

Hadronic Structure from Functional Methods

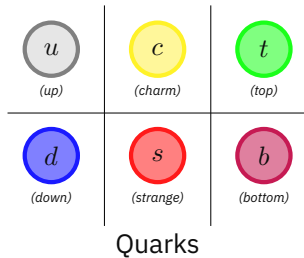
Eduardo Ferreira

ACHT2025, Budapest – May 5-7, 2025

1 **Functional Methods**

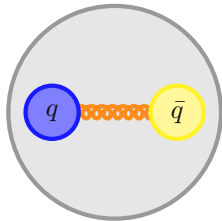
2 Hadronic Structure

Hadron Physics



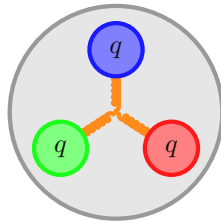
Gluons

Hadron Physics



Mesons

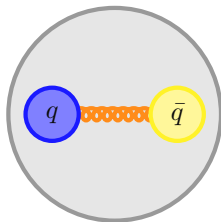
$\pi, \rho, K, \eta, \dots$



Baryons

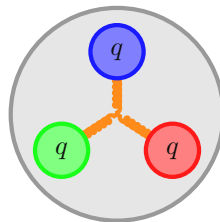
$p, n, \Delta, \Sigma, \Lambda, \dots$

Hadron Physics



Mesons

$\pi, \rho, K, \eta, \dots$



Baryons

$p, n, \Delta, \Sigma, \Lambda, \dots$

- Their dynamics and properties are determined by QCD
 - Atomic nuclei and nuclear stability, EMC effect
 - Radius; spin-, momentum-, charge distributions;
 - Interaction with external currents: $\{e^-, \nu, \dots\}N$ scattering, $N\pi$ scattering, nucleon compton scattering, meson photoproduction

(EIC White Paper), (Eichmann, Sanchis-Alepuz, Williams, Alkofer, Fischer; 2016), (European Muon Collaboration; 1983), ...

Dyson-Schwinger Equations

- Can relate n -point functions with each other:

The diagram illustrates the Dyson-Schwinger equation for the quark propagator. It shows a horizontal line with a circle in the middle, labeled with a superscript -1 , followed by an equals sign. To the right of the equals sign is a horizontal line with a superscript -1 , followed by a plus sign. To the right of the plus sign is a diagram representing a self-energy loop: a horizontal line with a circle, followed by a loop consisting of two circles connected by two wavy lines, and then a shaded circle at the end of the horizontal line.

- The **full quark propagator** is the bare propagator plus self-energy Σ
- Σ depends on gluon and qqg vertex

Dyson-Schwinger Equations

- Can relate n -point functions with each other:

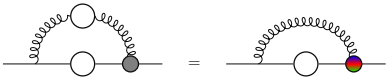
$$\begin{aligned}
 & \text{---} \bigcirc \text{---}^{-1} = \text{---}^{-1} + \text{---} \bigcirc \text{---}^{-1} \\
 & \text{~~~~~} \text{---} \bigcirc \text{---}^{-1} = \text{~~~~~} \text{~~~~~}^{-1} + \text{~~~~~} \bigcirc \text{~~~~~}^{-1} + \text{~~~~~} \bigcirc \text{~~~~~} \bigcirc \text{~~~~~}^{-1} \\
 & \quad + \text{~~~~~} \bigcirc \text{~~~~~} \bigcirc \text{~~~~~}^{-1} + \text{~~~~~} \bigcirc \text{~~~~~} \bigcirc \text{~~~~~} \bigcirc \text{~~~~~}^{-1} + \text{~~~~~} \bigcirc \text{~~~~~} \bigcirc \text{~~~~~} \bigcirc \text{~~~~~} \bigcirc \text{~~~~~}^{-1}
 \end{aligned}$$

The diagrams represent the Dyson-Schwinger equations for the fermion and photon propagators. In the fermion diagrams, white circles represent fermion loops and black dots represent photon insertion points. In the photon diagrams, white circles represent photon loops and black dots represent fermion insertion points. The curly lines represent photon propagators, and the straight lines represent fermion propagators.

- Coupled equations: n -point function depends on higher order correlations.

Truncations

- Calculations only possible with truncations
 - Chosen to reproduce the interesting physics of the system.
 - Common truncation: Maris-Tandy model (Maris, Roberts; 1997), (Maris, Tandy; 1999), (Maris, Tandy; 2000)



$$\frac{\alpha(k^2)}{k^2} = \pi\eta^7 \left(\frac{k^2}{\Lambda^2}\right)^2 e^{-\eta^2 \frac{k^2}{\Lambda^2}} + \frac{2\pi\gamma_m \left(1 - e^{-\frac{k^2}{\Lambda_t^2}}\right)}{\ln \left[e^2 - 1 + \left(1 + \frac{k^2}{\Lambda_{QCD}^2}\right)^2 \right]}$$

$$\text{gluon line with dot}^{-1} = \text{gluon line}^{-1} + C_F \text{gluon line with dot and loop}$$

$$\text{gluon line with dot}^{-1} = \text{gluon line}^{-1} + C_A \text{gluon line with dot and loop}$$

$$\text{gluon line with dot}^{-1} = \text{gluon line}^{-1} + TN_f \text{gluon line with dot and loop}$$

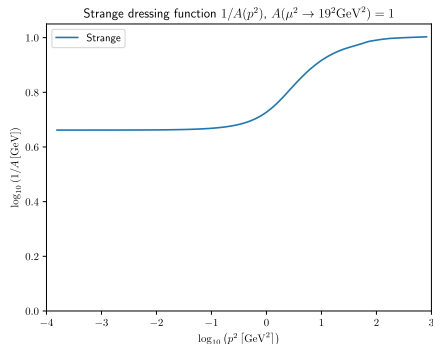
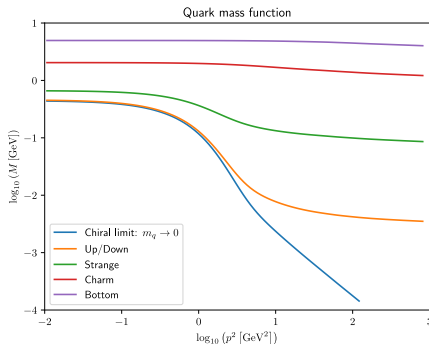
$$+ C_A \left(\text{gluon line with dot and loop} + \text{gluon line with dot and loop} \right)$$

$$\text{gluon line with dot and loop} = \text{gluon line with dot and loop} + \frac{1}{2} C_A^2 \left(\text{gluon line with dot and loop} + \text{gluon line with dot and loop} \right)$$

(Alkofer, Zierler; 2023)

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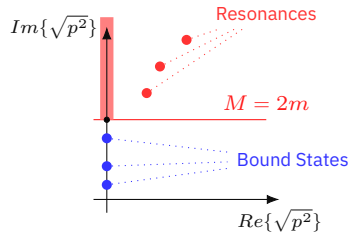
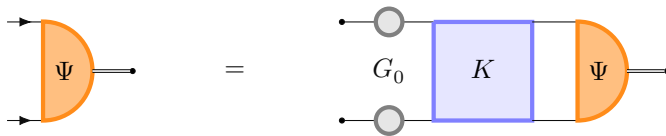
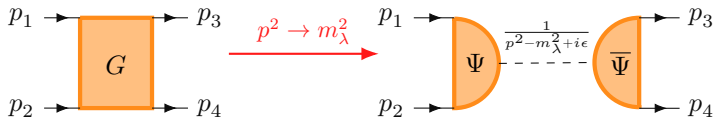


$$S(p) = \sigma_s(p^2) \mathbb{1} - i \sigma_v(p^2) \not{p} = \frac{1}{A(p^2)} \frac{-i \not{p} + M(p^2)}{p^2 + M(p^2)^2}$$

(Numerical results by: Raúl)

Bethe-Salpeter Equation

- The Bethe-Salpeter Wavefunction (BSWF) appears as the residue of a correlation function $G(p)$:



- Poles in correlation functions encode the theory's bound-state spectrum
- BSWF characterizes the bound-state

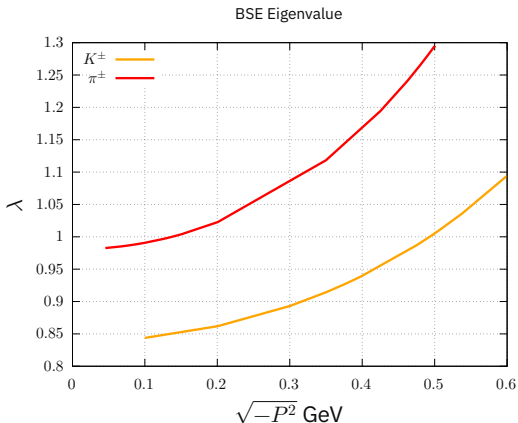
(Eichmann, Sanchis-Alepuz, Williams, Alkofer, Fischer; 2016)

Bethe-Salpeter Equation *Eigenvalue/vector equation*

- The BSE is usually solved as an eigenvalue equation:

$$\lambda_i(P^2) \psi_i = \mathbf{K} \mathbf{G}_0 \psi_i$$

- States found when $\lambda_i(P^2) \equiv 1$.
- Mass is given by $P^2 = -M^2$
- Eigenvalue/vector spectrum gives ground state and excited states.
- Like DSE, analytic structure is important, and **contour deformations** may be needed!

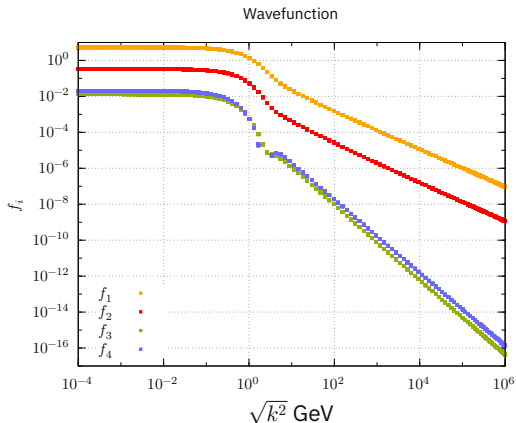


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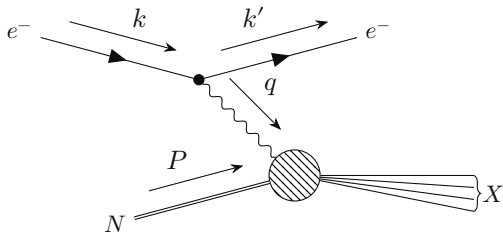
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1 Functional Methods

2 **Hadronic Structure**

Parton model



(PDG Section 18 (Structure Functions))

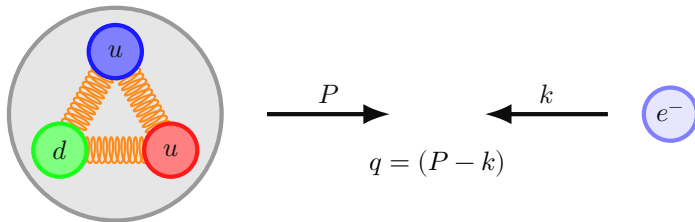
- Common experiment:
Deep inelastic scattering.
- Electron (or neutrino) probe interacts with nucleon via $\gamma, Z, W^\pm$.
- Main idea:
 - Hadrons are *bags* of **partons**.
 - Scattering occurs between the exchanged boson and a parton from the nucleon
- Cross-section separates:

$$\frac{d^2\sigma}{dx dq^2} \propto L_{\mu\nu} W^{\mu\nu}$$

- $W^{\mu\nu}$ described via PDFs.

PDFs/TMDs

- How to they show up?



$$W_{\mu\nu}(p, q, \sigma) \propto \int d^4z \langle p, \sigma | J_\mu^\dagger(z) J_\nu(0) | p, \sigma \rangle = \sum_i \tau(p, q, \sigma)_{\mu\nu}^i f_i(p, q)$$

PDFs/TMDs

- How to they show up?
- Take the collinear, q^2 very large limit.

$$\left| \text{Diagram} \right|^2 = \text{Diagram} \otimes \text{Diagram} + \mathcal{O}\left(\frac{k_{\perp}^2}{q^2}\right)$$

$$\sigma \propto \sum_f \int \frac{dx}{x} H(x, Q^2/\Lambda^2) q_f(x, \Lambda^2) + \mathcal{O}\left(\frac{k_{\perp}^2}{Q^2}\right)$$

$$q_f(x, \Lambda) \propto \int \frac{d^4\ell}{(2\pi)^4} \delta\left(x - \frac{\ell^+}{p^+}\right) \text{Tr} \left[\langle p | \bar{\psi}_f(\ell) \gamma^+ \psi_f(\ell) | p \rangle \right]$$

- Hard scale $Q^2 = -q^2$ introduced allows for factorization – power suppression $\mathcal{O}(Q^{-2})$
- Works also on other processes, like Drell-Yan.

PDFs/TMDs

- How to they show up?
- Take the collinear, q^2 very large limit.

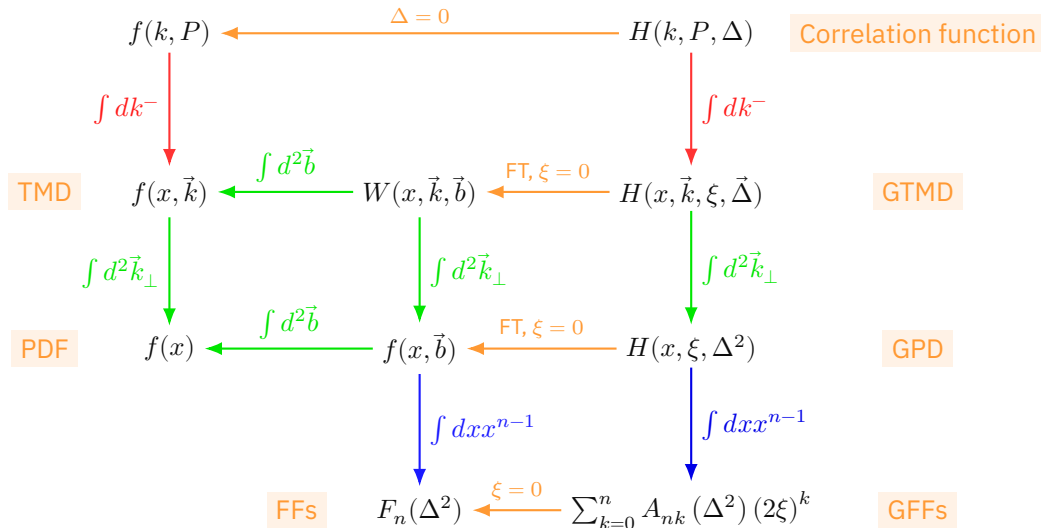
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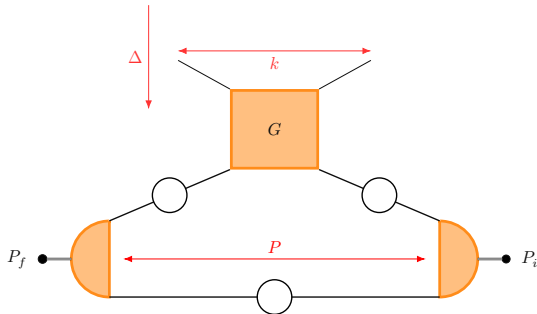
- Focus on the non-perturbative hadronic part.
- Can we build it from first-principles **FUN**ctional Methods calculations?

Hadronic quantities



Our Goal

- **Main Goal:** Get partonic distribution functions from hadron-hadron correlations via **FUN**ctional Methods



- G is the four-point quark correlation function, calculated with scattering equation.
- The BSWF is calculated via the meson BSE.
- First scalar toy model, then QCD.

(Mezrag; 2015), (Diehl, Gousset; 1998), (Tiburzi, Miller; 2003),
(Mezrag, Chang, Moutarde, Roberts, Rodríguez-Quintero, Sabatié, Schmidt; 2015),
many others, ...

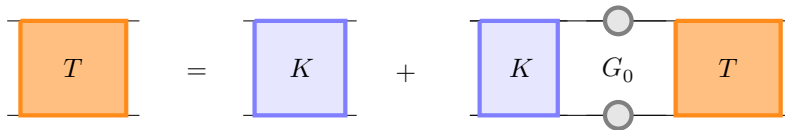
$$\mathcal{G}^{[\Gamma]}(P, k, \Delta) = \frac{1}{2} \text{Tr} \left[\int dk^- \int \frac{d^4 z}{2\pi^4} e^{ik \cdot z} \langle P_f | \bar{\psi}(z) \mathcal{W} \Gamma \psi(0) | P_i \rangle \right]$$

- Partonic distributions are calculated by integrating the correlator in k^- and taking appropriate traces.

4-point function

- 4-point function determined from scattering equation:

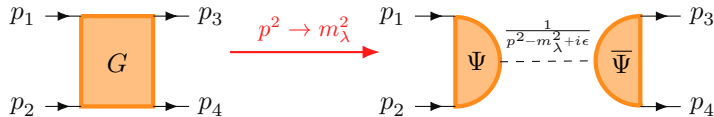
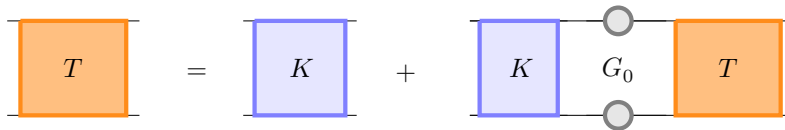
$$\mathbf{G} = \mathbf{G}_0 + \mathbf{G}_0 \mathbf{T} \mathbf{G}_0 \implies \mathbf{T} = \mathbf{K} + \mathbf{K} \mathbf{G}_0 \mathbf{T} \implies \mathbf{T} = (\mathbb{1} - \mathbf{K} \mathbf{G}_0)^{-1} \mathbf{K}.$$



4-point function

- 4-point function determined from scattering equation:

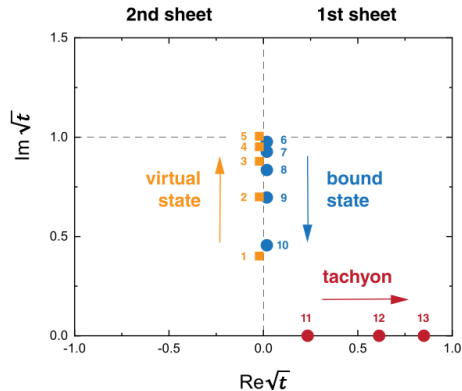
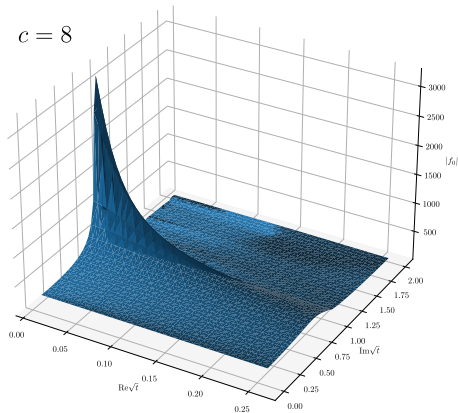
$$G = G_0 + G_0 T G_0 \Rightarrow T = K + K G_0 T \Rightarrow T = (\mathbb{1} - K G_0)^{-1} K.$$



- Same G_0 and K as in the BSE
- All dynamics of 2 ϕ particles:
 - Must produce bound state poles dynamically!**

4-point function

- Partial-wave expansion shows bound-state pole in the first Riemann sheet

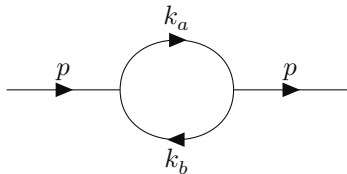


(Eichmann, Duarte, Peña, Stadler; 2019)

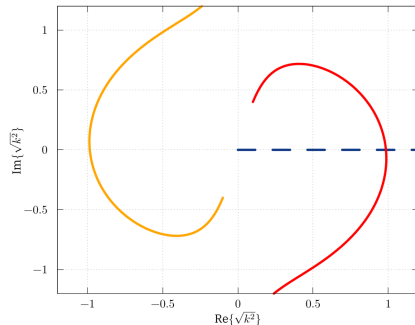
- Describes both long-range and short-range qq dynamics.

Analytic Structure is Important

- Calculating quantities in time-like momenta is **complicated** !
 - Analytic structure prevents naïve integration.
 - Poles and branch cuts are present.
- Calculate quantities in the P^2 complex plane.
- **Euclidean** $\Leftrightarrow$ Minkowski



$$I(p^2) = \int d^4k \frac{1}{k_a^2 - m^2 + i\epsilon} \frac{1}{k_b^2 - m^2 + i\epsilon}$$

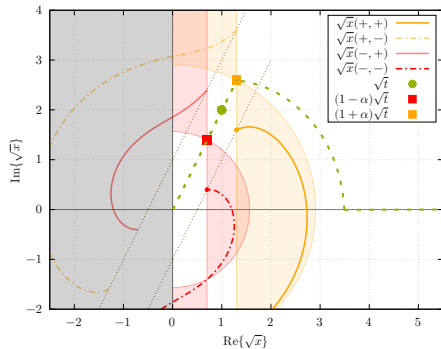


Analytic Structure of BSE $\text{Im} \sqrt{t} > \min \left(\frac{1}{1 \pm \alpha} \right)$

$$\mathbf{G}_0^{-1} = (l_+^2 + m^2)(l_-^2 + m^2)$$

- Branch cuts in complex $\sqrt{l}$ plane

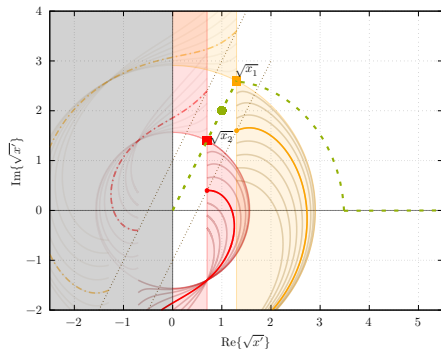
$$\sqrt{l}_\pm^\lambda = \mp(1 \pm \alpha)\sqrt{t} \left[z_l + i\lambda \sqrt{1 - z_l^2 + \frac{1}{(1 \pm \alpha)^2 t}} \right]$$



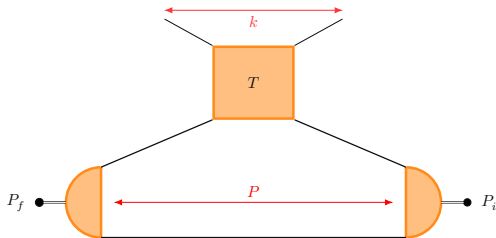
$$\mathbf{K}^{-1} = (q - l)^2 + \mu^2$$

- Branch cuts in complex $\sqrt{l}$ plane depend on path taken:

$$\sqrt{l} = \sqrt{\rho} \left(\Omega \pm i \sqrt{1 - \Omega^2 + \frac{\beta^2}{\rho}} \right)$$



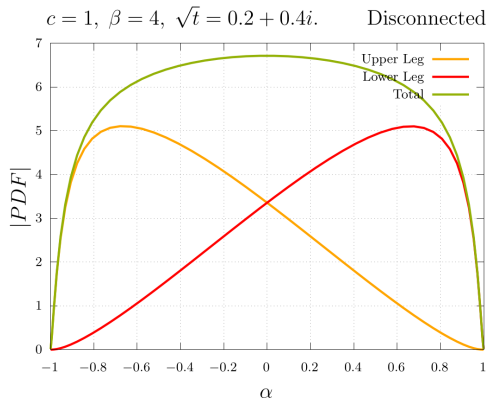
Put everything together



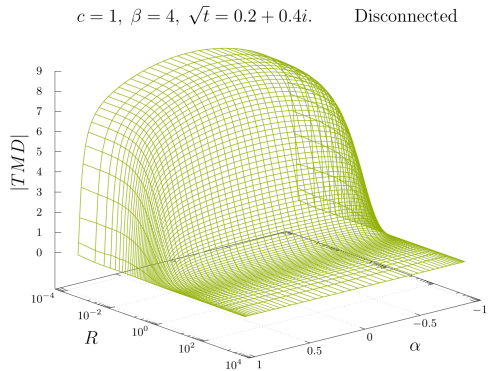
1. Calculate 4-point function
2. Calculate BSE
3. Do the loop integration
4. Project to LF

- We solve one triangle for each point in the R, α grid – **HPC Needed!**
 - $N_\alpha \times N_R \times N_Z$ 4-point functions!
- Obtained PDFs/TMDs for a system of two bound scalars ϕ .

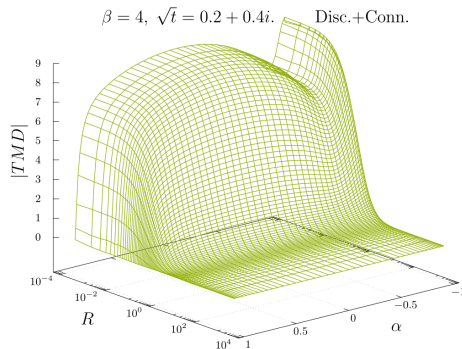
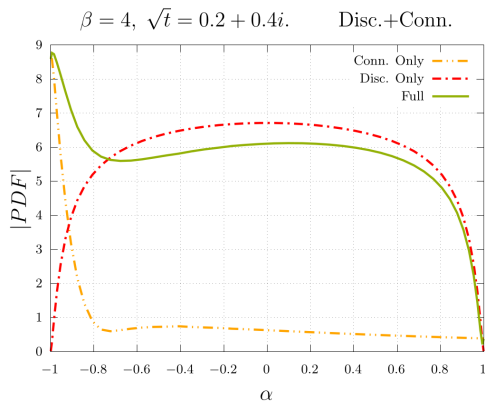
Results



■ Publication coming soon!



Results



■ Publication coming soon!

Conclusions

- Framework for calculation of PDFs/TMDs.
- Applications to hadronic structure calculations starting from self-consistent first principles calculations.
- Calculation of $q\bar{q}$ scattering equation – rich dynamics possible.
- Future applications to realistic QCD models soon.