EMISSION SHAPE EVOLUTION IN THE MONTE CARLO MODELS OF HEAVY-ION COLLISIONS

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JENSIS DE ROLA

TA*UNITER

OUTLINE

- Phenomenology
 - Correlation functions in femtoscopy
 - Lévy-type source function
- Monte Carlo models
 - EPOS₄ and UrQMD comparison
- D(
 ho) analysis
 - Results
 - EPOS₄, UrQMD at STAR energies
 - Other

CORRELATION FUNCTIONS IN FEMTOSCOPY

• Correlation function for bosons with the invariant momentum distributions N_1, N_2 : in this analysis: $\pi^+\pi^+$ and $\pi^-\pi^-$

$$C_2(p_1, p_2) = \frac{N_2(p_1, p_2)}{N_1(p_1)N_2(p_2)} \Longrightarrow C_2(q, K) \approx 1 + \left|\frac{\tilde{S}(q, K)}{\tilde{S}(0, K)}\right|^2,$$

where

$$\begin{split} N_2(p_1,p_2) &= \int S(x_2,p_2) |\Psi_2(x_1,x_2)|^2 \mathrm{d}^4 x_2 \mathrm{d}^4 x_1, \\ \tilde{S} &= \int S(x,k) e^{iqx} \mathrm{d}^4 x, \\ K &= \frac{k_1 + k_2}{2} \text{ (or } p_1,p_2), \\ q &= k_1 - k_2. \end{split}$$

• Event shape: $D(r,K) = \int S(x_1,K)S(x_2,K)d^4\rho = \int S\left(\rho + \frac{r}{2},K\right)S\left(\rho - \frac{r}{2},K\right)d^4\rho$ with average $\rho = \frac{x_1+x_2}{2}$ and relative $r = x_1 - x_2$

• With this,
$$C(q, K) = 1 + \int D(r, K)e^{-iqr} d^4r$$

THE LÉVY-TYPE SOURCE FUNCTION

• Lévy-stable distribution can be assumed and tested for the source function [Csörgő, Hegyi, Zajc Eur. Phys. J. C 36 (2004)], [Csörgő et al. AIP Conf. Proc. 828 (2006)]: $S \coloneqq L(x, \alpha, R),$

where

$$L = \frac{1}{(2\pi)^2} \int d^3 q e^{iqr} e^{-\frac{1}{2}|qR|^{\alpha}}$$

- Thus, the 2-particle correlation function $C_2(|k|) = 1 + \lambda e^{-(2R|k|)^{\alpha}},$ (note: $|k| = \frac{|q|}{2}$)
- The spatial distribution

$$D(\mathbf{r}) = L(\mathbf{r}, \alpha, 2^{\frac{1}{\alpha}R})$$

• In LCMS frame already *K*-independent: $D(r_{LCMS}) = \int d\Omega_{LCMS} dt D(r_{LCMS})$ (note: r_{LCMS} is the pair-separation vector in the lab frame)



[Plumberg & Heinz, Phys. Rev. C 98, 034910]



GAUSSIAN OR LÉVY SOURCE?

- Either assumption needs to be *tested*
- Gaussian: corresponds to $\alpha = 2$
- Experimental results in 1D: indicate non-Gaussian (α ≠ 2) behaviour, Lévy distribution describes the data better.[Kincses, Nagy, Csanád Communications Physics vol. 8 (55)], [Pórfy arXiv:2410.13975v2 [nucl-ex] 22 Apr 2025].

...as it will also appear in the simulations (?)



MONTE CARLO EVENT GENERATORS USED IN THIS ANALYSIS

EPOS



- Monte Carlo tool for simulating highenergy scatterings
- uses traditional S-matrix theory and modern concepts of perturbative QCD and saturation

K. Werner: Revealing a deep connection between factorization and saturation: New insight into modeling high-energy proton-proton and nucleus-nucleus scattering in the EPOS4 framework, Phys. Rev. C 108, 064903 – Published 6 December, 2023

K. Werner and B. Guiot *Perturbative QCD concerning light and heavy flavor in the EPOS4 framework*, Phys. Rev. C 108, 034904 (2023)

K. Werner Parallel scattering, saturation, and generalized Abramovskii-Gribov-Kancheli (AGK) theorem in the EPOS4 framework, with applications for heavy-ion collisions at $\sqrt{s_{NN}}$ of 5.02 TeV and 200 GeV, Phys. Rev. C 109, 034918 (2024)

K. Werner Core-corona procedure and microcanonical hadronization to understand strangeness enhancement in proton-proton and heavy ion collisions in the EPOS4 framework, Phys. Rev. C 109, 014910 (2024)

MONTE CARLO EVENT GENERATORS USED IN THIS ANALYSIS

EPOS₄



Ralewala

faiffa

nhoja Karjalan Runoja

Suomen kanfan muinosista ajoista.



HELEIRGIESÅ, 1835. Prántátty I. C. Frenckellin ja Pojan tyfóná.

In the latest version EPOS4:

- treatment of parton ladders completely redone
- saturation scales fixed by a prescription
- high precision applied in parallel scattering

K. Werner: Revealing a deep connection between factorization and saturation: New insight into modeling high-energy proton-proton and nucleus-nucleus scattering in the EPOS4 framework, Phys. Rev. C 108, 064903 – Published 6 December, 2023

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QUICK COMPARISON

EPOS4 [K. Werner Phys. Rev. C 109, 014910 (2024)], [Ayala et al. MDPI Particles vol. 6. (1)]

- Gribov—Regge-based + hydrodynamics
- Core-corona approach, viscous hydrodynamics, jet quenching
- QGP phase included
- Not a transport model per se
- Slower (hydro evolution)

UrOMD [Petersen, Hannah et al. arXiv:0805.0567 (2008).]

- Microscopic transport model
- Binary hadronic and string interactions
- No explicit QGP/hydro
- Excellent for studying hadronic transport and rescattering
- Faster (no hydro)

$D(\rho)$ ANALYSIS METHOD

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PION PAIR EVENT-BY-EVENT DISTRIBUTION

- Au+Au collisions simulated
- Separated into 10 $K_T = \frac{1}{2} \sqrt{K_x^2 + K_y^2}$ bins
- 1-dimensional projection along Bertsch—Pratt coordinates, cuts in η , p_T , q_{LCMS}
- Lévy-distribution fitted



PION PAIR EVENT-BY-EVENT DISTRIBUTION

- Lower energies (UrQMD): event averaging needed (low statistics)
- Averaging 100 events' $D(r_{LCMS})$ distributions before fitting



FIT PARAMETER AVG. & STD. DEV.

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- Histograms of fit parameters (α , R)
- Fit parameter mean and standard deviation obtained
- As a function of K_T :
 - *α* consistent with constant within uncertainty
 - R decreasing with K_T



EPOS4 $\alpha(K_T)$, $R(K_T)$

• In all analysed energies: *α* consistent with constant within uncertainty, *R* decreasing with *K*_T



EPOS4 $\alpha(K_T)$ centrality dependence

• Generally, a decreasing trend in parameter values with more peripheral collisions



EPOS4 $R(K_T)$ centrality dependence

- Generally, a decreasing trend in parameter values with more peripheral collisions
- Fits for both parameters tends to be less reliable towards higher centralities & higher K_T bins (lack of statistics without event averaging)



URQMD $\alpha(K_T)$, $R(K_T)$

- Same trend with UrQMD analysis
- Event averaging applied (D(r) fitted from 100-100 events)
- Still mostly too low statistics towards higher K_T bins



K_T -averaged results

- Histograms of fit parameters averaged (summed)
- Mean and standard deviation (~uncertainty) obtained







RESULT: EPOS4 & URQMD AT RHIC ENERGIES



RESULT: EPOS4 & URQMD AT RHIC ENERGIES

• Au+Au

- Quantitatively (mostly within the "std. dev. uncertainty") not far from the $\sqrt{s_{NN}} =$ 3 to 60 GeV data
- Larger differences toward larger energies of collider mode & smallest energies of fixed-target mode
- Different trend from the data both in UrQMD & EPOS₄ results



ANOTHER RESULT: URQMD AT NA61 ENERGIES

• Collisions 13A to 150A GeV/c energy (Ar+Sc)

[The NA61/SHINE Collaboration, arXiv:2503.22484v2 [nucl-ex] 31 Mar 2025]

- UrQMD 3D analysis done by B. Pórfy
- Quantitatively, UrQMD close to experimental data for $\sqrt{s_{NN}} = 6$ to 10 GeV
- Different trend compared to data



SUMMARY

- Au+Au simulations done at RHIC energies
- Source functions obtained from EPOS₄ and UrQMD
- Lévy-fits performed
- Results quantitatively close to experimental data but seemingly different trend
- Next steps: correlation afterburner, higher statistics



THANKYOU FORYOUR ATTENTION!





BACKUP SLIDES

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TODO: CORRELATION AFTERBURNER

CRAB: Correlation After Burner

- Generate correlation functions from semiclassical transport codes
- with impact-parameter averaging, experimental acceptances, experimental resolution, kinematic cuts
- Input: files of phase space points (final momentum & point of last interaction of generated particles e.g. in EPOS₄), impact parameters
- Efficient 3D processing possible
- Work-in-progress



$\mathsf{EPOS}_4 R(\langle K_T \rangle)$



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EPOS₄ $\alpha(\langle K_T \rangle)$



28/23

EVENT-BY-EVENT PION SOURCE FUNCTION LÉVY FITS

 $\sqrt{s_{NN}} = 7.7 \text{ GeV}$













 $\sqrt{s_{NN}} = 9.2 \text{ GeV}$





 $\sqrt{s_{NN}} = 11.5 \text{ GeV}$























 $\sqrt{s_{NN}} = 27 \text{ GeV}$



 α vs R, 0.175<K_T<0.225, $\sqrt{s_{NN}}$ =7.7GeV α vs R, 0.325<K_T<0.375, $\sqrt{s_{_{NN}}}$ =7.7GeV α vs R, 0.225<K_T<0.275, $\sqrt{s_{NN}}$ =7.7GeV α vs R, 0.275<K_T<0.325, $\sqrt{s_{NN}}$ =7.7GeV α vs R, 0.375<K_T<0.425, $\sqrt{s_{NN}}$ =7.7GeV <u>ل</u> لا س 0.2 0.4 0.6 0.8 1 1.2 1.4 1.6 1.8 2 0.2 0.4 0.6 0.8 1 1.2 1.4 1.6 1.8 2 D 0.2 0.4 0.6 0.8 1 1.2 1.4 1.6 1.8 0 0.2 0.4 0.6 0.8 1 1.2 1.4 1.6 1.8 2 0.2 0.4 0.6 0.8 1 1.2 1.4 1.6 1.8 α vs R, 0.425<K_T<0.475, $\sqrt{s_{NN}}$ =7.7GeV α vs R, 0.475<K_T<0.525, $\sqrt{s_{NN}}$ =7.7GeV α vs R, 0.525<K_T<0.575, $\sqrt{s_{NN}}$ =7.7GeV α vs R, 0.575<K_T<0.625, $\sqrt{s_{NN}}$ =7.7GeV α vs R, all K_T, $\sqrt{s_{NN}}$ =7.7GeV R [fm] R [fm] 0 0.2 0.4 0.6 0.8 1 1.2 1.4 1.6 1.8 2 0 0.2 0.4 0.6 0.8 1 1.2 1.4 1.6 1.8 2 0 0.2 0.4 0.6 0.8 1 1.2 1.4 1.6 1.8 2 0 0.2 0.4 0.6 0.8 1 1.2 1.4 1.6 1.8 2 1 1.2 1.4 0.8 1.6

 $\alpha \vee S R \sqrt{S_{NN}} = 7.7 \text{ GeV}$



$\alpha \vee S R \sqrt{S_{NN}} = 9.2 \text{ GeV}$



 $\alpha \vee S R \sqrt{S_{NN}} = 11.5 \text{ GeV}$



 $\alpha \vee S R \sqrt{S_{NN}} = 14.5 \text{ GeV}$



$\alpha \vee S R \sqrt{S_{NN}} = 19.6 \text{ GeV}$

 α vs R, 0.175<K_T<0.225, $\sqrt{s_{NN}}$ =27GeV α vs R, 0.225<K_T<0.275, $\sqrt{s_{NN}}$ =27GeV α vs R, 0.275<K_T<0.325, $\sqrt{s_{NN}}$ =27GeV α vs R, 0.325<K_T<0.375, $\sqrt{s_{NN}}$ =27GeV α vs R, 0.375<K_T<0.425, $\sqrt{s_{NN}}$ =27GeV <u>الم</u> 1.8 10 10 0.2 0.4 0.6 0.8 1 1.2 1.4 1.6 1.8 2 0.2 0.4 0.6 0.8 1 1.2 1.4 1.6 1.8 2 0 0.2 0.4 0.6 0.8 1 1.2 1.4 1.6 1.8 2 0.2 0.4 0.6 0.8 1 1.2 1.4 1.6 1.8 α vs R, 0.425<K_T<0.475, $\sqrt{s_{NN}}$ =27GeV α vs R, 0.475<K_T<0.525, $\sqrt{s_{NN}}$ =27GeV α vs R, 0.525<K_T<0.575, $\sqrt{s_{NN}}$ =27GeV α vs R, 0.575<K_T<0.625, $\sqrt{s_{NN}}$ =27GeV α vs R, all K_T, $\sqrt{s_{NN}}$ =27GeV R [fm] 0 0.2 0.4 0.6 0.8 1 1.2 1.4 1.6 1.8 2 0 0.2 0.4 0.6 0.8 1 1.2 1.4 1.6 1.8 2 0 0.2 0.4 0.6 0.8 1 1.2 1.4 1.6 1.8 2 0 0.2 0.4 0.6 0.8 1 1.2 1.4 1.6 1.8 2 1.2 1.4 0.8

 $\alpha \vee S R \sqrt{S_{NN}} = 27 \text{ GeV}$

IMPACT PARAMETER (BIM) AND NPART DISTRIBUTIONS, $\sqrt{s_{NN}} = 7.7 \text{ GeV}$



IMPACT PARAMETER (BIM) AND NPART DISTRIBUTIONS, $\sqrt{s_{NN}} = 9.2 \text{ GeV}$



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IMPACT PARAMETER (BIM) AND NPART DISTRIBUTIONS, $\sqrt{s_{NN}} = 11.5 \text{ GeV}$



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IMPACT PARAMETER (BIM) AND NPART DISTRIBUTIONS, $\sqrt{s_{NN}} = 14.5$ GeV



IMPACT PARAMETER (BIM) AND NPART DISTRIBUTIONS, $\sqrt{s_{NN}} = 19.6 \text{ GeV}$



2500

Impact parameter (BIM) and NPART distributions, $\sqrt{s_{NN}} = 27 \text{ GeV}$







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