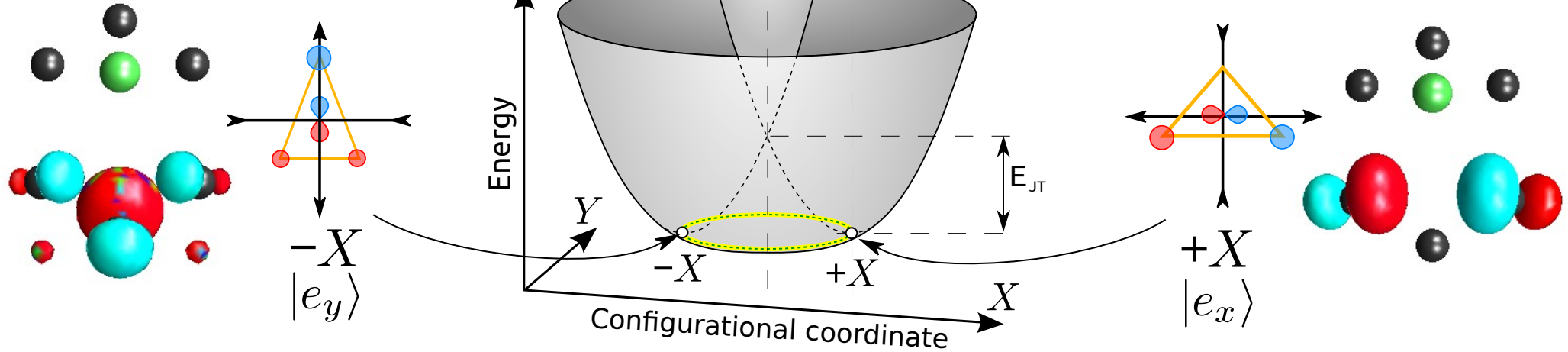
$|^3A_2\rangle$:
 m_S - electronic spin: +1,0,-1

orbital degeneracy:

$|^3E\rangle$:
 m_L - orbital quantum number: +1, -1 or e_x, e_y
 m_S - electronic spin: +1,0,-1

Jahn teller effect within $|^3E\rangle$

(ae_x) or (ae_y) orbital degeneracy:
two-particle system



orbitals induce strain on geometry

“E” 2x *degenerate* vibration mode
interacts with
“e” 2x *degenerate* orbital

ladder operators
for vibration modes:

$$\hat{X}, \hat{Y} = \frac{1}{\sqrt{2}} (a_{X,Y}^\dagger + a_{X,Y})$$

$$\hat{H}_{DJT} = F \left[\begin{pmatrix} 1 & \\ & -1 \end{pmatrix} \hat{X} + \begin{pmatrix} & 1 \\ 1 & \end{pmatrix} \hat{Y} \right]$$

Structure of ZFS tensor within $|3E\rangle$

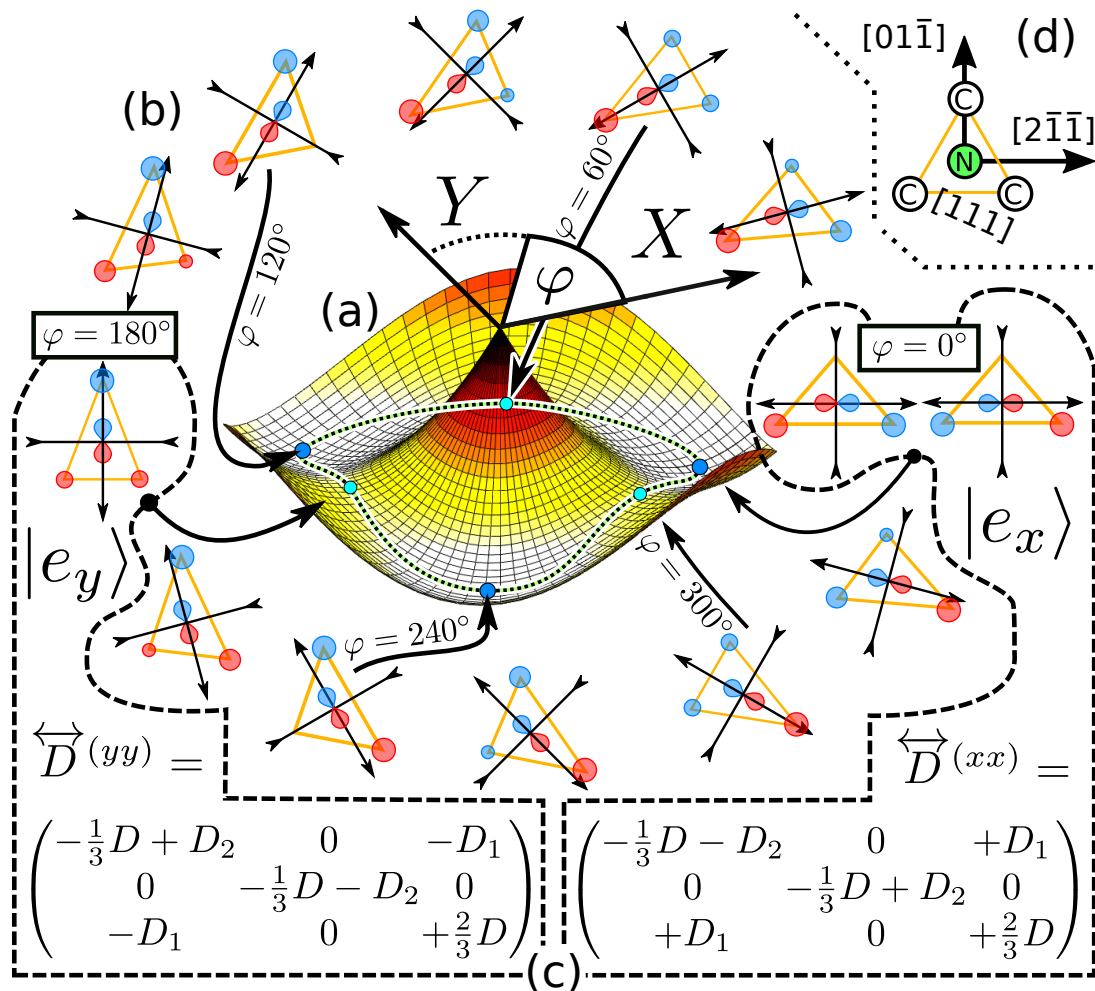
ZFS for the ground state:

$$\overleftrightarrow{D} = \begin{pmatrix} -\frac{1}{3}D & 0 & 0 \\ 0 & -\frac{1}{3}D & 0 \\ 0 & 0 & +\frac{2}{3}D \end{pmatrix}$$

$$\hat{H} = \overleftrightarrow{S} \overleftrightarrow{D} \overleftrightarrow{S}$$

☞ JT distortion induces distorted $|e_x\rangle, |e_y\rangle$ orbitals

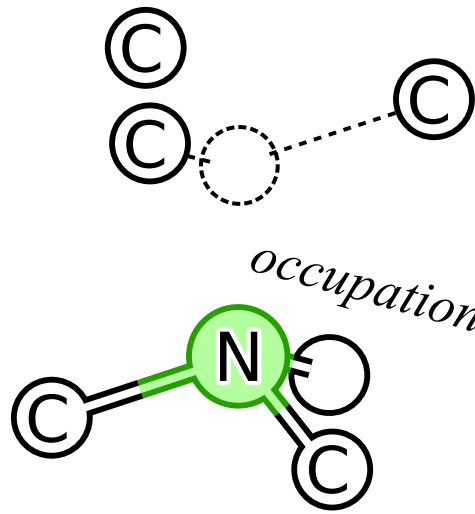
☞ JT motion induces large changes in ZFS!
 or in hyperfine: \overleftrightarrow{A}
 or in quadrupole: \overleftrightarrow{Q} tensors



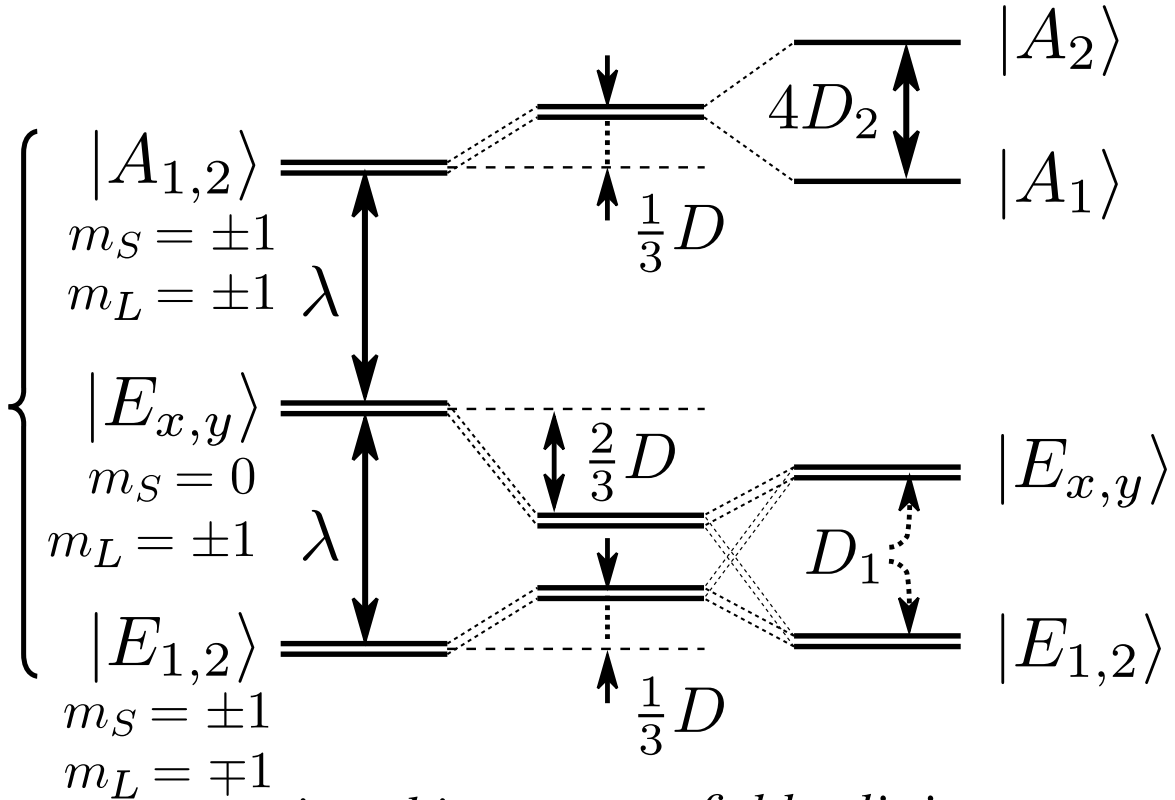
The $|^3E\rangle$ excited triplet (*fine structure*)

m_L - orbital quantum number $|e_{\pm}\rangle = (|e_x\rangle \pm |e_y\rangle)/\sqrt{2}$

m_S - electronic spin



occupation: $|^3E\rangle$
(ae_{\pm})



spin-orbit

zero-field splitting

or in short: ZFS

group theoretical approach for NV's ZFS:

M W Doherty et al 2011 New J. Phys. 13 025019

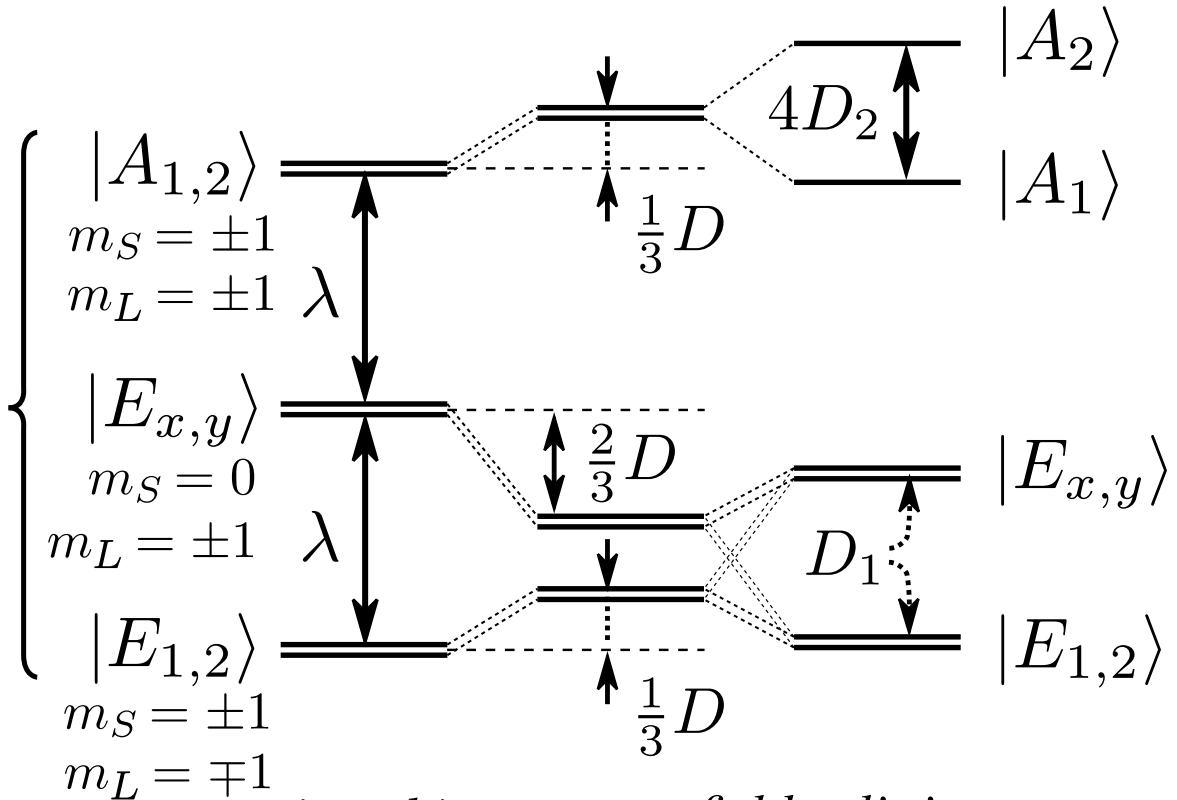
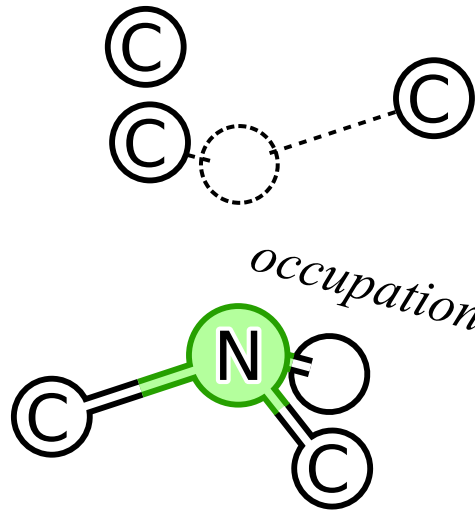
J R Maze et al 2011 New J. Phys. 13 025025

The $|^3E\rangle$ excited triplet (*fine structure*)

m_L - orbital quantum number $|e_{\pm}\rangle = (|e_x\rangle \pm |e_y\rangle)/\sqrt{2}$

m_S - electronic spin

m_I - ^{14}N nuclear spin



spin-orbit

zero-field splitting

or in short: ZFS

group theoretical approach for NV's ZFS:

M W Doherty et al 2011 New J. Phys. 13 025019

J R Maze et al 2011 New J. Phys. 13 025025

Spin Hamiltonian for the $|^3A_2\rangle$ ground state



$$\hat{H} = D^{(g)} \left(\hat{S}_z^2 - \frac{1}{3} S(S+1) \right) + Q^{(g)} \left(\hat{I}_z^2 - \frac{1}{3} I(I+1) \right) + A_{\parallel}^{(g)} \hat{S}_z \hat{I}_z + \frac{1}{2} A_{\perp}^{(g)} \left(\hat{S}_+ \hat{I}_- + \hat{S}_- \hat{I}_+ \right)$$

spin-spin ZFS
nuclear quadrupole
hyperfine interaction

$$+ \sum_{ij} A_{ij}^{C_1} \hat{S}_i \hat{I}_j^{C_1} + \dots$$

$\hat{S}_z, \hat{S}_{\pm}, |m_S = \pm 1, 0\rangle$ electronic S=1 spin of NV(-)

^{13}C hyperfine
 $A_C \sim \text{MHz}$

$\hat{I}_z, \hat{I}_{\pm}, |m_I = \pm 1, 0\rangle$ nuclear I=1 spin of ^{14}N

not discussed now

Spin Hamiltonian for the $|^3E\rangle$ excited level



spin-orbit

$$\hat{H}_0 = \lambda^{(e)} \hat{\sigma}_z \hat{S}_z +$$

$$D^{(e)} \left(\hat{S}_z^2 - \frac{1}{3} S(S+1) \right) + Q^{(e)} \left(\hat{I}_z^2 - \frac{1}{3} I(I+1) \right) + A_{\parallel}^{(e)} \hat{S}_z \hat{I}_z + \frac{1}{2} A_{\perp}^{(e)} \left(\hat{S}_+ \hat{I}_- + \hat{S}_- \hat{I}_+ \right)$$

spin-spin ZFS

nuclear quadrupole

hyperfine interaction

$$\hat{W} = D_1^{(e)} \left[(\hat{S}_z \hat{S}_+ + \hat{S}_+ \hat{S}_z) \hat{\sigma}_- + (\hat{S}_z \hat{S}_- + \hat{S}_- \hat{S}_z) \hat{\sigma}_+ \right] + D_2^{(e)} \left[\hat{S}_-^2 \hat{\sigma}_- + \hat{S}_+^2 \hat{\sigma}_+ \right] +$$

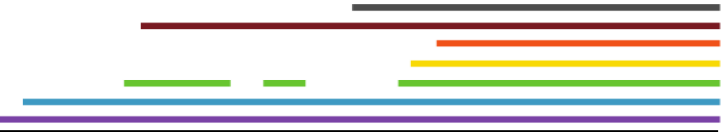
$$Q_1^{(e)} \left[(\hat{I}_z \hat{I}_+ + \hat{I}_+ \hat{I}_z) \hat{\sigma}_- + (\hat{I}_z \hat{I}_- + \hat{I}_- \hat{I}_z) \hat{\sigma}_+ \right] + Q_2^{(e)} \left[\hat{I}_-^2 \hat{\sigma}_- + \hat{I}_+^2 \hat{\sigma}_+ \right] +$$

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orbital flip: $\hat{\sigma}_{\pm} |e_{\pm}\rangle = |e_{\mp}\rangle$

nuclear+electronic spin-orbit

Spin Hamiltonian for the $|^3E\rangle$ excited level



spin-orbit

$$\hat{H}_0 = \lambda^{(e)} \hat{\sigma}_z \hat{S}_z +$$

$$D^{(e)} \left(\hat{S}_z^2 - \frac{1}{3} S(S+1) \right) + Q^{(e)} \left(\hat{I}_z^2 - \frac{1}{3} I(I+1) \right) + A_{\parallel}^{(e)} \hat{S}_z \hat{I}_z + \frac{1}{2} A_{\perp}^{(e)} \left(\hat{S}_+ \hat{I}_- + \hat{S}_- \hat{I}_+ \right)$$

spin-spin ZFS

nuclear quadrupole

hyperfine interaction

$$\Delta m_I = \pm 1$$

^{14}N nuclear spin flips

$$\Delta m_I = \pm 2$$

$$\hat{W} = D_1^{(e)} \left[(\hat{S}_z \hat{S}_+ + \hat{S}_+ \hat{S}_z) \hat{\sigma}_- + (\hat{S}_z \hat{S}_- + \hat{S}_- \hat{S}_z) \hat{\sigma}_+ \right] + D_2^{(e)} \left[\hat{S}_-^2 \hat{\sigma}_- + \hat{S}_+^2 \hat{\sigma}_+ \right] +$$

$$Q_1^{(e)} \left[(\hat{I}_z \hat{I}_+ + \hat{I}_+ \hat{I}_z) \hat{\sigma}_- + (\hat{I}_z \hat{I}_- + \hat{I}_- \hat{I}_z) \hat{\sigma}_+ \right] + Q_2^{(e)} \left[\hat{I}_-^2 \hat{\sigma}_- + \hat{I}_+^2 \hat{\sigma}_+ \right] +$$

$$A_1^{(e)} \left[(\hat{I}_z \hat{S}_+ + \hat{I}_+ \hat{S}_z) \hat{\sigma}_- + (\hat{I}_z \hat{S}_- + \hat{I}_- \hat{S}_z) \hat{\sigma}_+ \right] + A_2^{(e)} \left[(\hat{S}_- \hat{I}_- \hat{\sigma}_- + \hat{S}_+ \hat{I}_+ \hat{\sigma}_+) \right]$$

orbital flip: $\hat{\sigma}_{\pm} |e_{\pm}\rangle = |e_{\mp}\rangle$

nuclear+electronic spin-orbit

Quality of *ab initio* results

ab initio DFT can predict the spin Hamiltonian parameters $\pm 20\%$ precision



theoretical (experimental)

<i>spin-orbit</i>	$\lambda_1^{(e)} = 4.8^a (5.3^b) \text{ GHz}$	$D^{(g)} = 2.98^{\text{p.w.}} (2.87^e) \text{ GHz}$	<i>spin-spin</i>	
	$A_{\parallel}^{(g)} = -1.7^c (-2.14^d) \text{ MHz}$	$D^{(e)} = 1.67^{\text{p.w.}} (1.42^e) \text{ GHz}$		
<i>hyperfine</i>	$A_{\perp}^{(g)} = -1.7^c (-2.70^d) \text{ MHz}$	$D_1^{(e)} = -290^{\text{p.w.}} (200/\sqrt{2}^e) \text{ MHz}$	<i>ZFS</i>	
	$A_{\parallel}^{(e)} = -41^{\text{p.w.}} (-40^f) \text{ MHz}$	$D_2^{(e)} = +901^{\text{p.w.}} (1550/2^e) \text{ MHz}$		
	$A_{\perp}^{(e)} = -27^{\text{p.w.}} (-23^g) \text{ MHz}$	$Q^{(g)} = -5.37^h (-4.95^h) \text{ MHz}$		
	$A_1^{(e)} = -92^{\text{p.w.}} (\text{n.a.}) \text{ kHz}$	$Q^{(e)} = -3.91^{\text{p.w.}} (\text{n.a.}) \text{ MHz}$		<i>quadrupole interaction</i>
	$A_2^{(e)} = +49^{\text{p.w.}} (\text{n.a.}) \text{ kHz}$	$Q_1^{(e)} = +17.3^{\text{p.w.}} (\text{n.a.}) \text{ kHz}$		
<i>radiative lifetime of $^3E\rangle$</i>	$\tau_{\text{rad.}}^{-1} = (12 \text{ ns}) (83^i) \text{ MHz}$	$Q_2^{(e)} = +17.2^{\text{p.w.}} (\text{n.a.}) \text{ kHz}$		

(g) ground state triplet: $|^3A_2\rangle$

(e) excited level triplet: $|^3E\rangle$

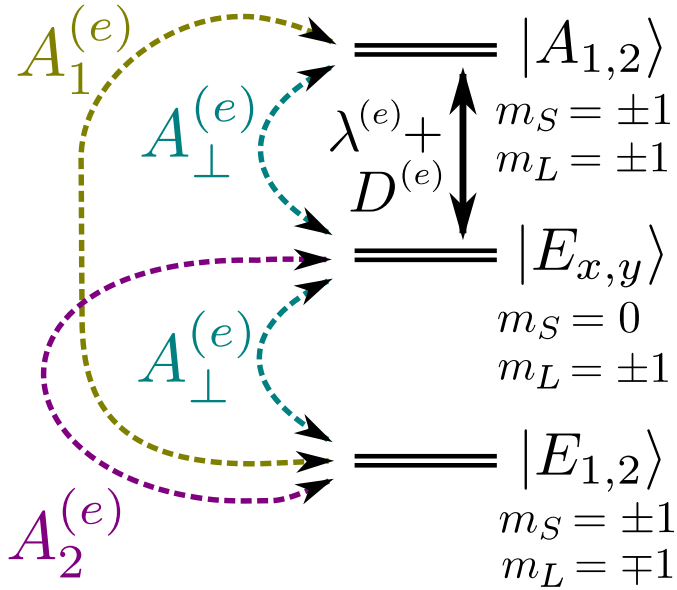
“p” and “q” Ham reduction factors within Jahn-Teller theory are required

p.w.: present work

Processes that flip the ^{14}N nuclear spin

(e) - excited ^3E state

(a) hyperfine interaction

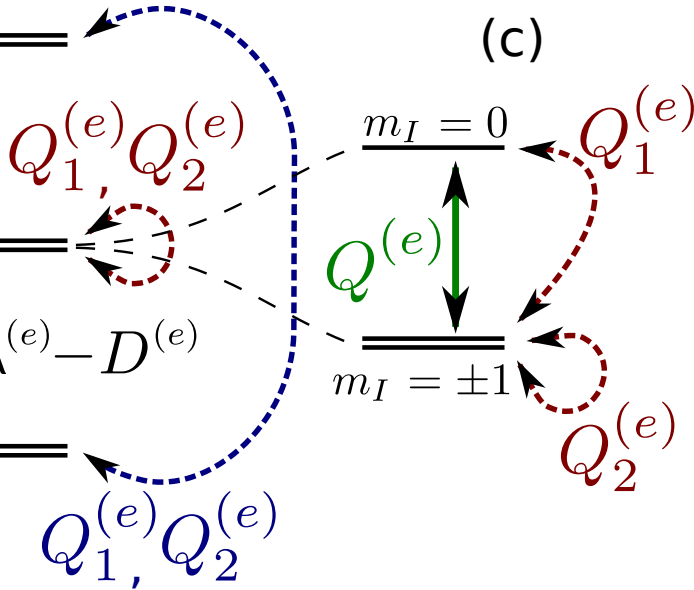


$$A_2^{(e)} (\hat{S}_- \hat{I}_- \hat{\sigma}_- + \dots)$$

$$A_1^{(e)} (\hat{I}_z \hat{S}_+ \hat{\sigma}_- + \dots)$$

$|A_{1,2}/D| \sim 10^{-6}$ negligible...

(b) quadrupolar interaction



flips ^{14}N spin by $\Delta m_I = \pm 1$ flip

$$A_{\perp}^e \hat{S}_+ \hat{I}_-$$

$$|A_{\perp}/D| \sim 0.01$$

$$Q_1^{(e)} (\hat{I}_z \hat{I}_+ \hat{\sigma}_- + \dots)$$

$$|Q_1/Q| \sim 0.004$$

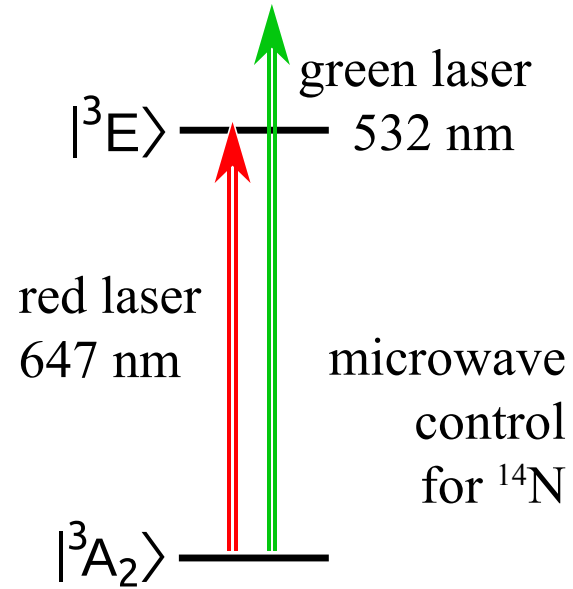
$\Delta m_I = \pm 2$ flip

$$Q_2^{(e)} (\hat{I}_-^2 \hat{\sigma}_- + \dots)$$

ESLAC: excited state level anticrossing

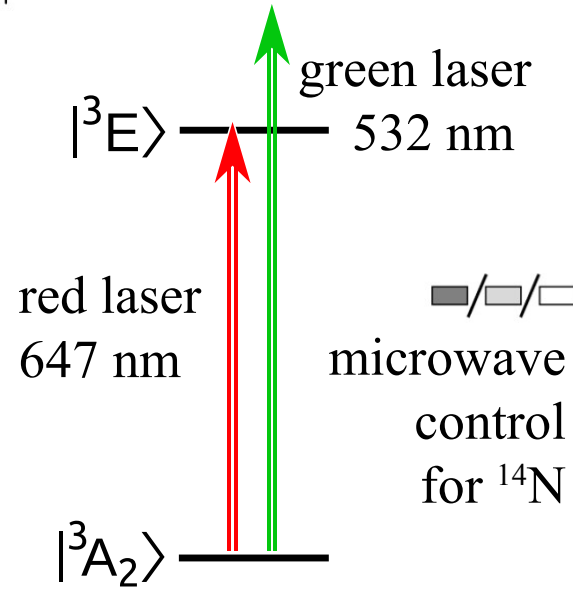
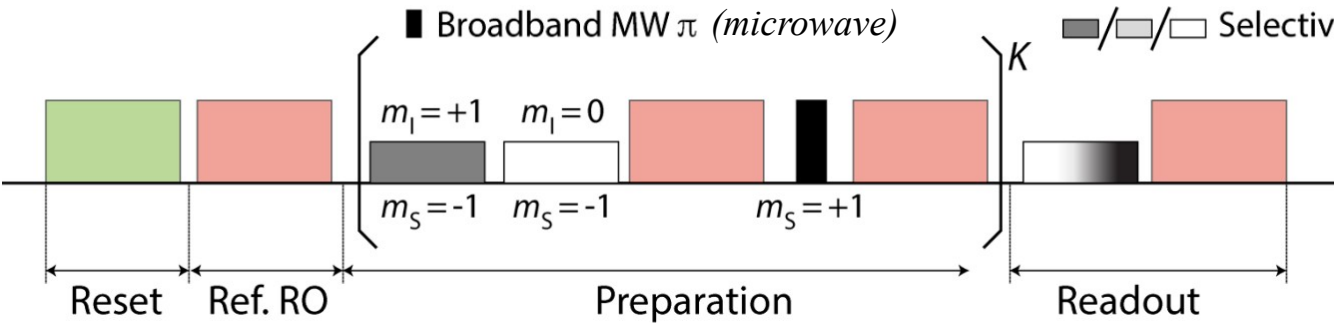
without applied magnetic field

^{14}N nuclear spin manipulation

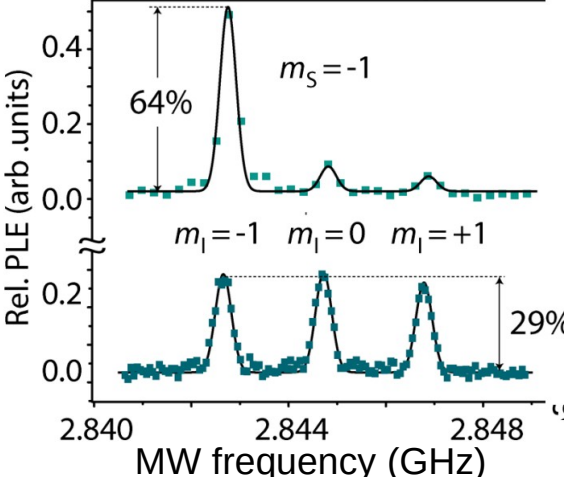


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experimental

Nuclear spin flip probabilities



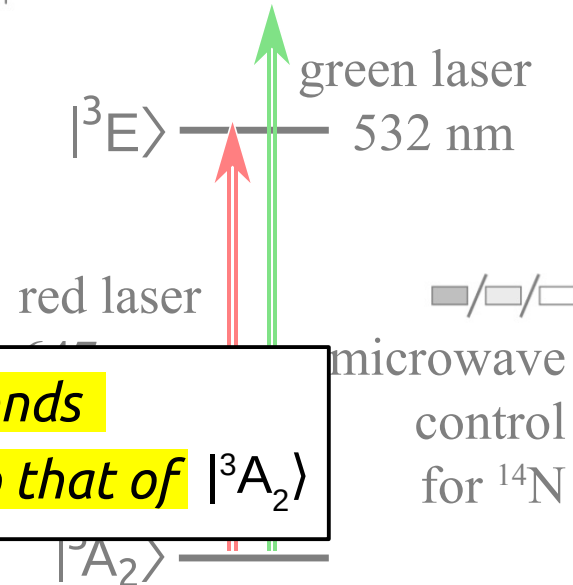
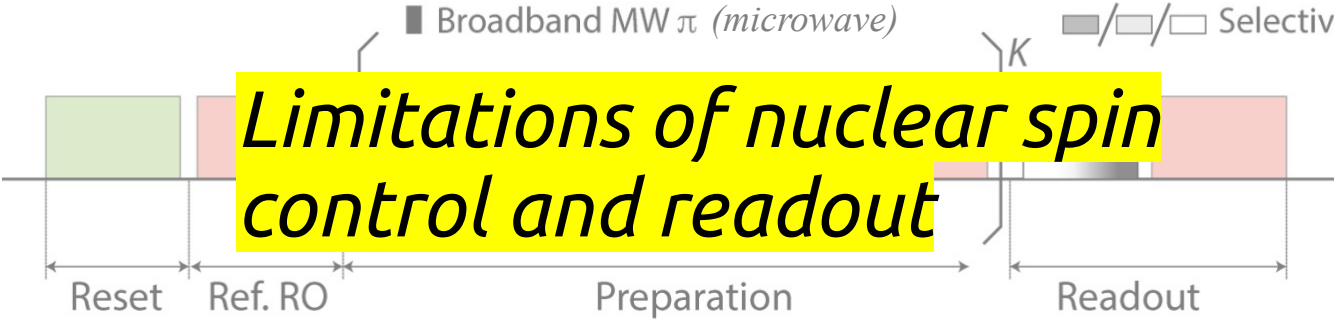
- 1) **Reset:** Initialize NV(-)'s spin to $m_S = 0$
- 2) **Preparation:** Initialize ^{14}N nuclear spin to $m_I = +1$ or -1 or 0
- 3) **Readout:** Optical readout of m_S and m_I both ($\sim 20 \mu\text{s}$)



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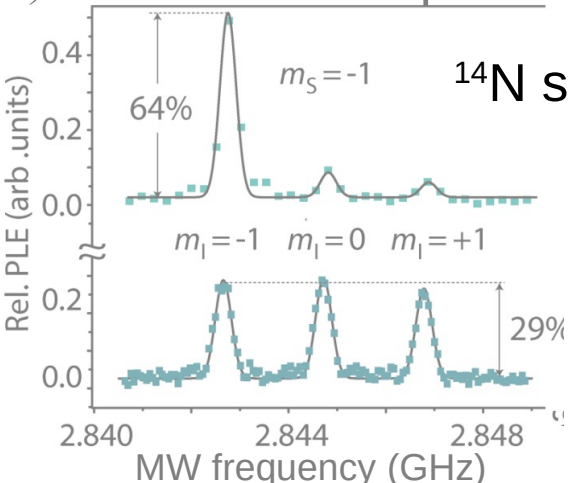
Nuclear spin flip probabilities

Limitations of nuclear spin control and readout



- 1) **Reset:** Initialize NV(-)'s spin to $m_S = 0$
- 2) **Preparation:** Initialize ^{14}N nuclear spin
- 2) **Readout:** Optical readout of m_S

this is not hours or seconds compared to T_1 time to that of $|^3\text{A}_2\rangle$



^{14}N spin relaxes: $\sim 20 \mu\text{s}$ total stay within $|^3\text{E}\rangle$ or $N \sim 1000$ opt. cycles (red laser)

see also: npj Quantum Inf 3, 33 (2017)

transition	theory	expt.
$ m_I = \pm 1\rangle \rightarrow m_I = \mp 1\rangle$	0.11	0.18(3)
$ m_I = \pm 1\rangle \rightarrow m_I = 0\rangle$	0.08	0.16(3)
$ m_I = 0\rangle \rightarrow m_I = \pm 1\rangle$	0.16	0.25(3)

Summary

☞ Complete ab-initio theory for \overleftrightarrow{D} , \overleftrightarrow{Q} , \overleftrightarrow{A} tensors (spin-spin ZFS, quadrupole, hyperfine)

→ not only trivial D , Q , A_{\perp} , A_{\parallel} terms

→ including the nontrivial, orbital driven D_1 , D_2 , Q_1 , Q_2 , A_1 , A_2 parameters

☞ During optical cycles within $|^3E\rangle$:

→ “new” ^{14}N nuclear spin relaxation channels open

$$\Delta m_I = \pm 1$$

by $A_{\perp}^e \hat{S}_+ \hat{I}_-$
(hyperfine)

$$\Delta m_I = \pm 2$$

by $Q_2^{(e)} (\hat{I}_-^2 \hat{\sigma}_- + \dots)$
(nuclear quadrupole)

orbital-nuclear spin interaction

☞ Limitations for optical control and readout for nuclear spin

→ ^{14}N nuclear spin relaxes much faster in $|^3E\rangle$ than that in $|^3A_2\rangle$

→ longer than $\sim \mu\text{s}$ total stay within $|^3E\rangle$ (*readout & preparation*)

→ or more than $N > \sim 1000$ opt. cycles

Summary

- ☞ Complete ab-initio theory for \overleftrightarrow{D} , \overleftrightarrow{Q} , \overleftrightarrow{A} tensors (spin-spin ZFS, quadrupole, hyperfine)
→ not only trivial D , Q , A_{\perp} , A_{\parallel} terms
→ including the nontrivial, orbital driven D_1 , D_2 , Q_1 , Q_2 , A_1 , A_2 parameters

Outlook: Apply on “G4V” defects
SiV(-), GeV(-), SnV(-), PbV(-)

- ☞ During optical cycles within $|^3E\rangle$:
→ “new” ^{14}N nuclear spin relaxation channels open

$$\Delta m_I = \pm 1 \quad \text{by } A_{\perp}^e \hat{S}_+ \hat{I}_- \quad (\text{hyperfine})$$

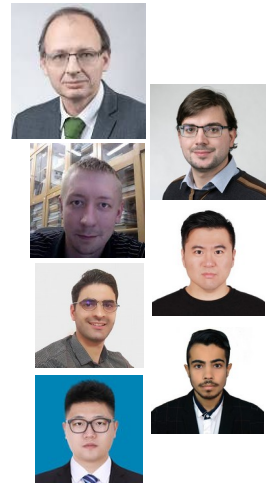
$$\Delta m_I = \pm 2 \quad \text{by } Q_2^{(e)} (\hat{I}_-^2 \hat{\sigma}_- + \dots) \quad (\text{nuclear quadrupole})$$

orbital-nuclear spin interaction

- ☞ Limitations for optical control and readout for nuclear spin
→ ^{14}N nuclear spin relaxes much faster in $|^3E\rangle$ than that in $|^3A_2\rangle$
→ longer than $\sim \mu\text{s}$ total stay within $|^3E\rangle$ (readout & preparation)
→ or more than $N > \sim 1000$ opt. cycles

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Song Li
Meysam Mohseni
Nima Ghafaricherati
Bian Guodong



**Thank you for your
kind attention!**

experimental:
Phys. Rev. Lett. 131, 236901 (2023)

theoretical model:

doi.org/10.48550/arXiv.2402.19418

*Resonant versus nonresonant spin
readout of a NV center in diamond under
cryogenic conditions*

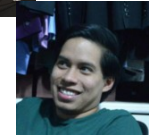
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^{13}C hyperfine for G4V centers in diamond

*Hyperfine on ^{13}C sites will be entangled to **orbital** degrees of freedom too!*

$$\hat{W} = \overleftrightarrow{S} \overleftrightarrow{A}_0 \overrightarrow{I} + q(\overleftrightarrow{S} \overleftrightarrow{A}_x \overrightarrow{I} \hat{\sigma}_z + \overleftrightarrow{S} \overleftrightarrow{A}_y \overrightarrow{I} \hat{\sigma}_x)$$

three 3x3 hyperfine matrices:

\overleftrightarrow{A}_0 : “normal” hyperfine

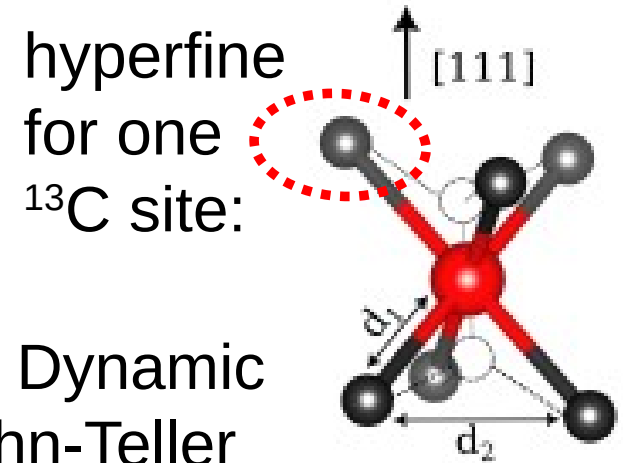
$\overleftrightarrow{A}_x, \overleftrightarrow{A}_y$: “orbital” hyperfine for one ^{13}C site

should be visible below $\sim 20\text{ K}$

$> 50\text{ K}$ orbital averaging should occur

similarly to that of NV's 3E excited state

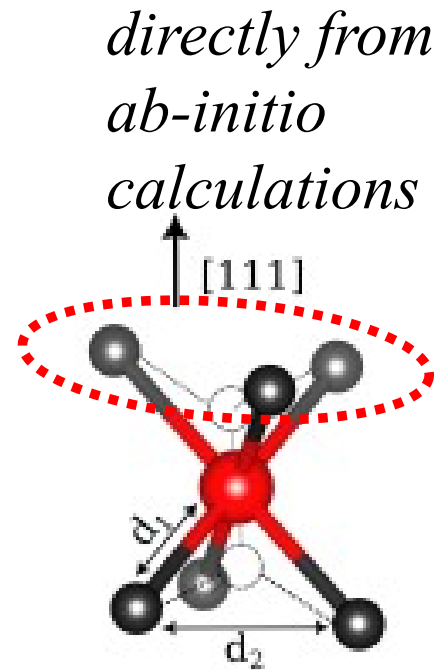
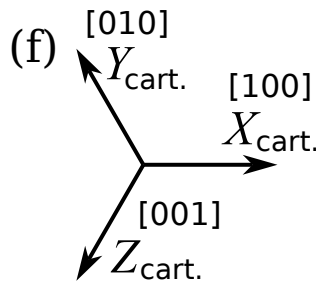
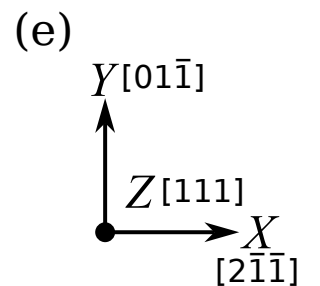
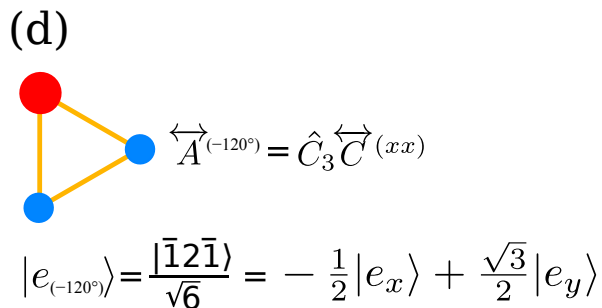
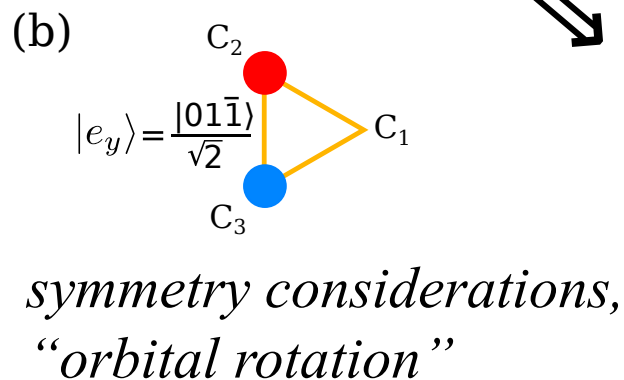
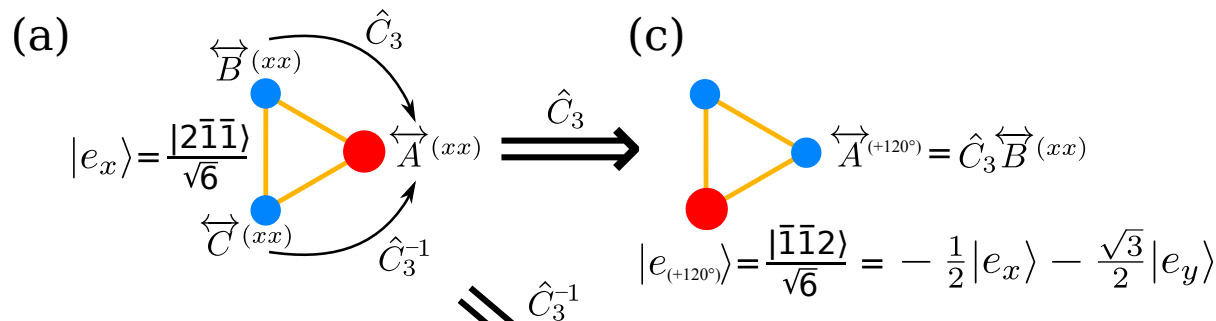
$\langle \hat{\sigma}_z \rangle = \langle \hat{\sigma}_x \rangle = 0$ (thermal average)



$$\hat{\sigma}_x = \begin{pmatrix} & 1 \\ 1 & \end{pmatrix} = |e_x\rangle\langle e_y| + |e_y\rangle\langle e_x|$$
$$\hat{\sigma}_z = \begin{pmatrix} 1 & \\ & -1 \end{pmatrix} = |e_x\rangle\langle e_x| - |e_y\rangle\langle e_y|$$

How to interpret ab-initio data

Three different \overleftrightarrow{A} , \overleftrightarrow{B} , \overleftrightarrow{C} hyperfine tensors on 3 equivalent ^{13}C sites



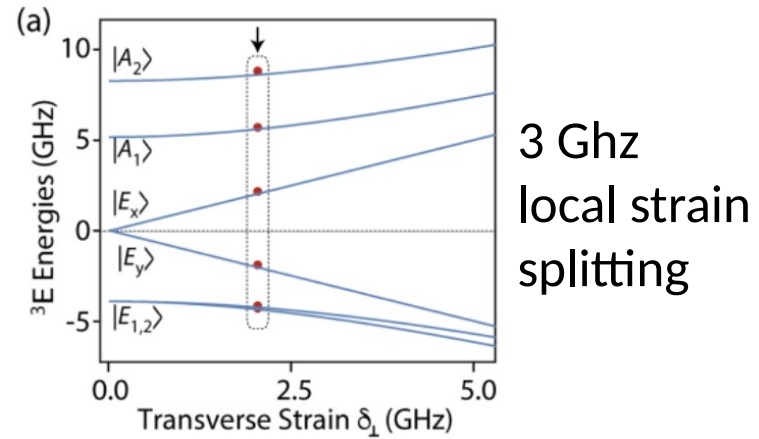
static Jahn-Teller distortion $\rightarrow e_x, e_y$ orbitals split the hyperfine tensors too!

Nuclear spin flip probabilities

Coherent time evolution for Q_2

$$p(|E_y\rangle \otimes |\mp\rangle \rightarrow |E_y\rangle \otimes |\pm\rangle) = (Q_2^{(e)} n \tau_{\text{rad}})^2 = 0.109,$$

$$p(|E_y\rangle \otimes |\mp\rangle \rightarrow |E_x\rangle \otimes |\pm\rangle) = 0,$$



Fermi's golden rule for hyperfine transitions

$$p(|E_x\rangle \otimes |0\rangle \rightarrow |E_{1,2}\rangle \otimes |\pm 1\rangle) = 2 \times \left(\frac{A_{\perp}^{(e)}}{\lambda^{(e)} - D^{(e)}} \right)^2 \times n = 0.112,$$

$$p(|E_x\rangle \otimes |0\rangle \rightarrow |A_1\rangle \otimes |\pm 1\rangle) = 2 \times \frac{1}{2} \left(\frac{A_{\perp}^{(e)}}{\lambda^{(e)} + D^{(e)} - D_2^{(e)}} \right)^2 \times n = 0.017,$$

$$p(|E_x\rangle \otimes |0\rangle \rightarrow |A_2\rangle \otimes |\pm 1\rangle) = 2 \times \frac{1}{2} \left(\frac{A_{\perp}^{(e)}}{\lambda^{(e)} + D^{(e)} - D_2^{(e)}} \right)^2 \times n = 0.021,$$