

Simphony

**our python package to simulate
point defect dynamics**

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Budapest, Hungary

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Funding acknowledgment



AI4QT



PROJECT
FINANCED FROM
THE NRDI FUND



AI4QT partners (Hungary, Germany)



Simphony



Qutility



qutility.io

Our python package to simulate point defect quantum dynamics
for quantum technology applications

Quantum Computing

*Universal high-fidelity quantum gates
for spin-qubits in diamond*

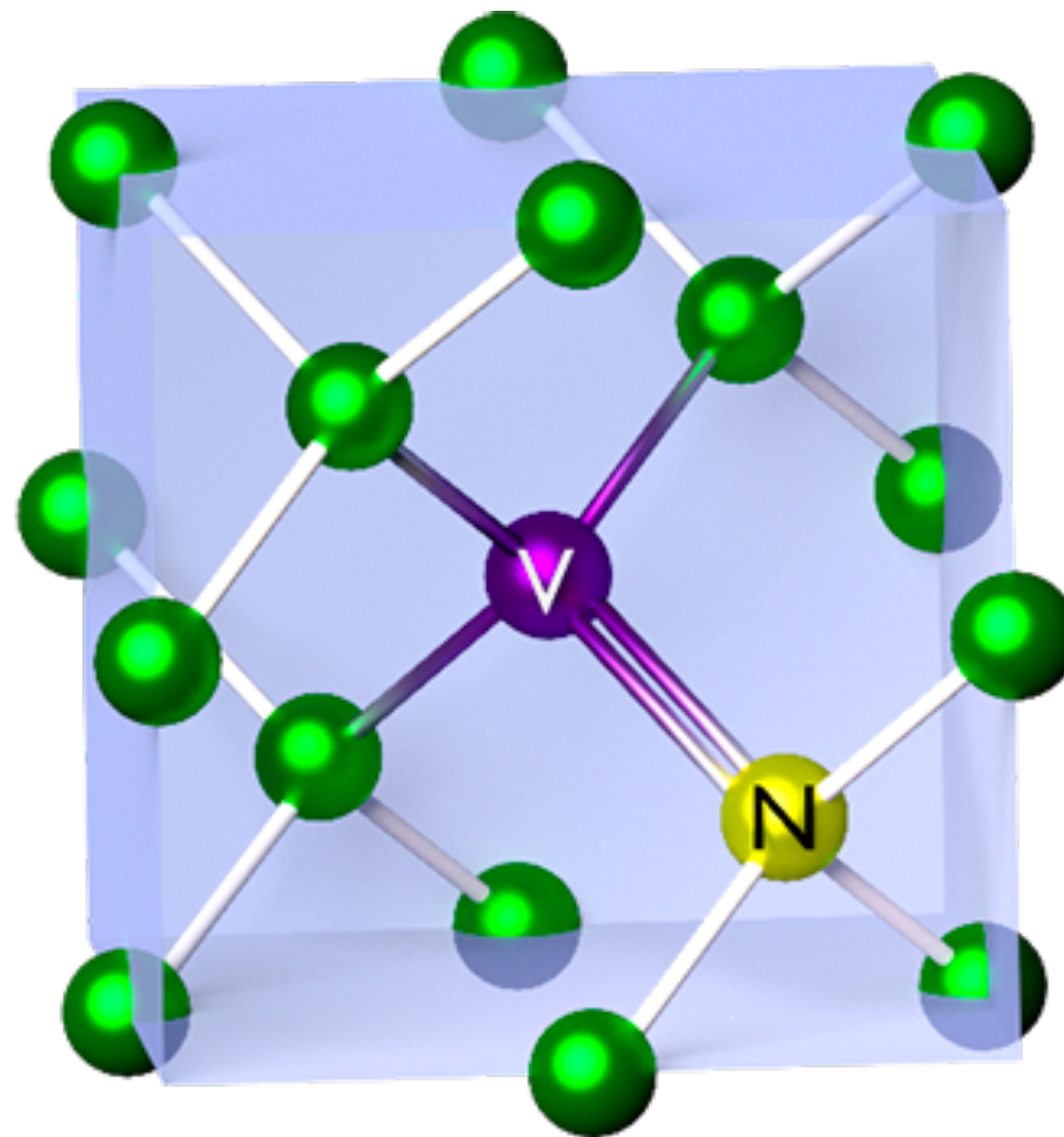
<https://arxiv.org/abs/2403.10633>

Taminiau group, Delft

*High-Fidelity Electron Spin Gates
in a Scalable Diamond Quantum Register*

<https://arxiv.org/abs/2406.04199>

Jelezko group, Ulm



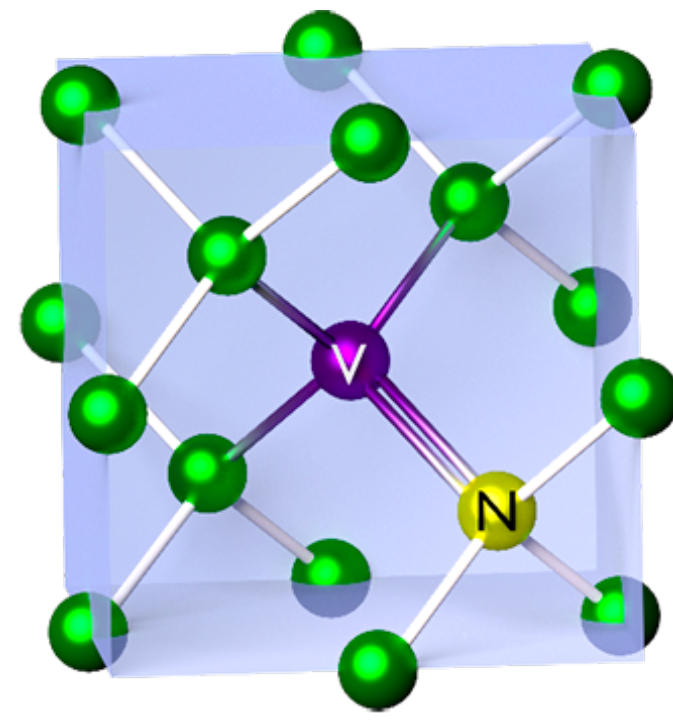
Quantum Sensing

*Nanoscale magnetism and magnetic
phase transitions in atomically thin CrSBr*

<https://arxiv.org/abs/2312.09279>

Maletinsky group, Basel

Simphony

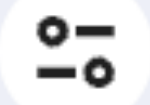


- Simphony is our python package to simulate point defect quantum dynamics for quantum technology applications
- Simphony is our “Minimal Viable Product” => **We are looking for collaborators!**

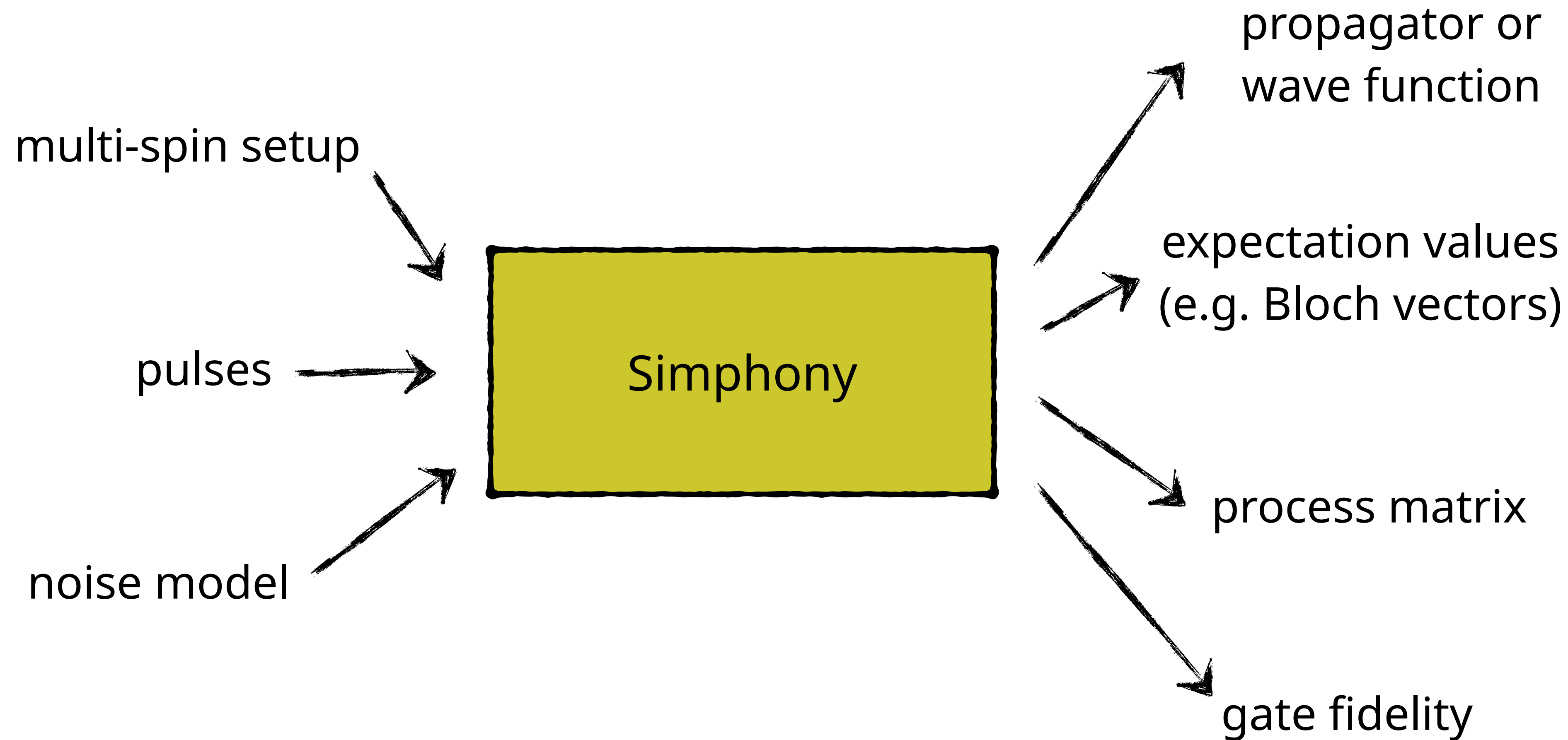
Input-output structure of Symphony



Qutality



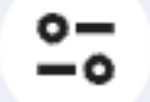
qutality.io



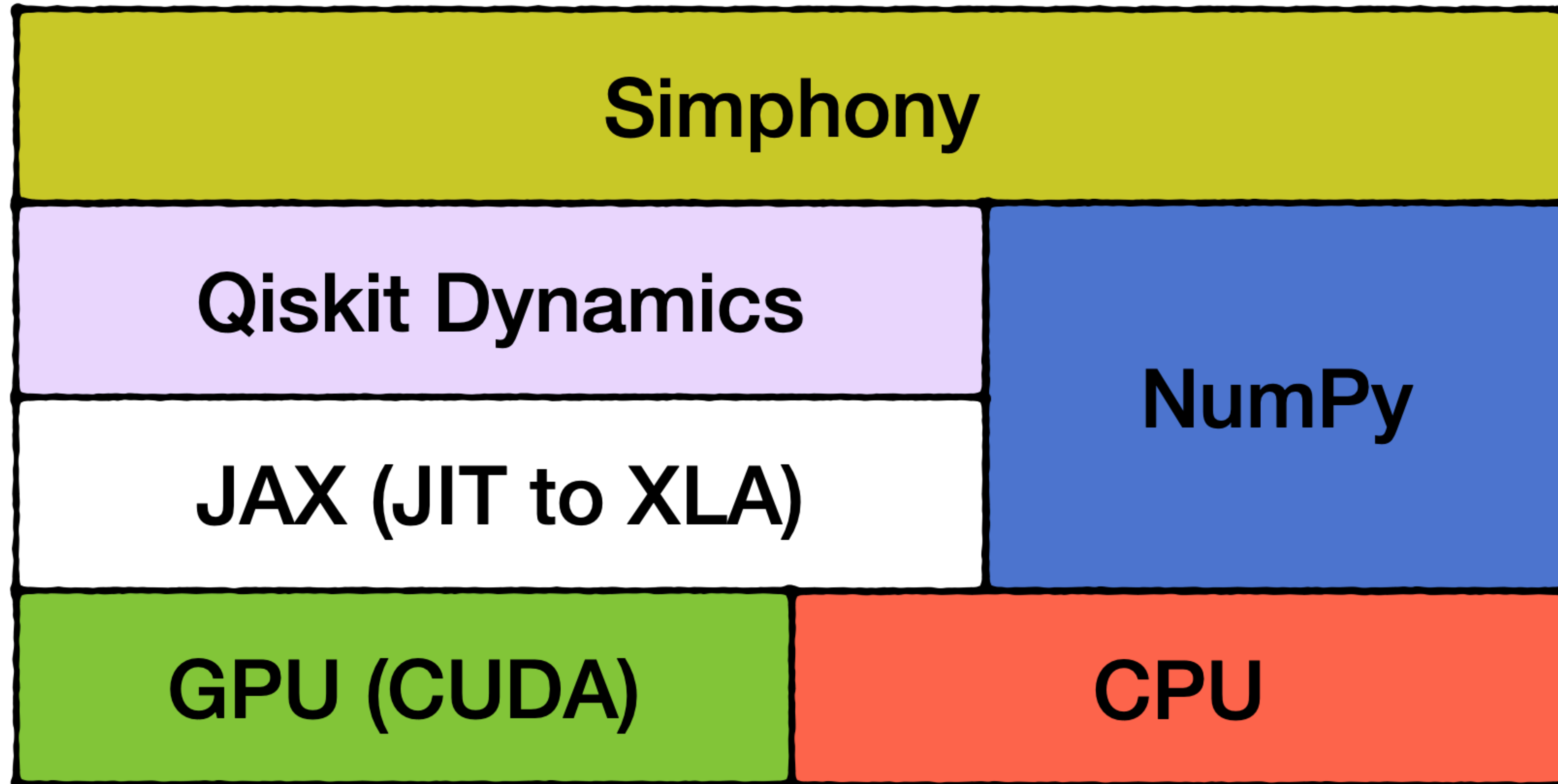
Dependency diagram of Symphony



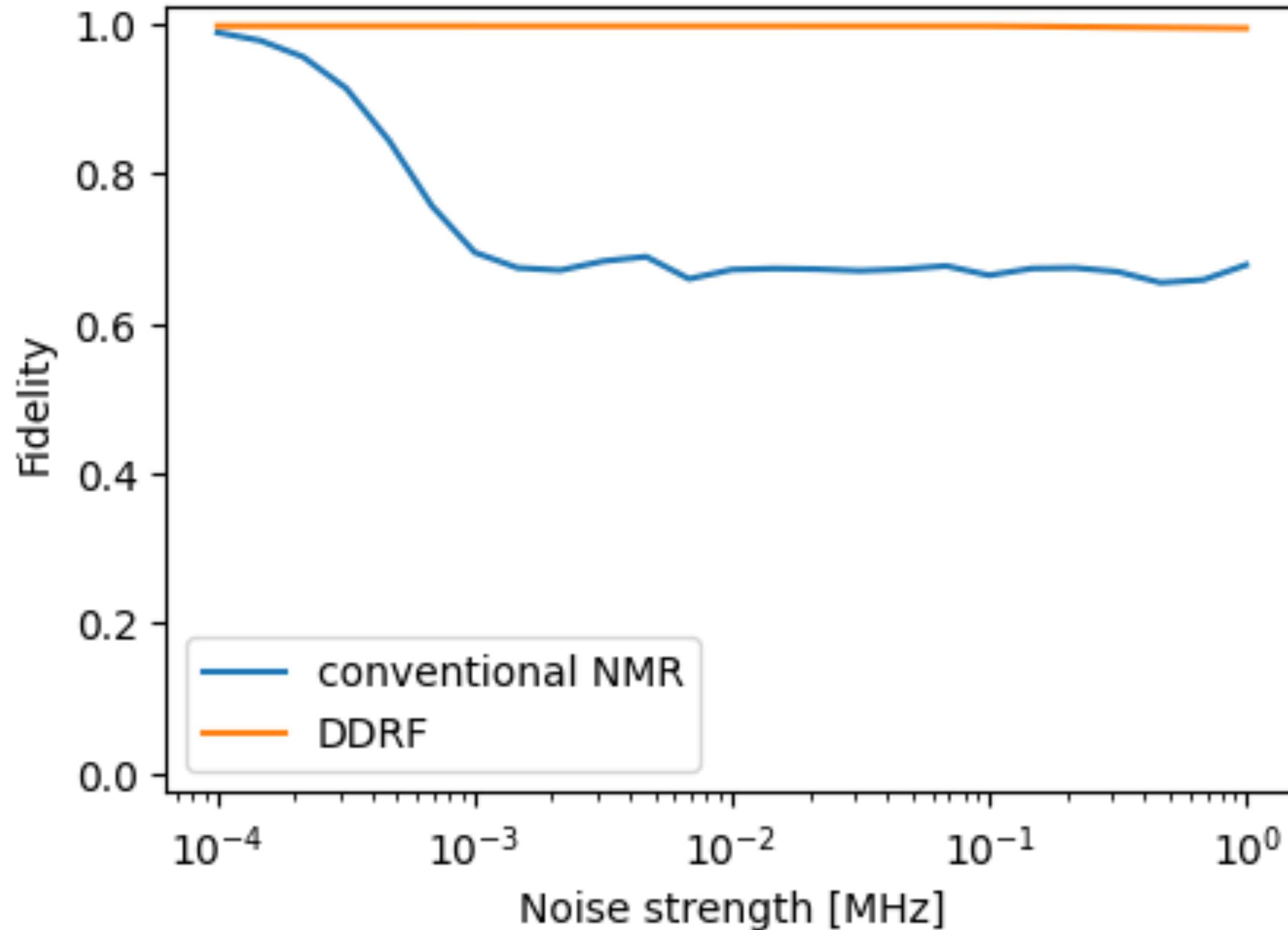
Qutlity



qutlity.io



Example: noise-resistant controlled rotation



Multi-spin setup

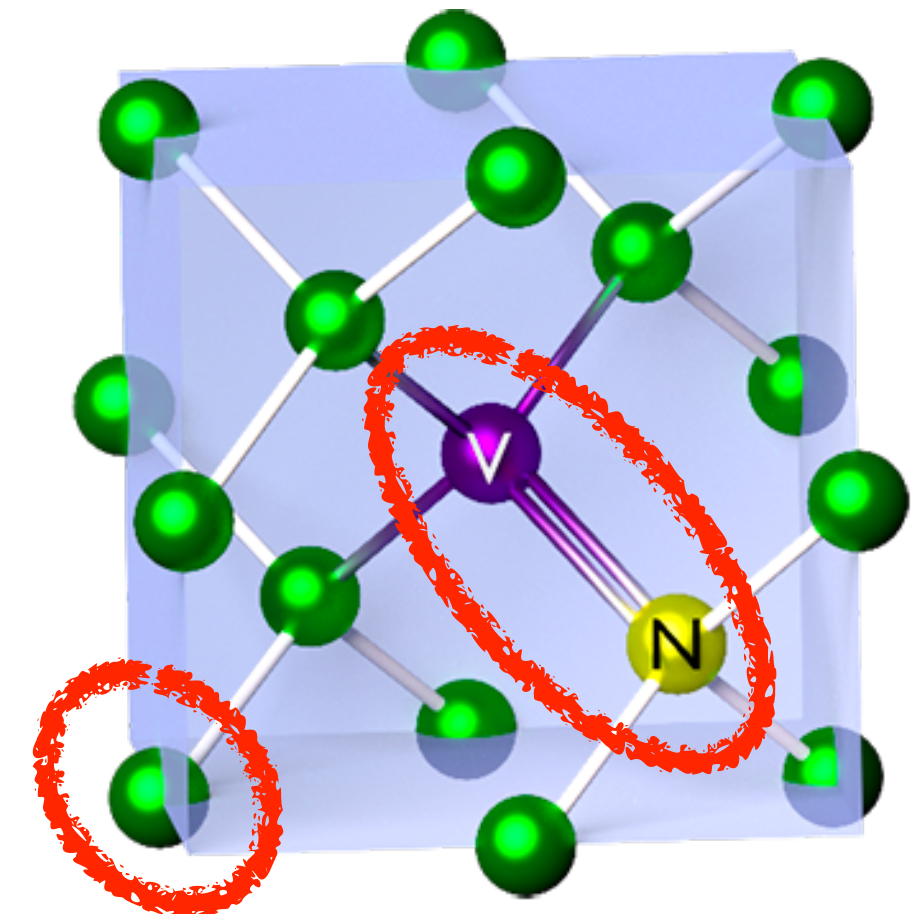
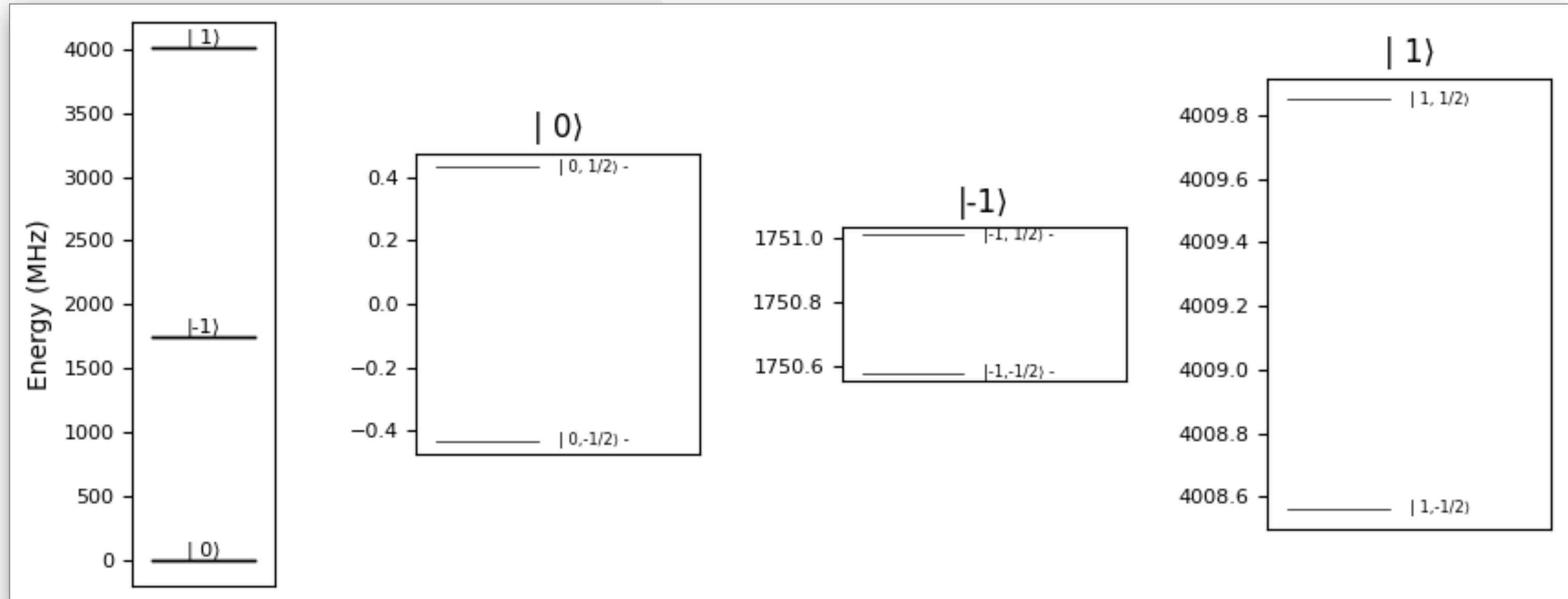
```
spin_e = simphony.Spin(  
    dimension = 3,  
    name = 'NV-e',  
    gyromagnetic_ratio = 28020, # MHz/T  
    qubit_subspace = (0, -1),  
    zero_field_splitting = 2880 # MHz  
)
```

```
spin_C = simphony.Spin(  
    dimension = 2,  
    name = '13C',  
    gyromagnetic_ratio = 10.71, # MHz/T  
    qubit_subspace = (-1/2, 1/2),  
    zero_field_splitting = 0  
)
```

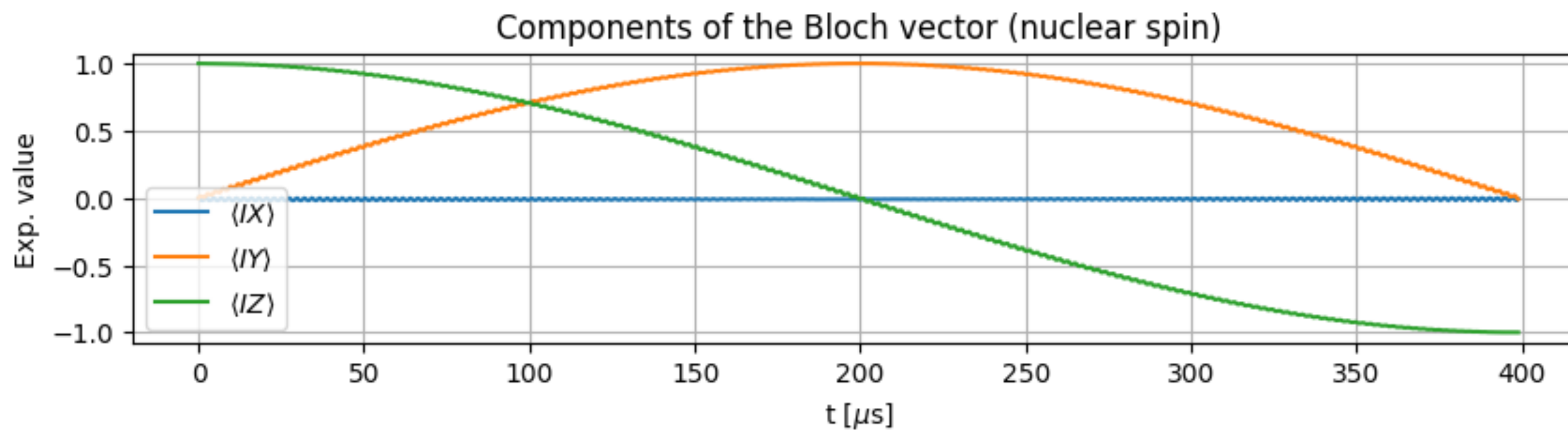
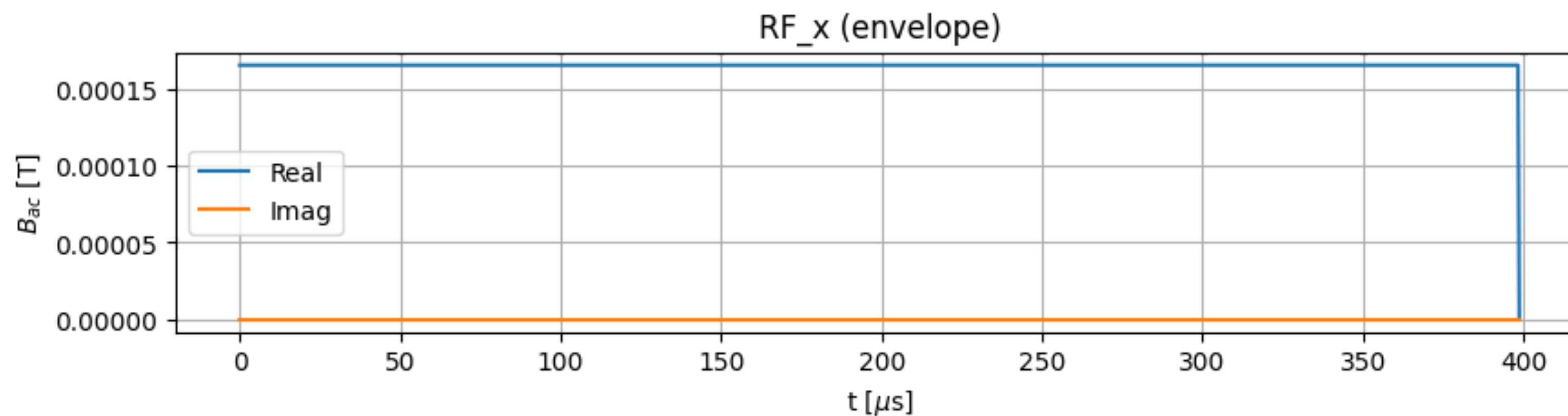
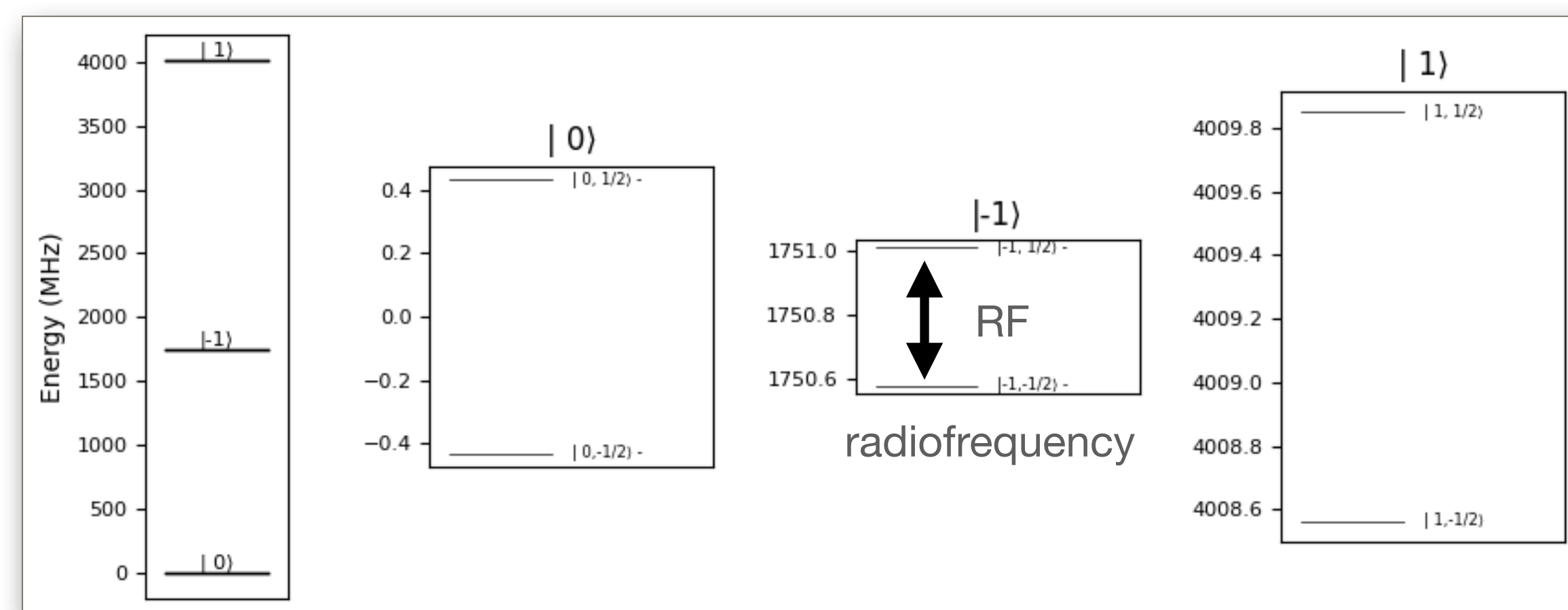
```
hyperfine = simphony.Interaction(spin_e, spin_C)  
hyperfine.zz = 0.213154 # MHz  
hyperfine.xz = 0.003 # MHz  
hyperfine.zx = 0.003 # MHz
```

```
static_field = simphony.StaticField()  
static_field.z = 0.0403 # T
```

```
driving_field_MW = simphony.DrivingField(direction = [1, 0, 0], name = 'MW_x')  
driving_field_RF = simphony.DrivingField(direction = [1, 0, 0], name = 'RF_x')
```

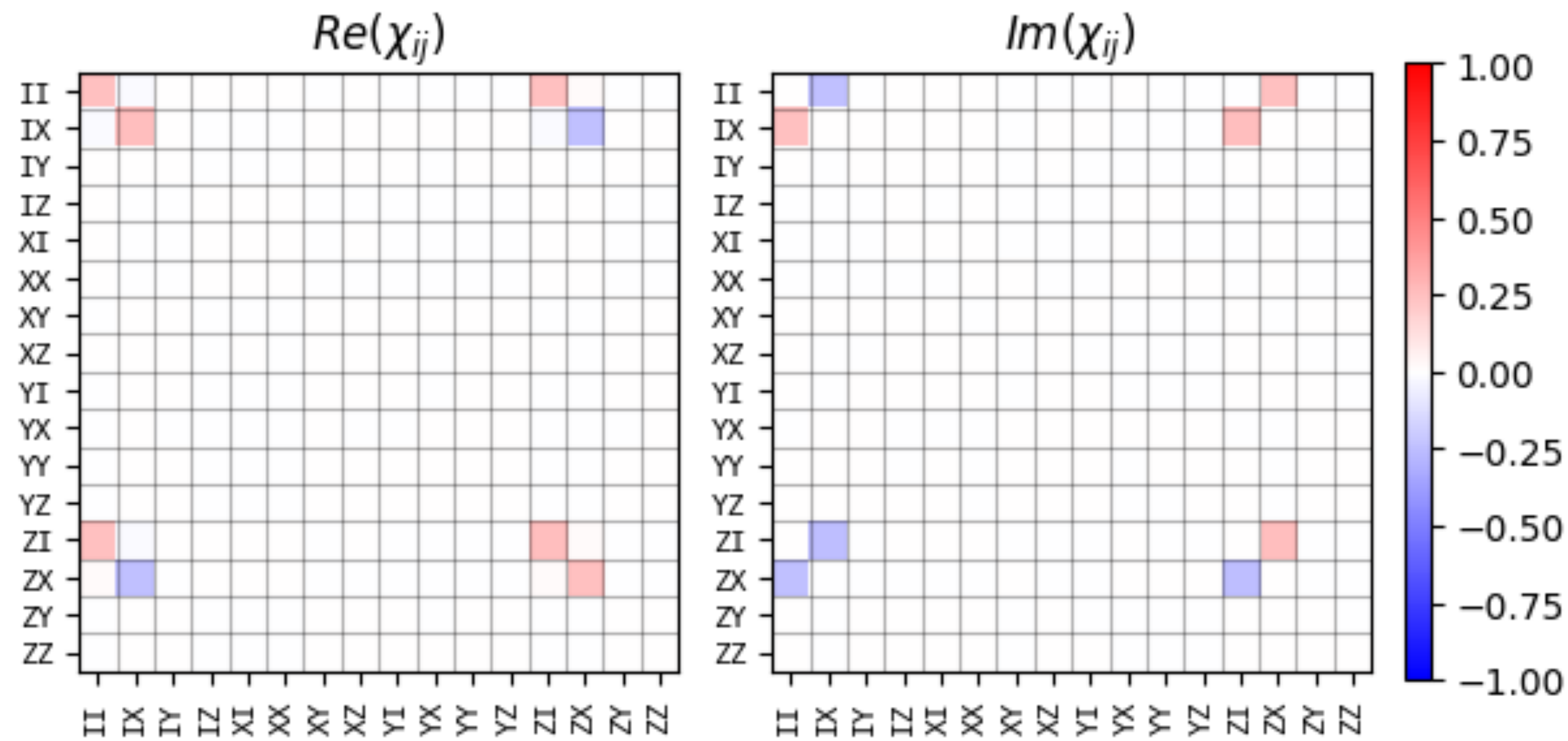


Conventional nuclear magnetic resonance



Conventional nuclear magnetic resonance

Process matrix



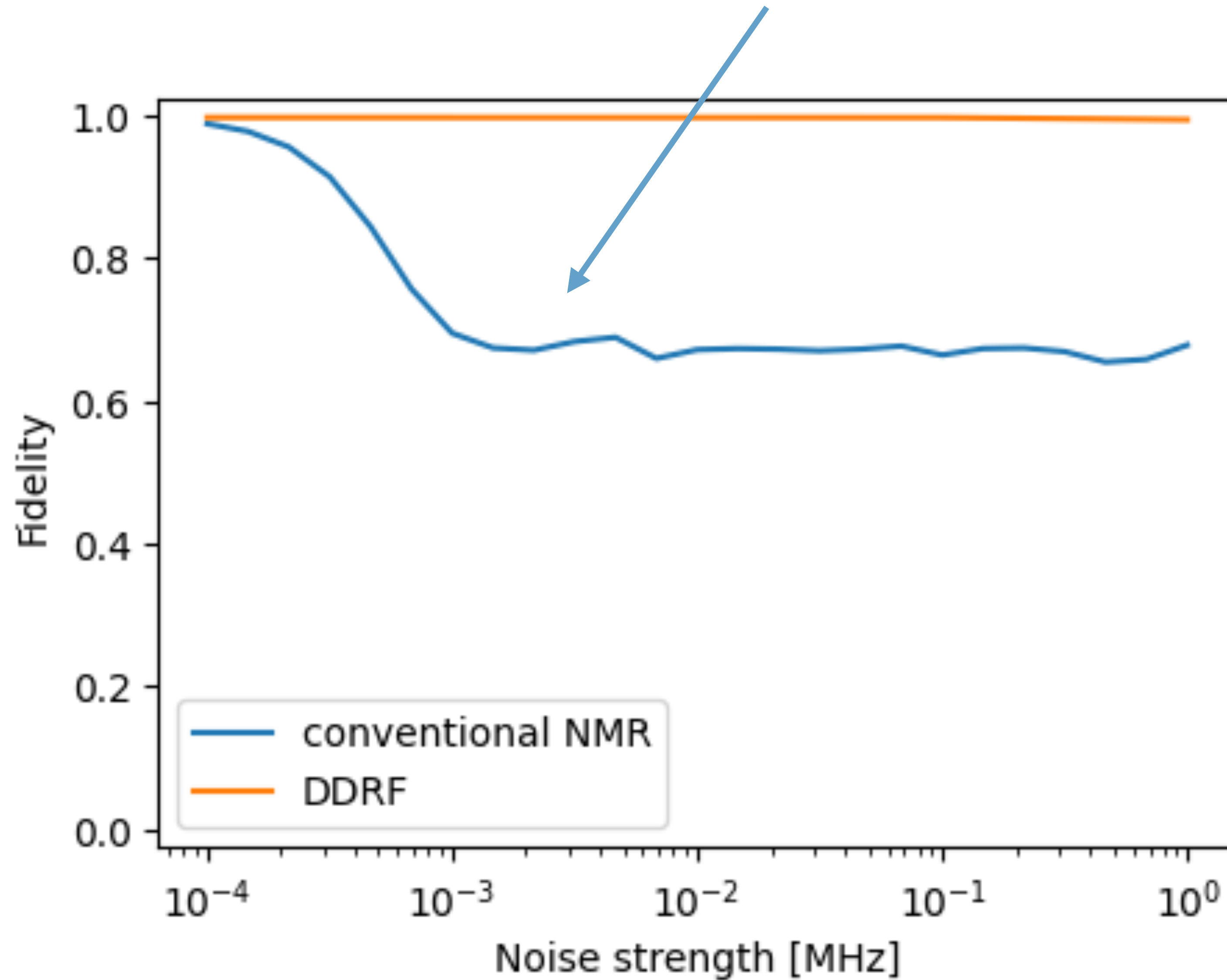
Ideal gate:

$RX(\pi)$ controlled on $m_S = -1$

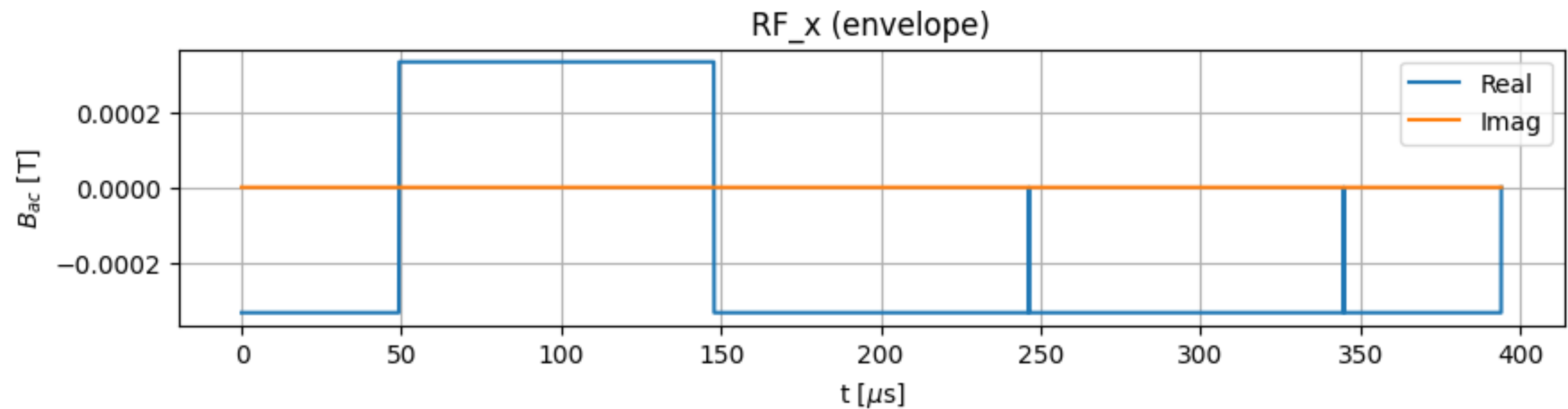
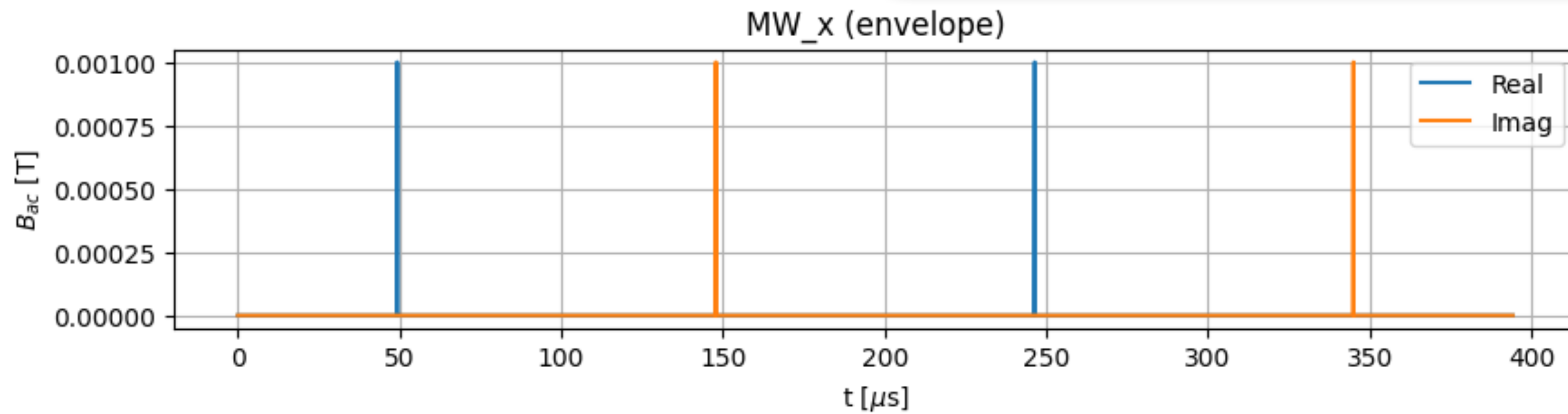
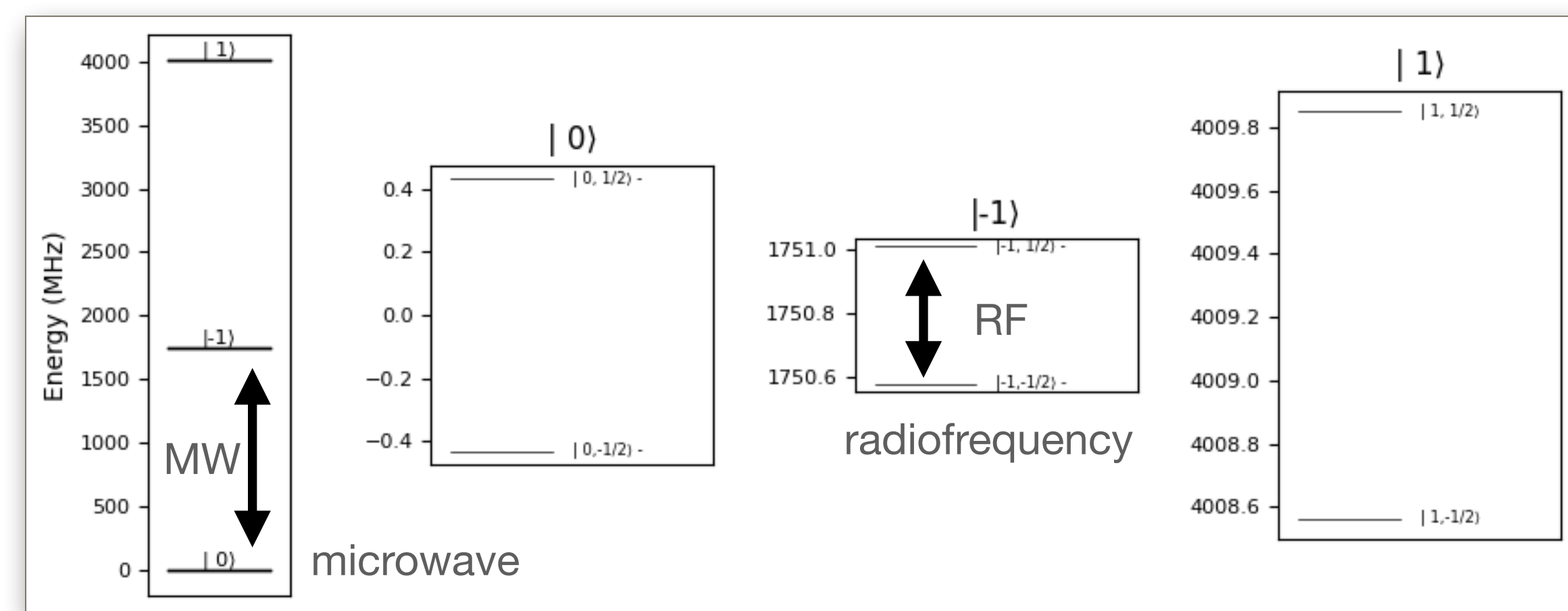
Average gate fidelity:

99.87%

Conventional NMR: vulnerable to noise

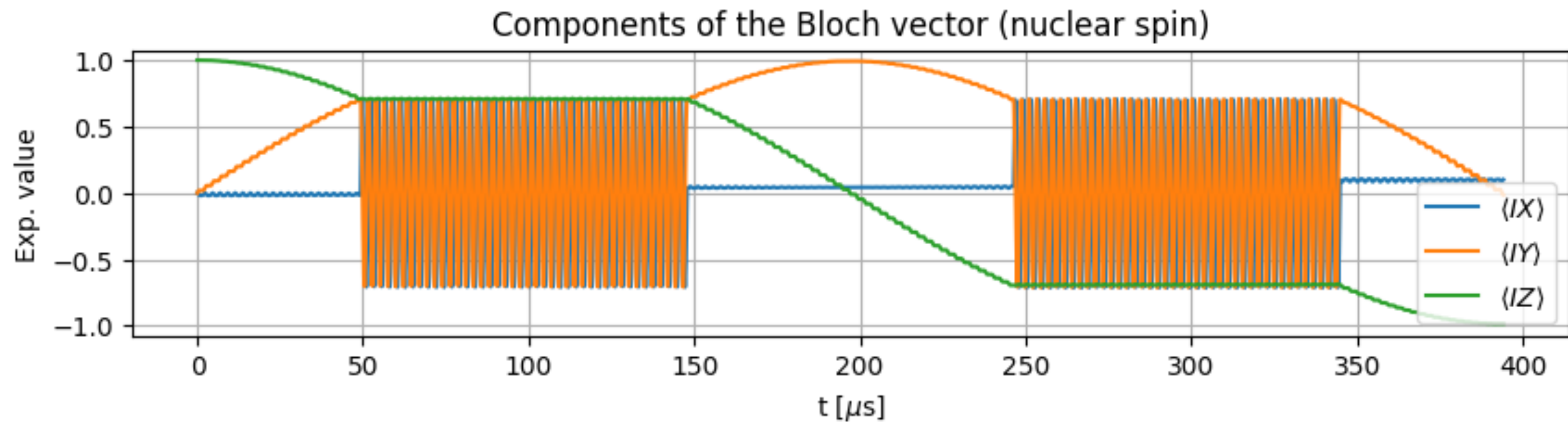
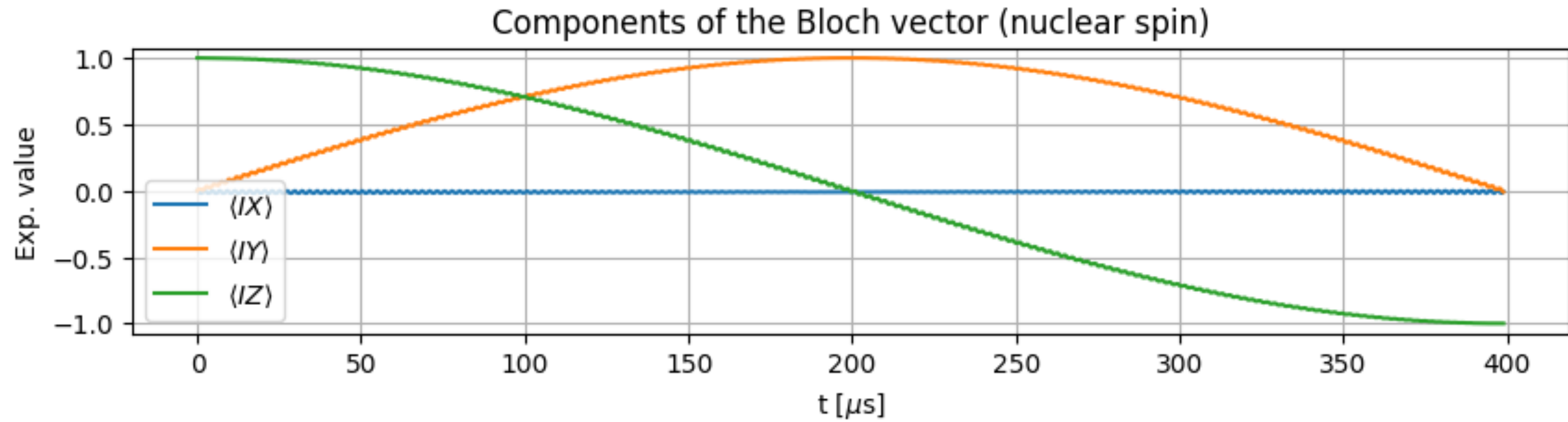


DDRF pulse sequence

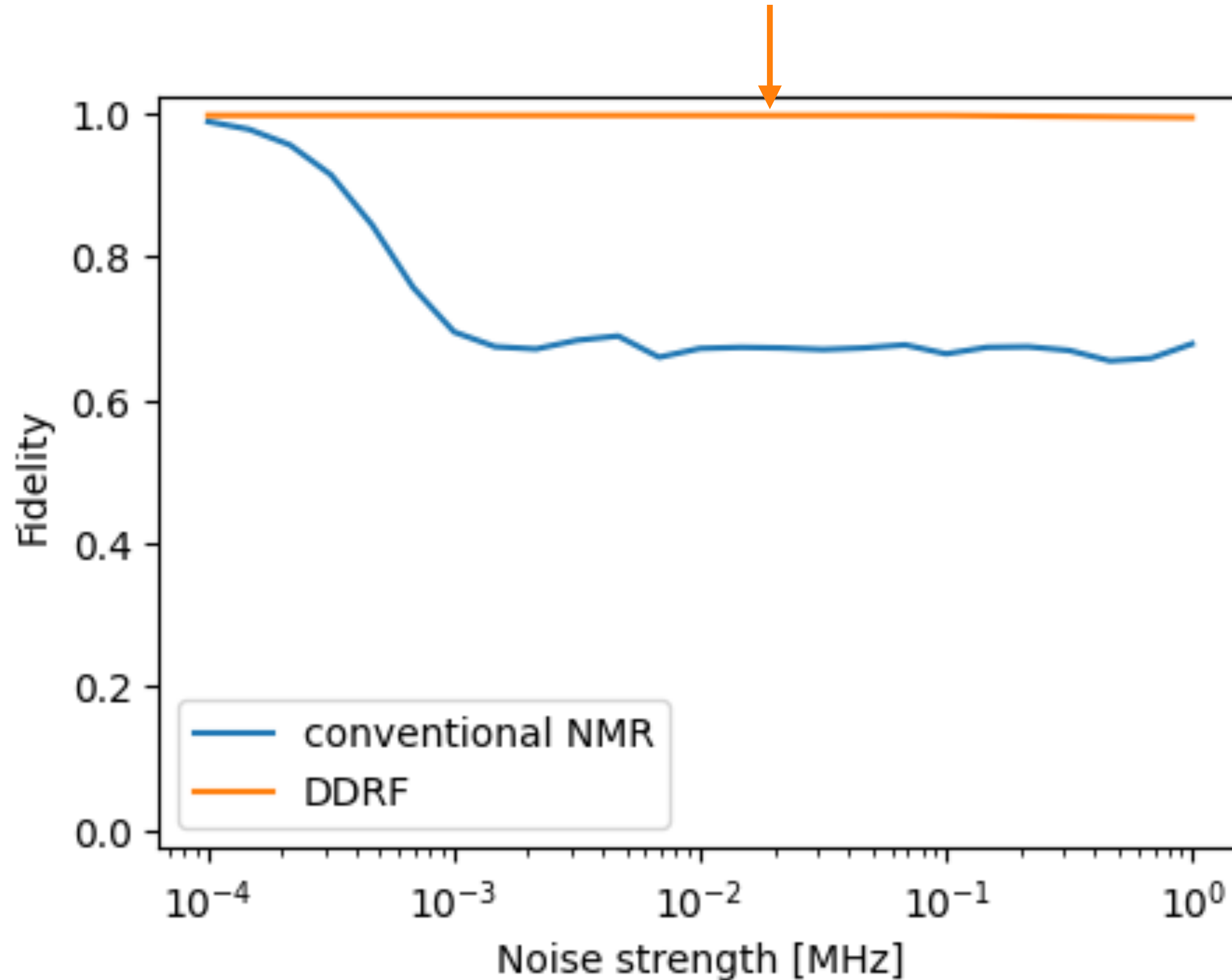


Spin dynamics due to DDRF pulse sequence

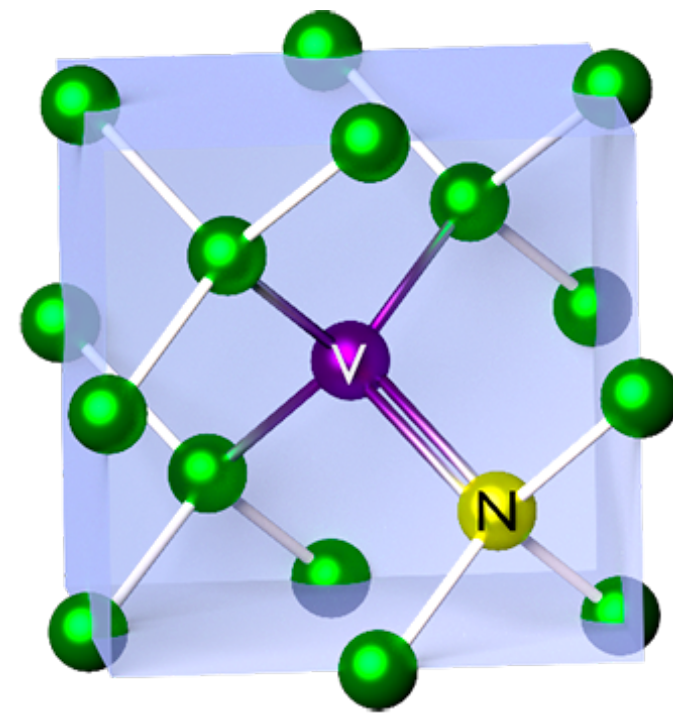
Conventional NMR



DDRF pulse sequence: **robust against noise**



Simphony



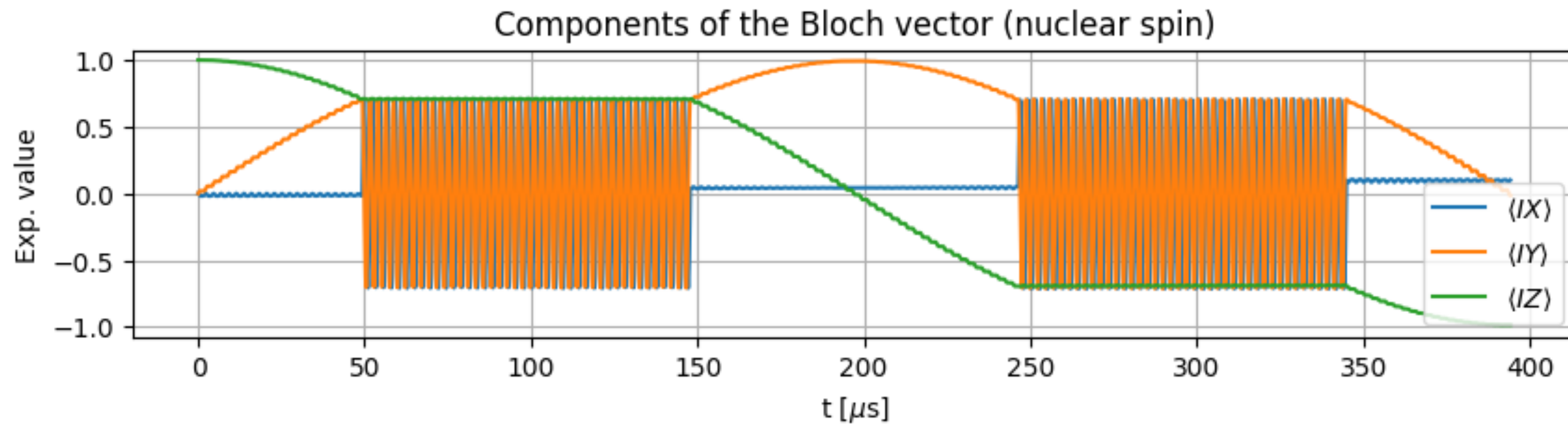
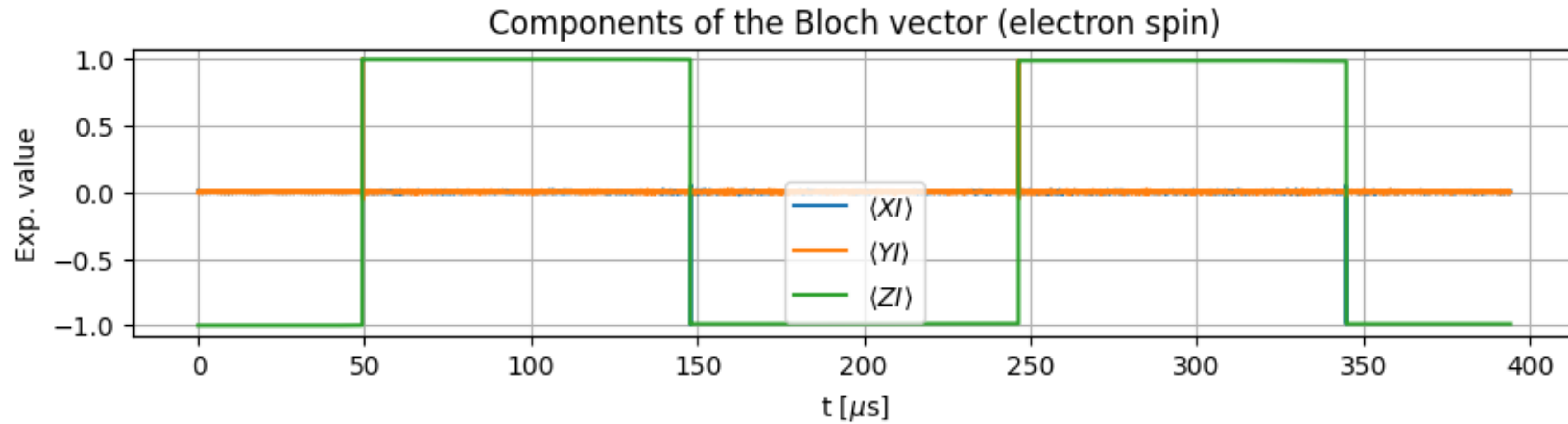
Roadmap

- Integration of ^{13}C hyperfine database.
- Integration of point-defect Hamiltonians beyond NV centers.
- Scalable registers of coupled NV centers.
- Lindblad-type noise models.
- Non-Markovian system-bath interaction.
- Photophysics dynamics: initialisation and readout.

Summary

- Simphony is our python package to simulate point defect spin dynamics for quantum technology applications.
- Simphony is our “Minimal Viable Product” => **We are looking for collaborators!**

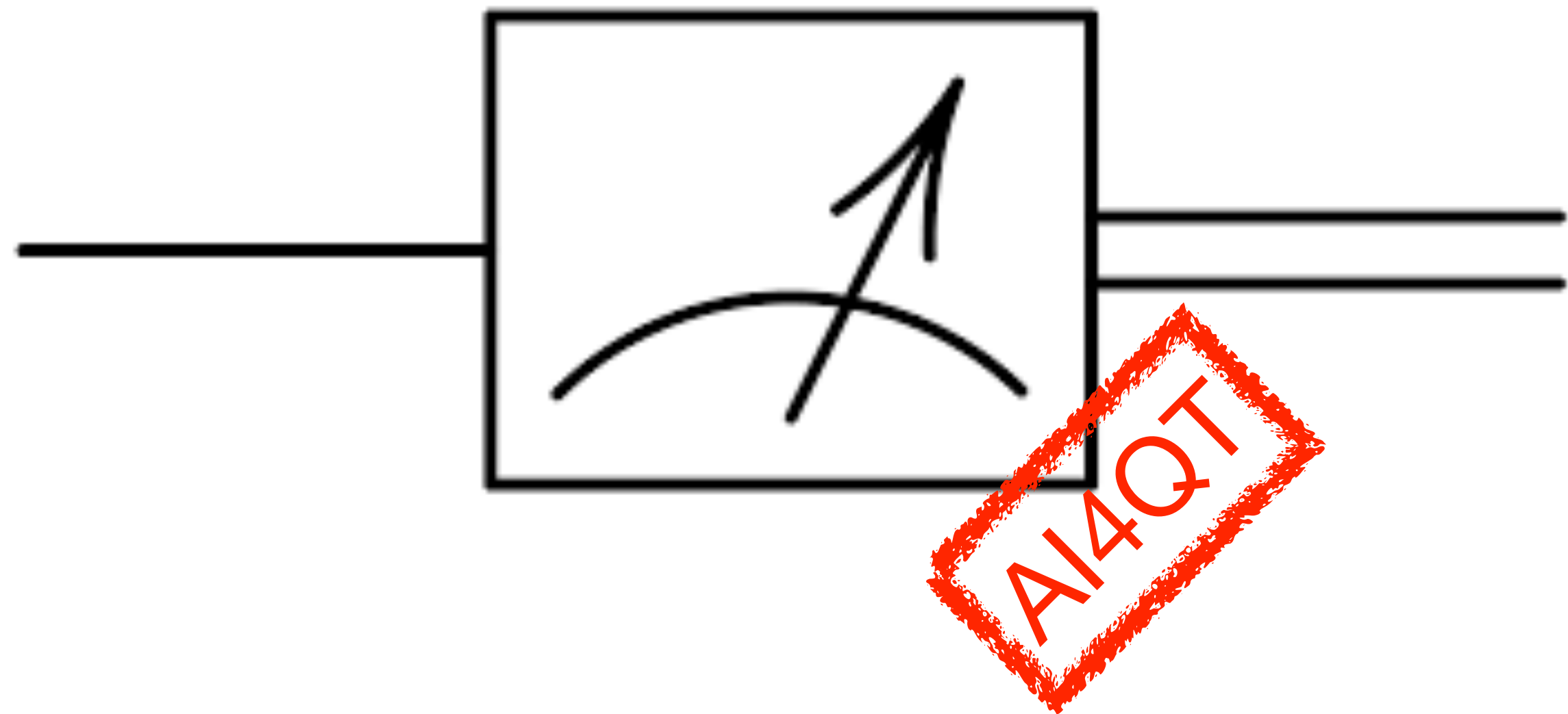
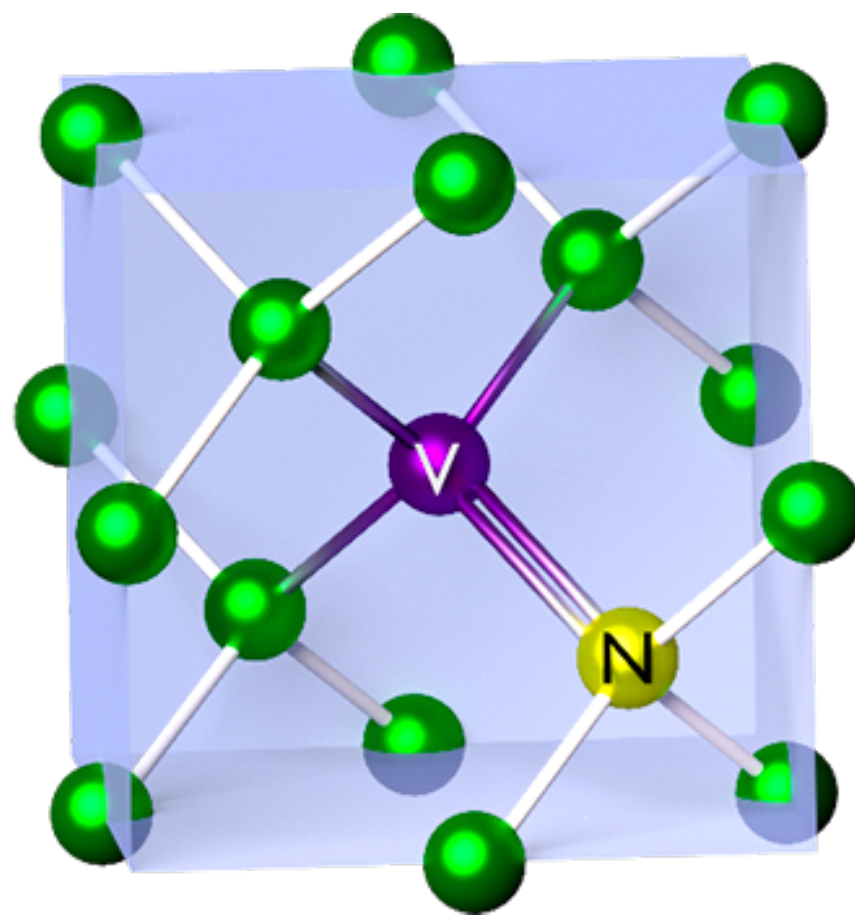
DDRF pulse sequence



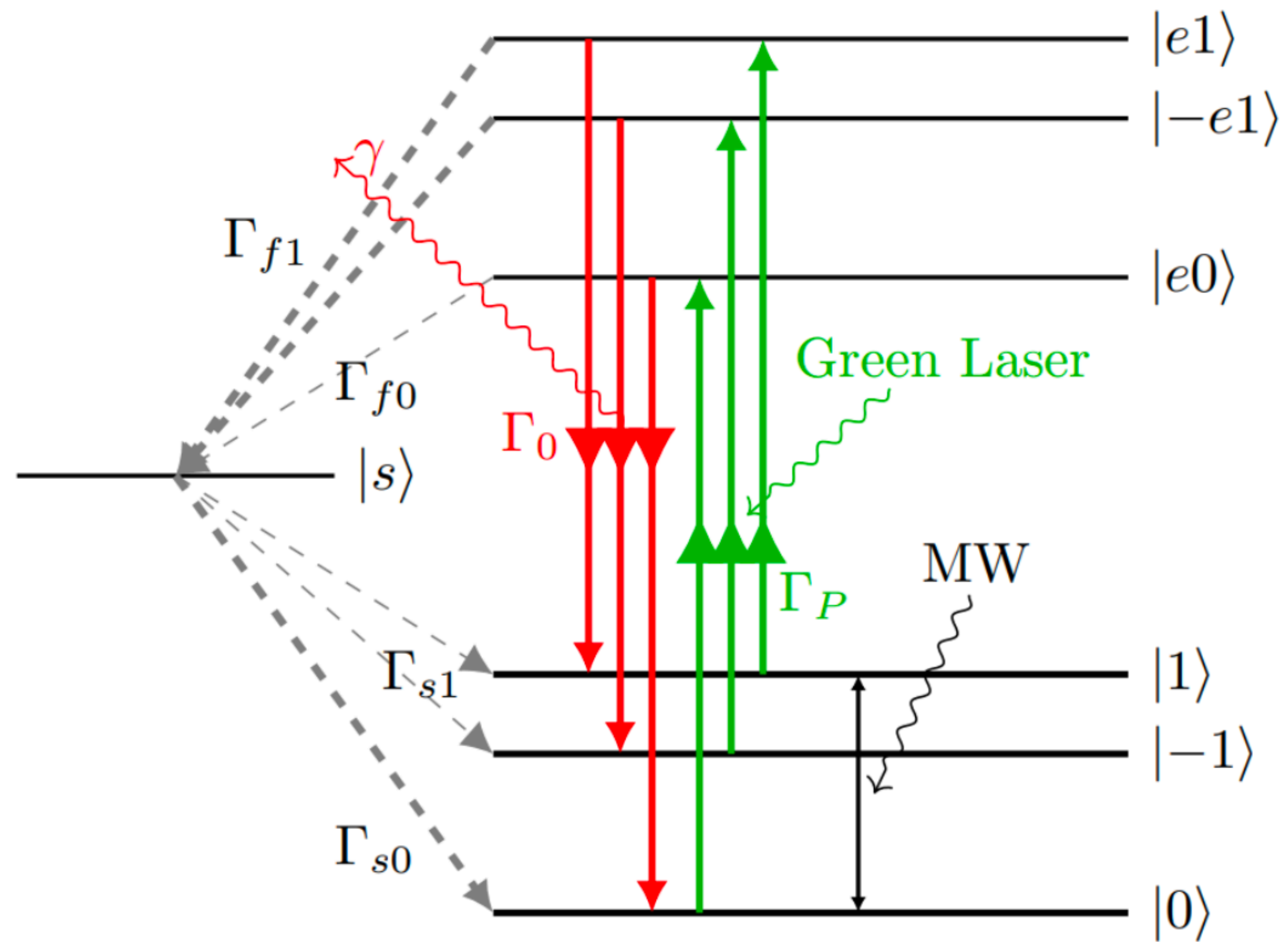
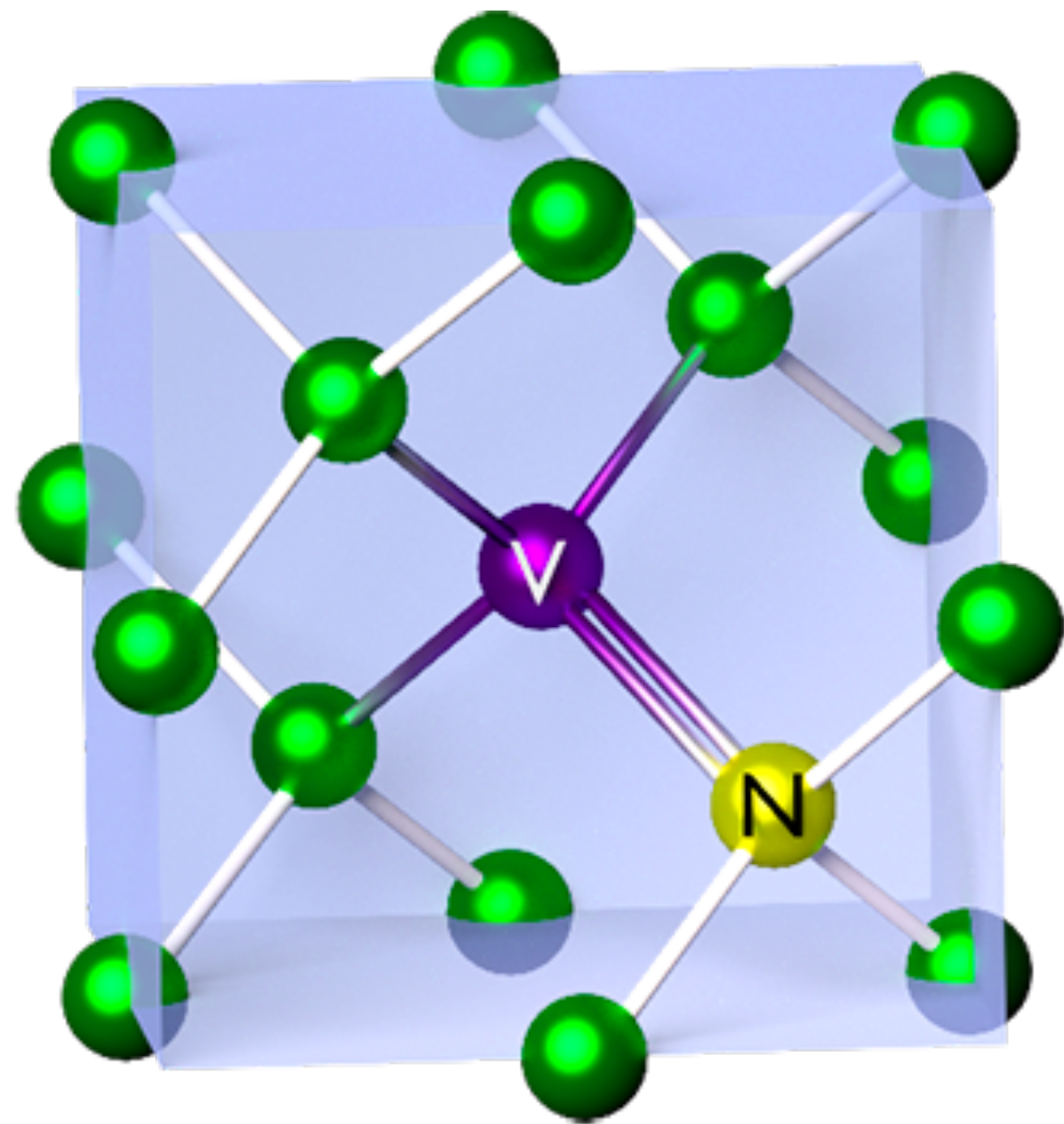
Towards an AI tool to enhance qubit readout in NV-based registers

How could this tool improve readout?

- Single-shot readout at elevated temperatures.
- In multi-shot readout (standard for room temperature): reduce the number of shots needed for a given precision.

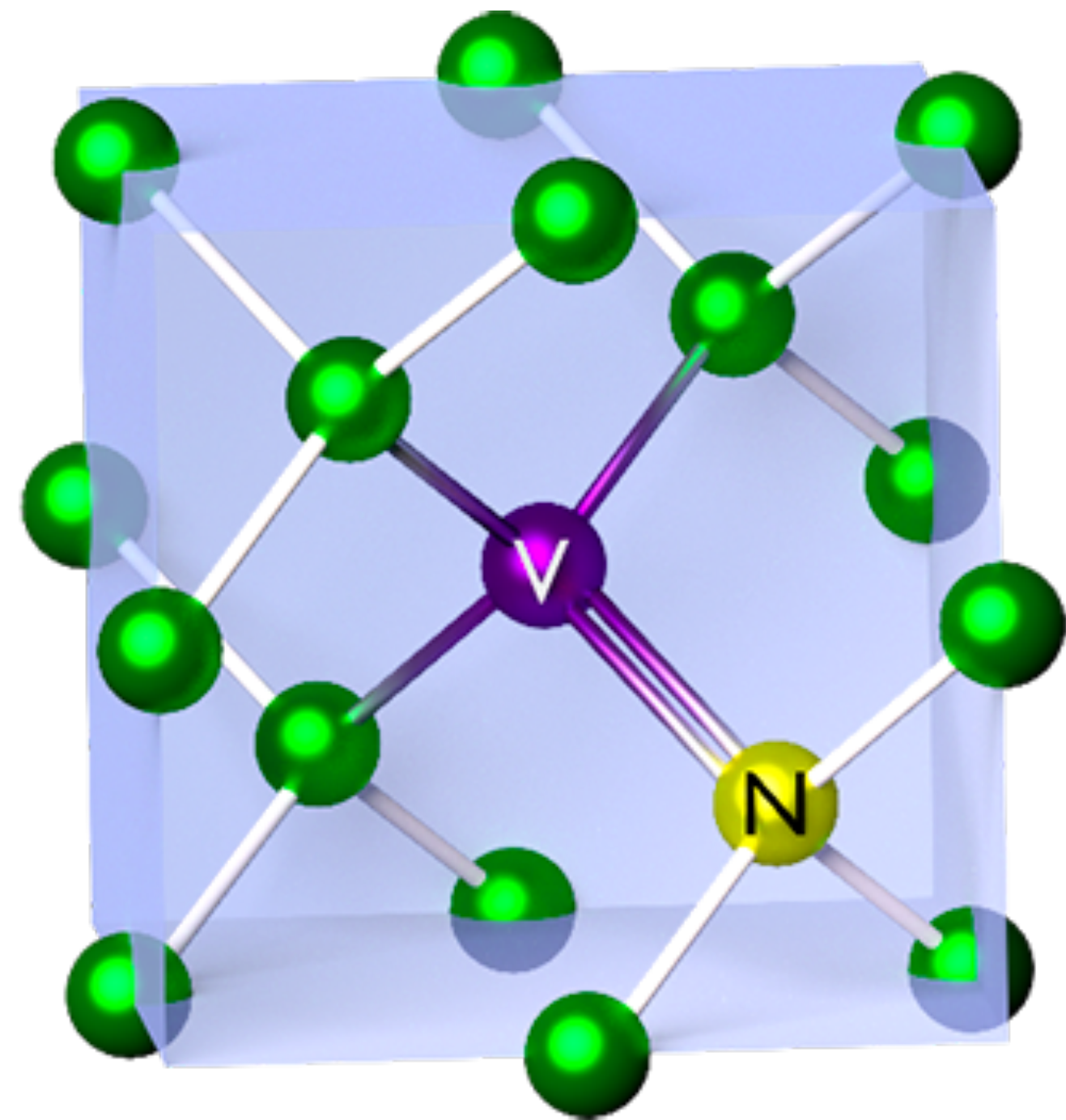


Towards an AI tool to enhance qubit readout in NV-based registers



ical values of the decay rates are $\Gamma_0 \simeq 63$ MHz, $\Gamma_{f0} \simeq 12$ MHz, $\Gamma_{f1} \simeq 80$ MHz, $\Gamma_{s1} \simeq 2.4$ MHz and $\Gamma_{s0} \simeq 3.3$ MHz

Minimal model to understand fundamental limitations of readout



green laser

red photoluminescent
photons

IDEAL PHOTODETECTOR
each PL photon reaches it
it detects all photons with certainty
it measures photon arrival times

Model: photon-number-resolved master equation

$$\dot{\rho}_0^n = -i\frac{\Omega}{2}(\rho_{01}^n - \rho_{10}^n) - \frac{\gamma_1}{2}(\rho_0^n - \rho_1^n) - \frac{\gamma_1}{2}(\rho_0^n - \rho_{-1}^n) - \Gamma_P\rho_0^n + \Gamma_0\rho_{e0}^{n-1} + \Gamma_{s0}\rho_s^n,$$

$$\dot{\rho}_1^n = i\frac{\Omega}{2}(\rho_{01}^n - \rho_{10}^n) + \frac{\gamma_1}{2}(\rho_0^n - \rho_1^n) - \Gamma_P\rho_1^n + \Gamma_0\rho_{e1}^{n-1} + \Gamma_{s1}\rho_s^n,$$

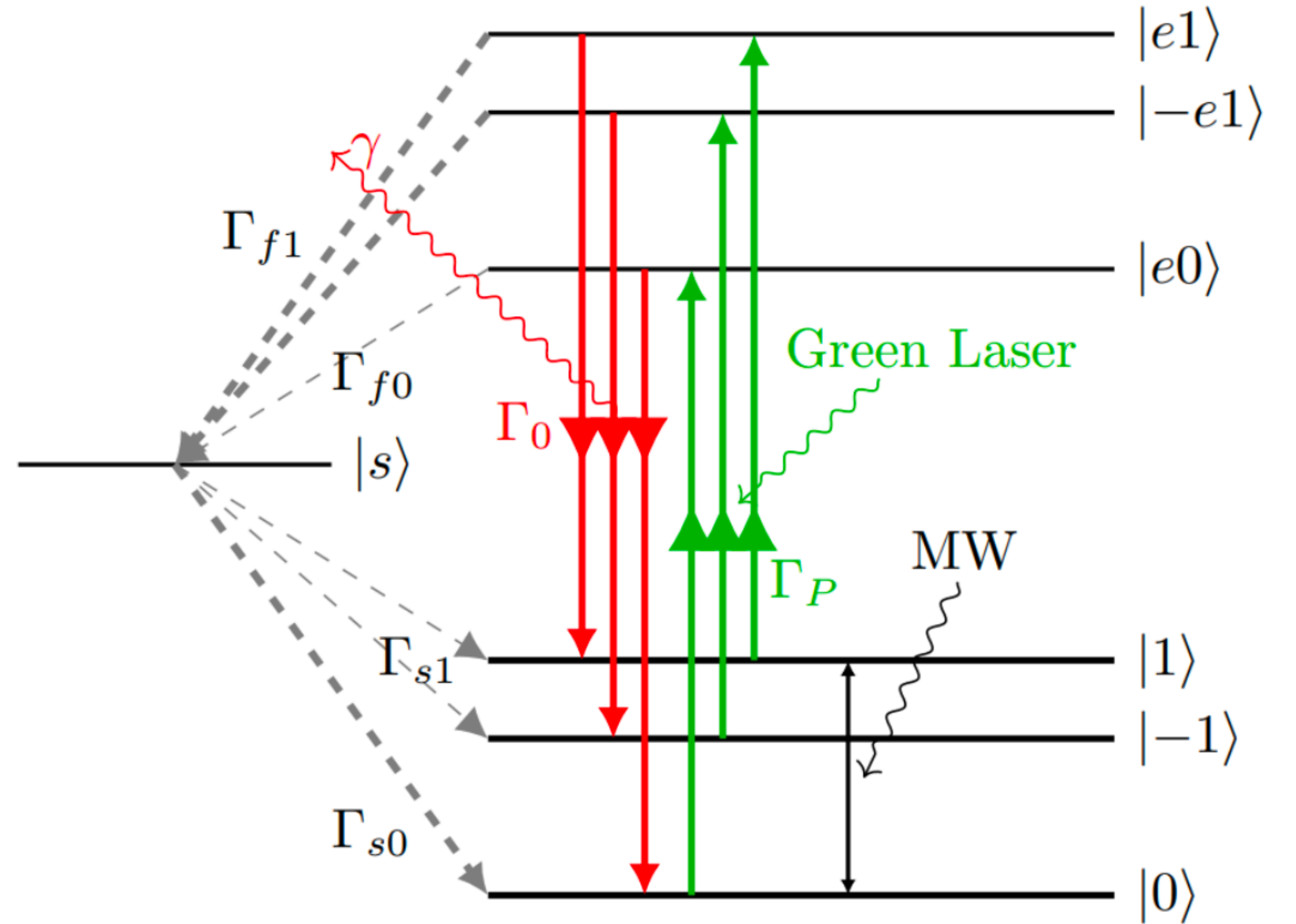
$$\dot{\rho}_{-1}^n = \frac{\gamma_1}{2}(\rho_0^n - \rho_{-1}^n) - \Gamma_P\rho_{-1}^n + \Gamma_0\rho_{-e1}^{n-1} + \Gamma_{s1}\rho_s^n,$$

$$\dot{\rho}_{e0}^n = \Gamma_P\rho_0^n - \Gamma_0\rho_{e0}^n - \Gamma_{f0}\rho_{e0}^n,$$

$$\dot{\rho}_{e1}^n = \Gamma_P\rho_1^n - \Gamma_0\rho_{e1}^n - \Gamma_{f1}\rho_{e1}^n,$$

$$\dot{\rho}_{-e1}^n = \Gamma_P\rho_{-1}^n - \Gamma_0\rho_{-e1}^n - \Gamma_{f1}\rho_{-e1}^n,$$

$$\dot{\rho}_s^n = \Gamma_{f1}\rho_{e1}^n + \Gamma_{f1}\rho_{-e1}^n + \Gamma_{f0}\rho_{e0}^n - (\Gamma_{s0} + \Gamma_{s1})\rho_s^n.$$

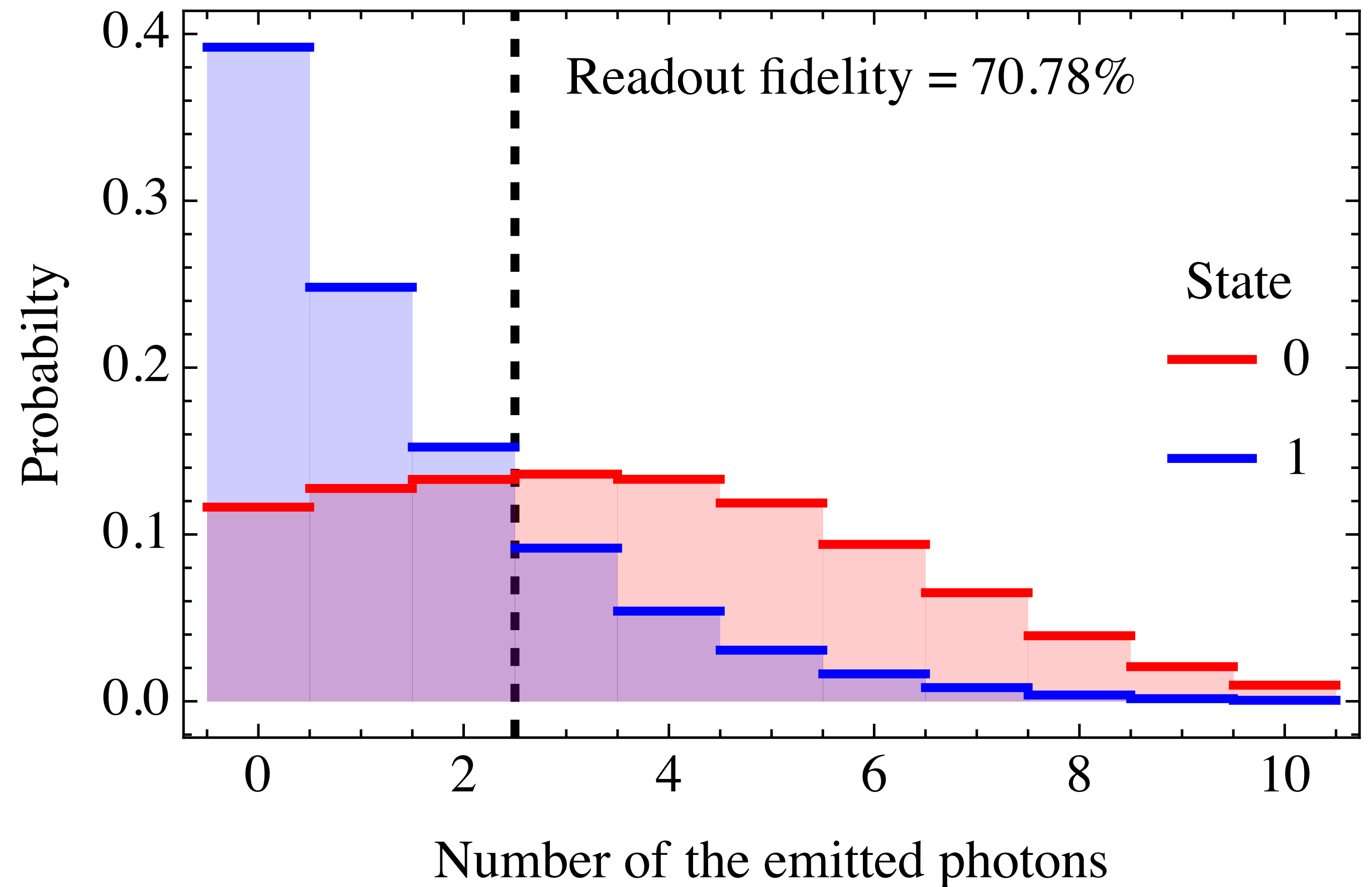


ical values of the decay rates are $\Gamma_0 \simeq 63$ MHz, $\Gamma_{f0} \simeq 12$ MHz, $\Gamma_{f1} \simeq 80$ MHz, $\Gamma_{s1} \simeq 2.4$ MHz and $\Gamma_{s0} \simeq 3.3$ MHz

Photon count reveals partial information about qubit state

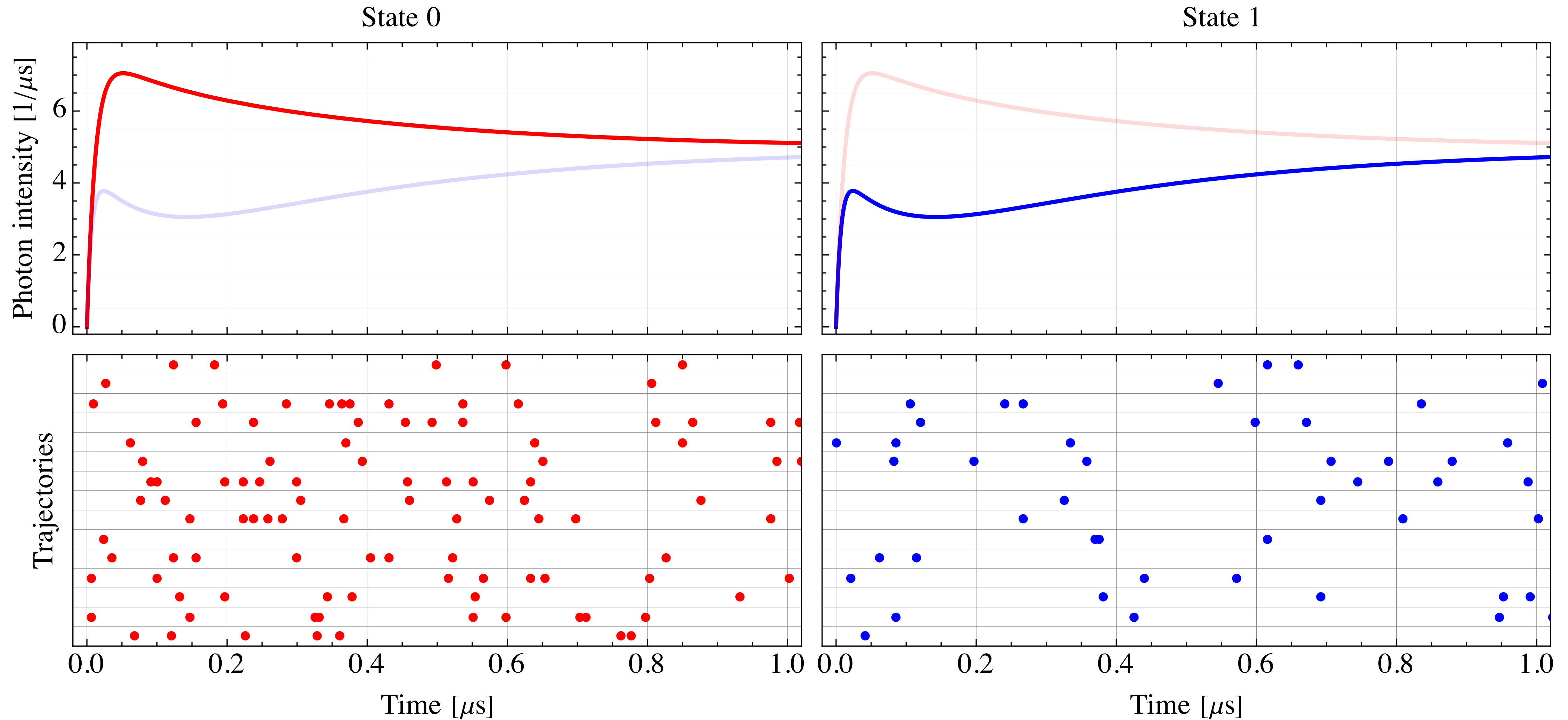
- (1) fixed laser intensity
- (2) fixed 'integration time'

red/blue:
photon-count histograms
for 0/1 qubit basis states



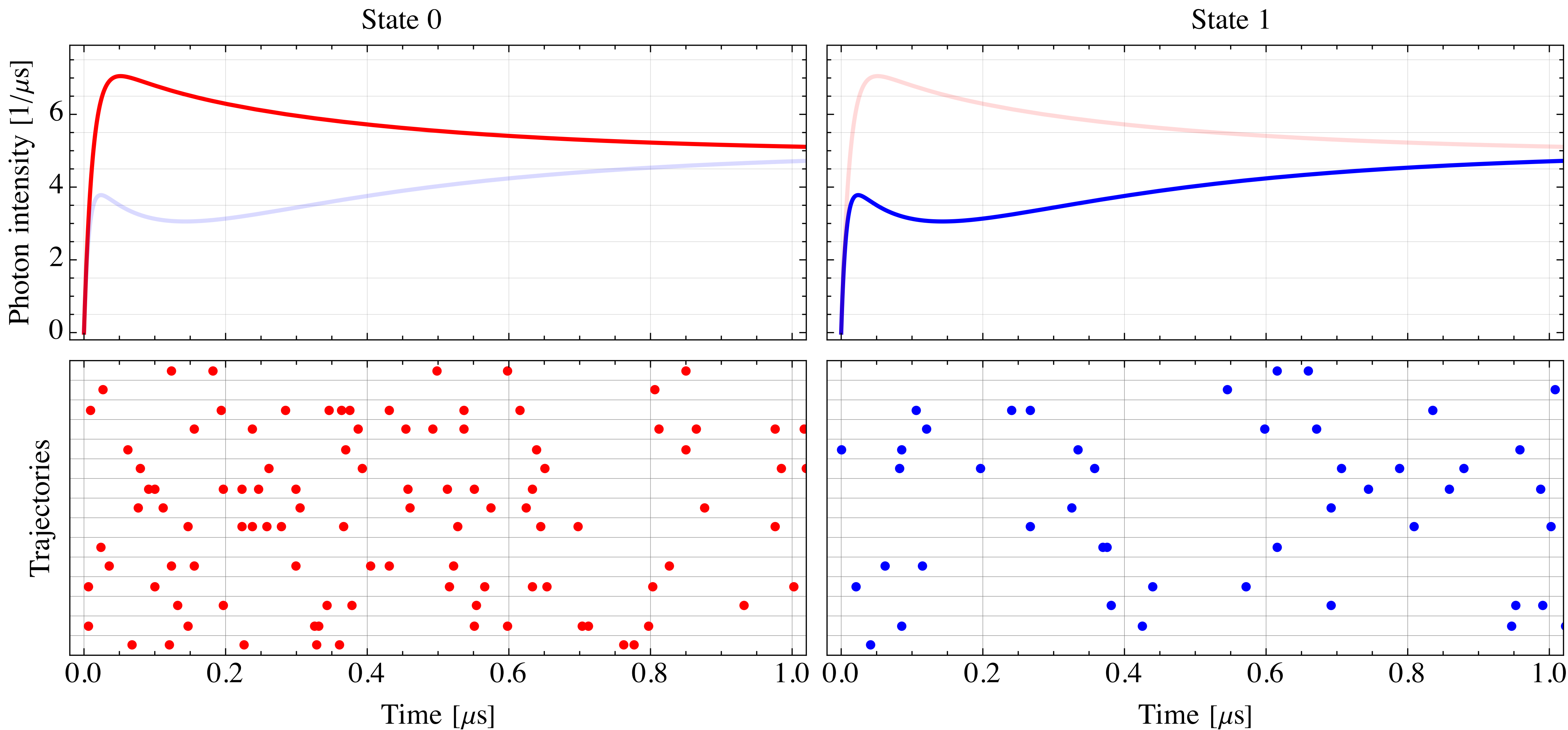
Laser intensity = 10 x PL decay rate.
Integration time = 0.1 us

Photon detection trajectories from the master equation



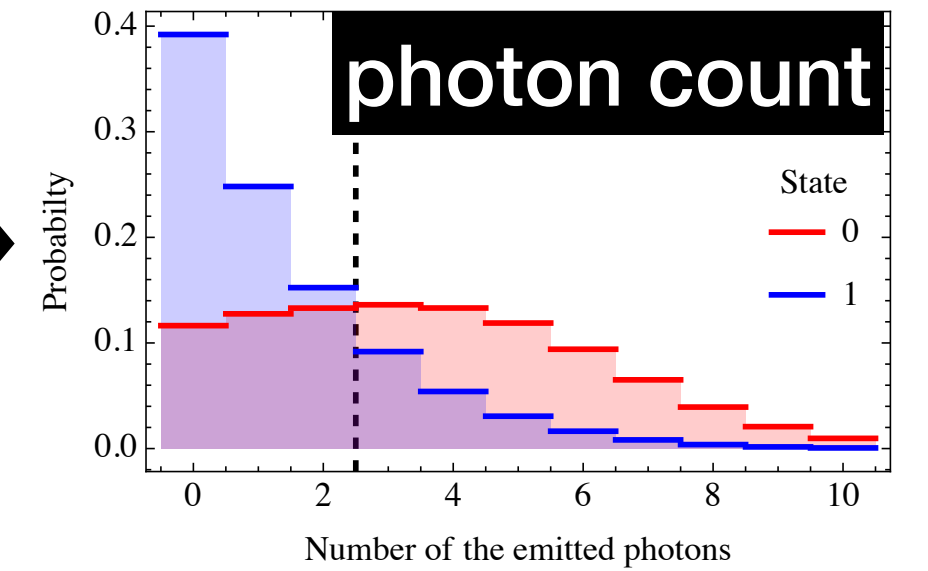
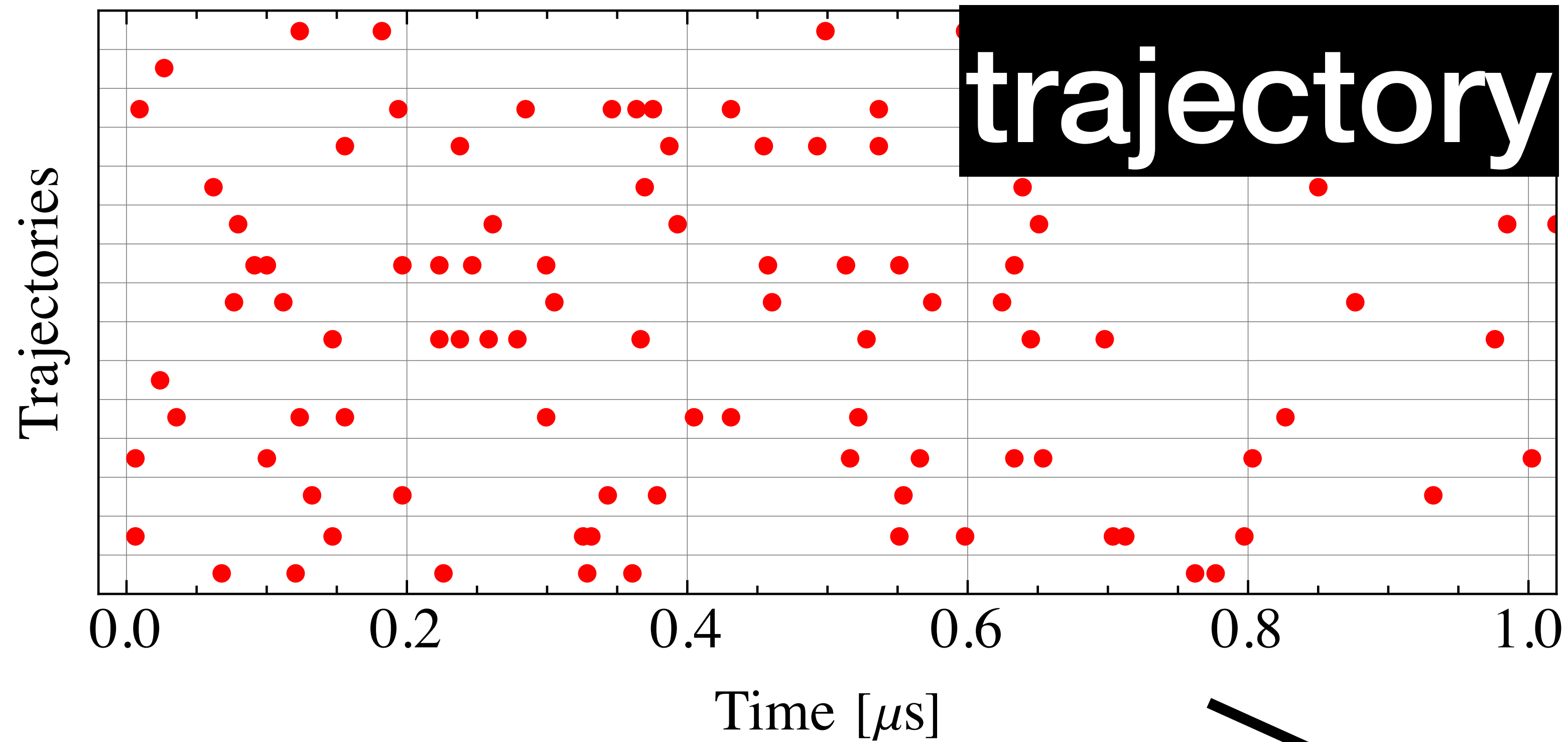
Laser intensity = 0.16 x PL decay rate.

There is more information in the trajectory than in the photon count!



Laser intensity = 0.16 x PL decay rate.

There is more information in the trajectory than in the photon count!



We try to squeeze out a more precise qubit readout from the more complete information present in the trajectories.

inferred bit, 0 or 1

Photon count trajectories from master equation

$$\begin{aligned}\dot{\rho}_0^n &= -i\frac{\Omega}{2}(\rho_{01}^n - \rho_{10}^n) - \frac{\gamma_1}{2}(\rho_0^n - \rho_1^n) - \frac{\gamma_1}{2}(\rho_0^n - \rho_{-1}^n) \\ &\quad - \Gamma_P \rho_0^n + \Gamma_0 \rho_{e0}^{n-1} + \Gamma_{s0} \rho_s^n,\end{aligned}$$

$$\begin{aligned}\dot{\rho}_1^n &= i\frac{\Omega}{2}(\rho_{01}^n - \rho_{10}^n) + \frac{\gamma_1}{2}(\rho_0^n - \rho_1^n) \\ &\quad - \Gamma_P \rho_1^n + \Gamma_0 \rho_{e1}^{n-1} + \Gamma_{s1} \rho_s^n,\end{aligned}$$

$$\dot{\rho}_{-1}^n = \frac{\gamma_1}{2}(\rho_0^n - \rho_{-1}^n) - \Gamma_P \rho_{-1}^n + \Gamma_0 \rho_{-e1}^{n-1} + \Gamma_{s1} \rho_s^n,$$

- (1) Given transition rates $W_{i \leftarrow j}$ and initial time $t = 0$
- (2) Given state j
- (3) Random realization of the jump times $\tau_i \sim \text{Exp}(1/W_{i \leftarrow j})$ ($\forall i$)
- (4) $j \leftarrow \text{argmin}_i \tau_i$ is the new state and $t \leftarrow t + \min \tau_i$
- (5) if $t > t_{\max}$ break else goto (2)

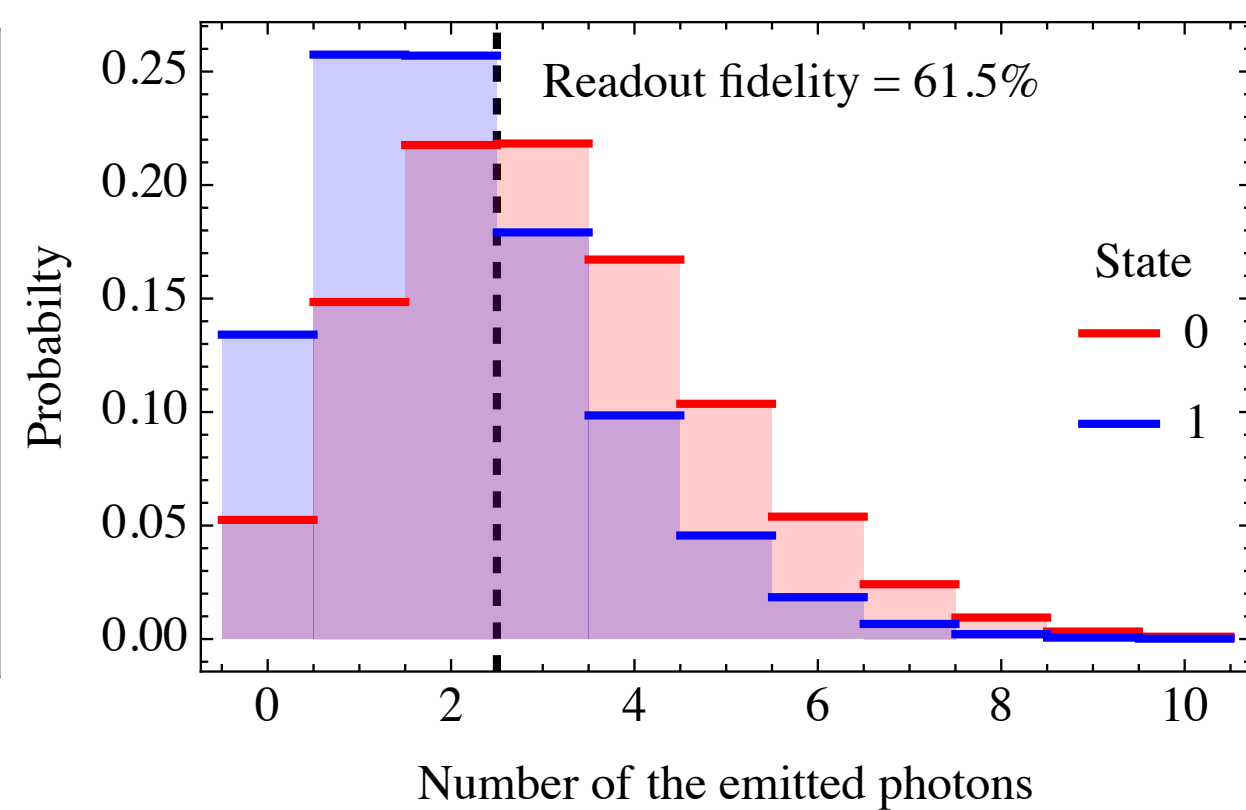
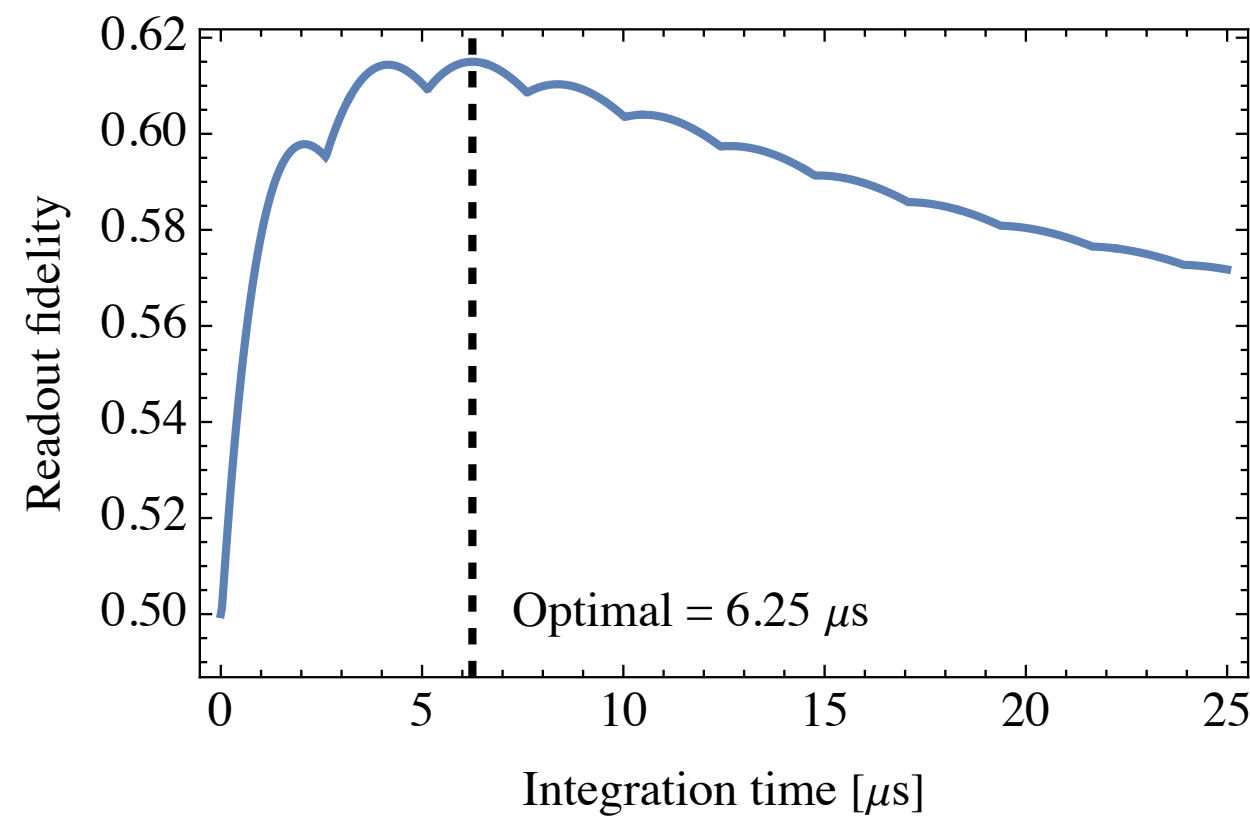
$$\dot{\rho}_{e0}^n = \Gamma_P \rho_0^n - \Gamma_0 \rho_{e0}^n - \Gamma_{f0} \rho_{e0}^n,$$

$$\dot{\rho}_{e1}^n = \Gamma_P \rho_1^n - \Gamma_0 \rho_{e1}^n - \Gamma_{f1} \rho_{e1}^n,$$

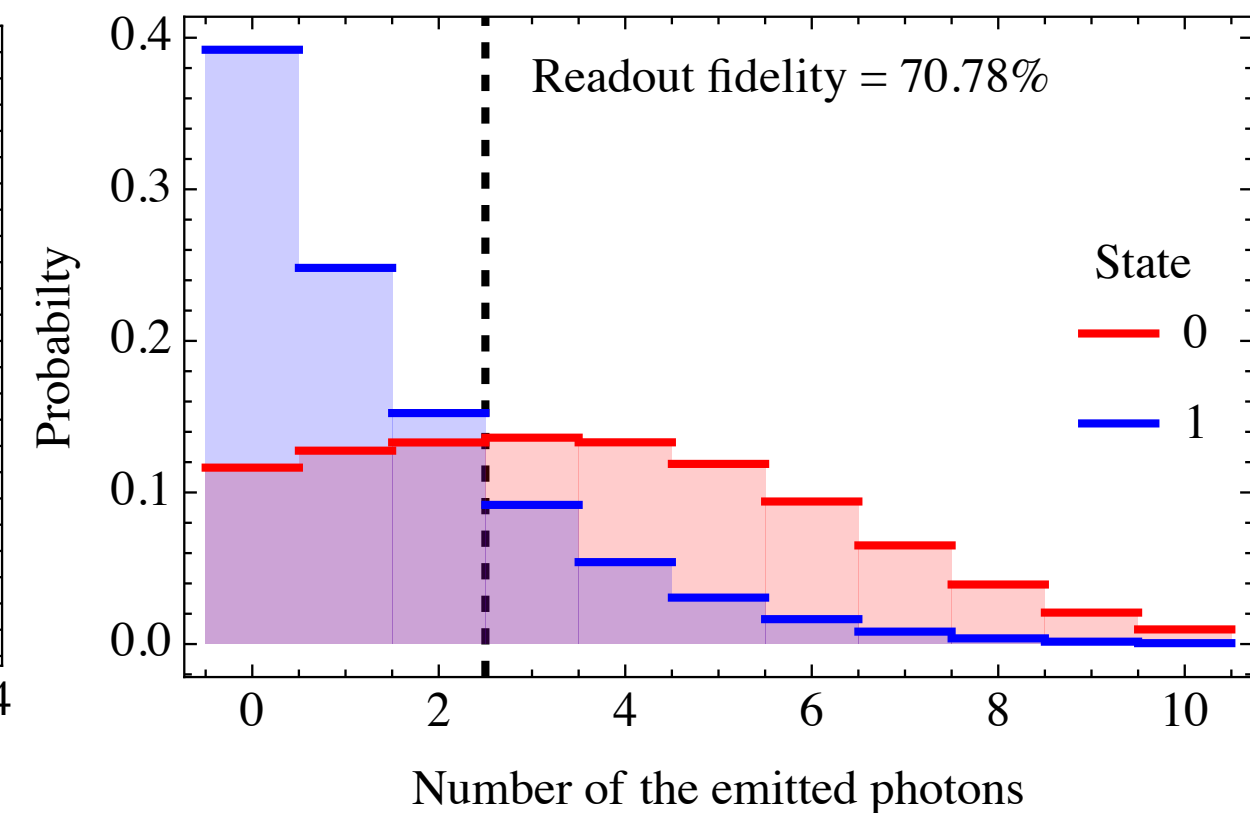
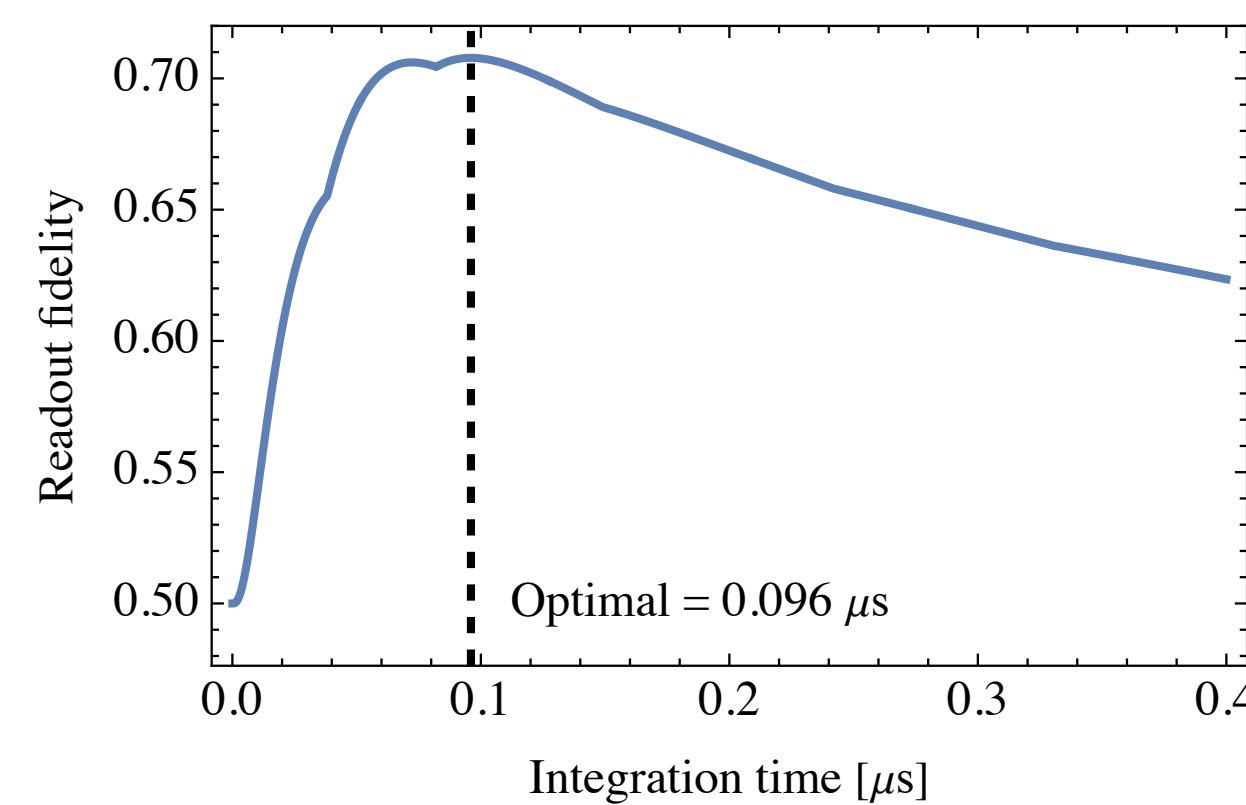
$$\dot{\rho}_{-e1}^n = \Gamma_P \rho_{-1}^n - \Gamma_0 \rho_{-e1}^n - \Gamma_{f1} \rho_{-e1}^n,$$

$$\dot{\rho}_s^n = \Gamma_{f1} \rho_{e1}^n + \Gamma_{f1} \rho_{-e1}^n + \Gamma_{f0} \rho_{e0}^n - (\Gamma_{s0} + \Gamma_{s1}) \rho_s^n.$$

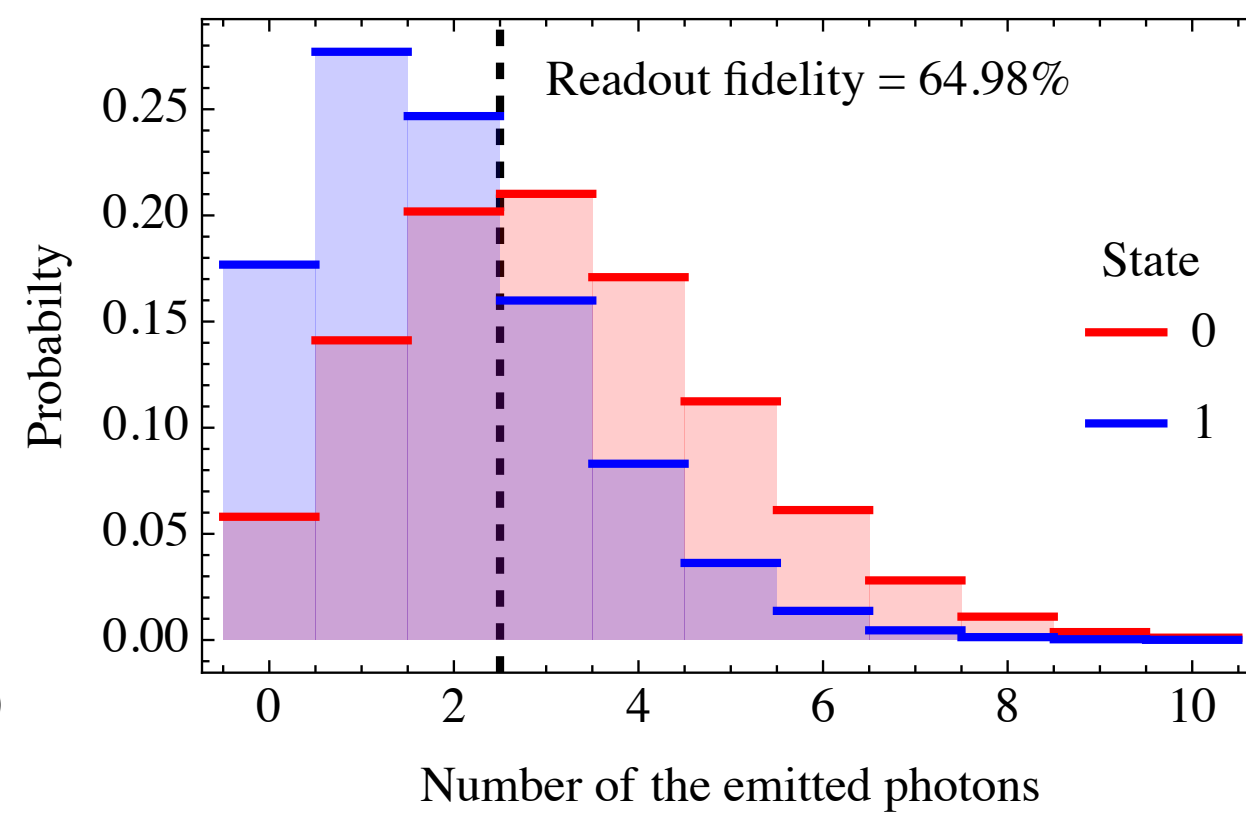
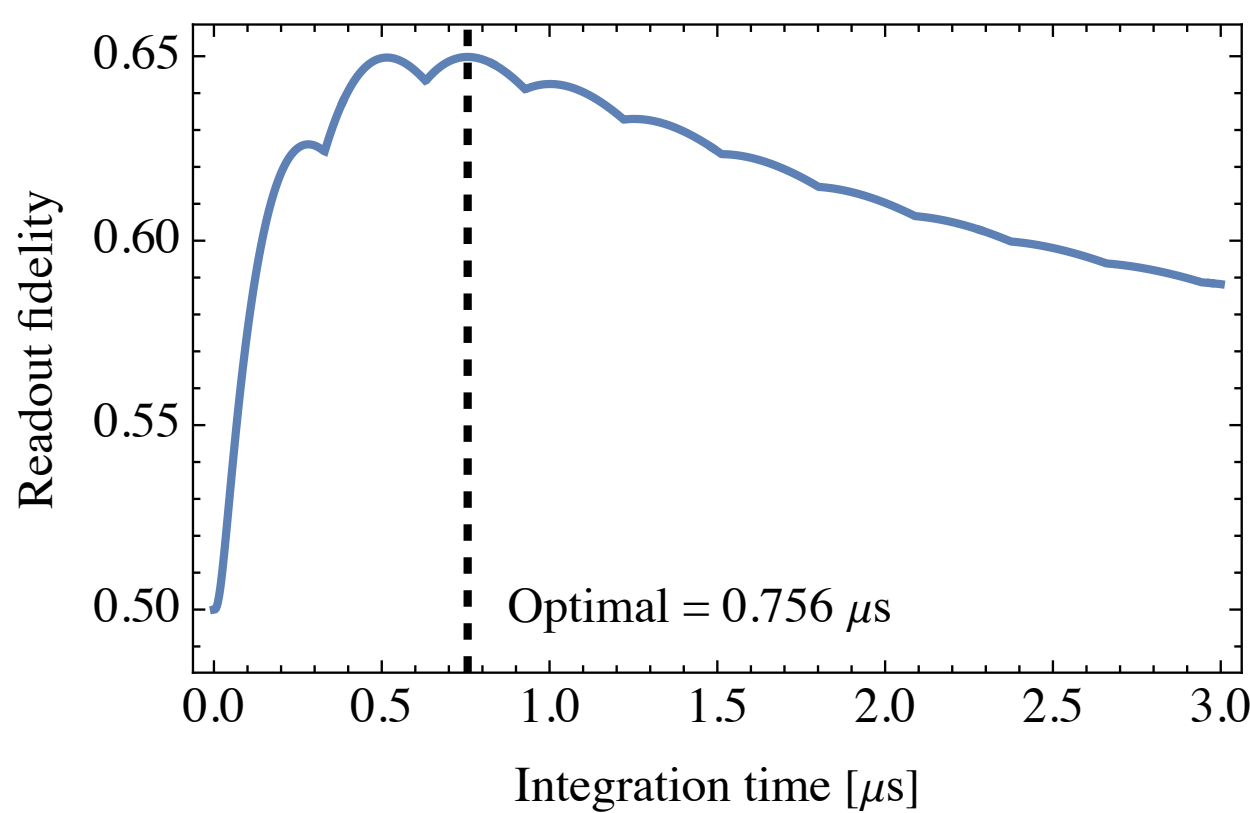
Laser intensity (excitation rate) = 0.01 (in the unit of direct decay rate from excited state to ground state)



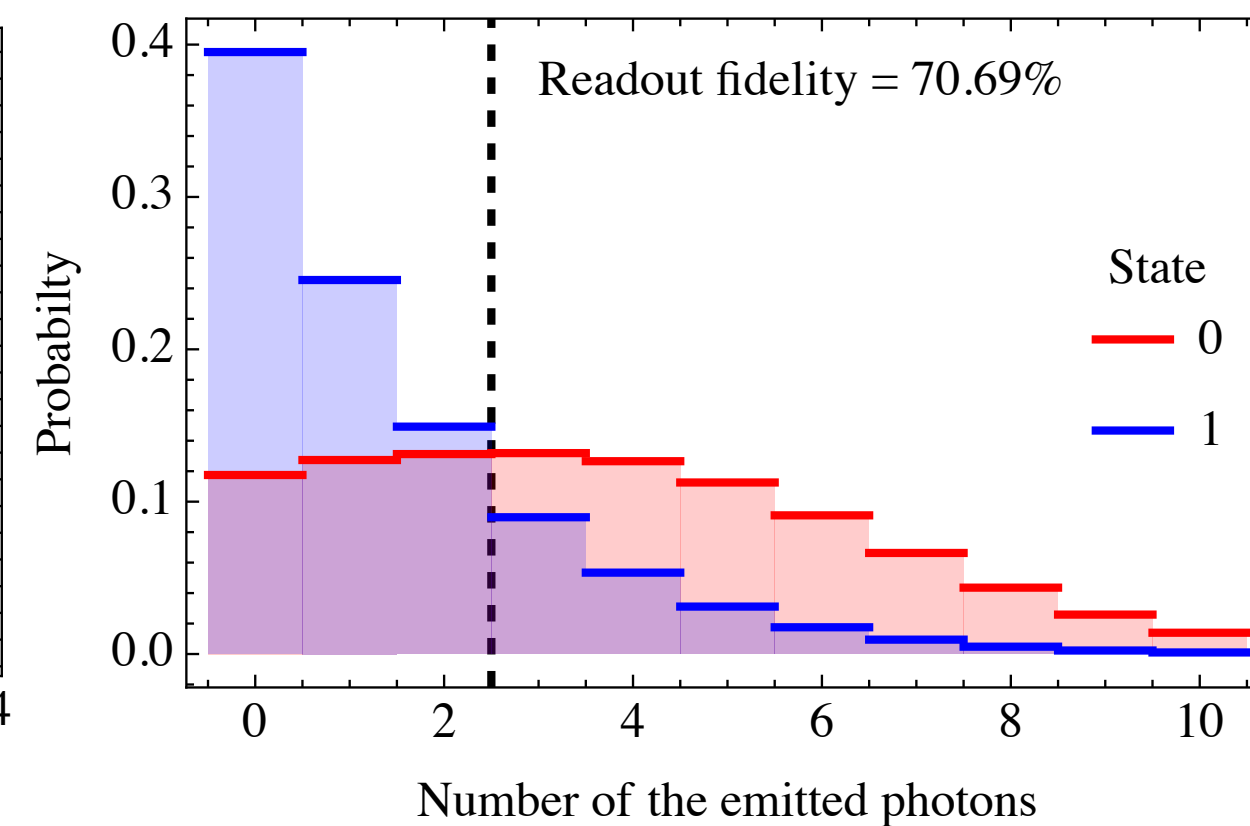
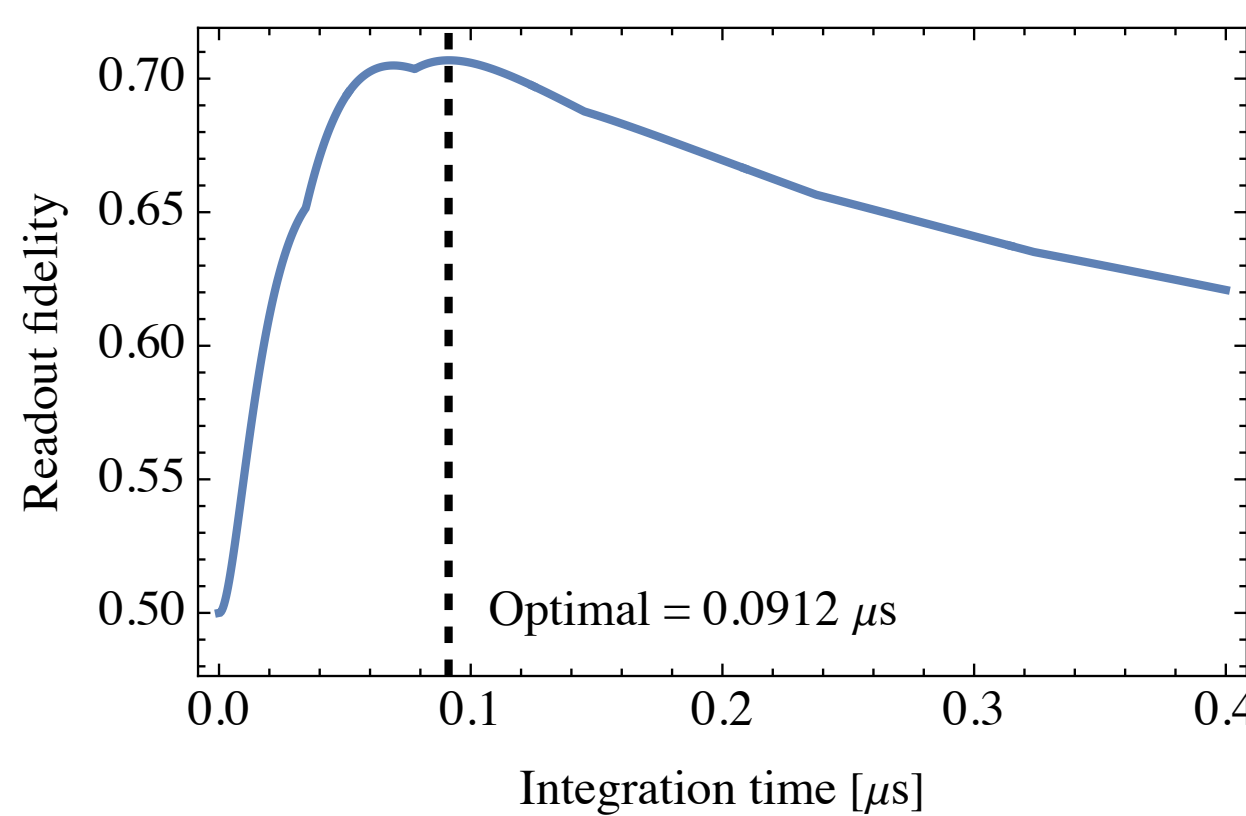
Laser intensity (excitation rate) = 10 (in the unit of direct decay rate from excited state to ground state)



Laser intensity (excitation rate) = 0.1 (in the unit of direct decay rate from excited state to ground state)



Laser intensity (excitation rate) = 100 (in the unit of direct decay rate from excited state to ground state)



Laser intensity (excitation rate) = 1 (in the unit of direct decay rate from excited state to ground state)

