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# Impact of elliptic flow on higher order flow coefficients in heavy-ion collisions at the LHC



HUN  
REN



Based On:

S. Prasad, N. Mallick, S. Tripathy, R. Sahoo, Phys. Rev. D **107** (2023) 074011.

S. Prasad, A. Menon K. R., R. Sahoo, N. Mallick, Phys. Lett. B **868** (2025) 139753.

Presented By:

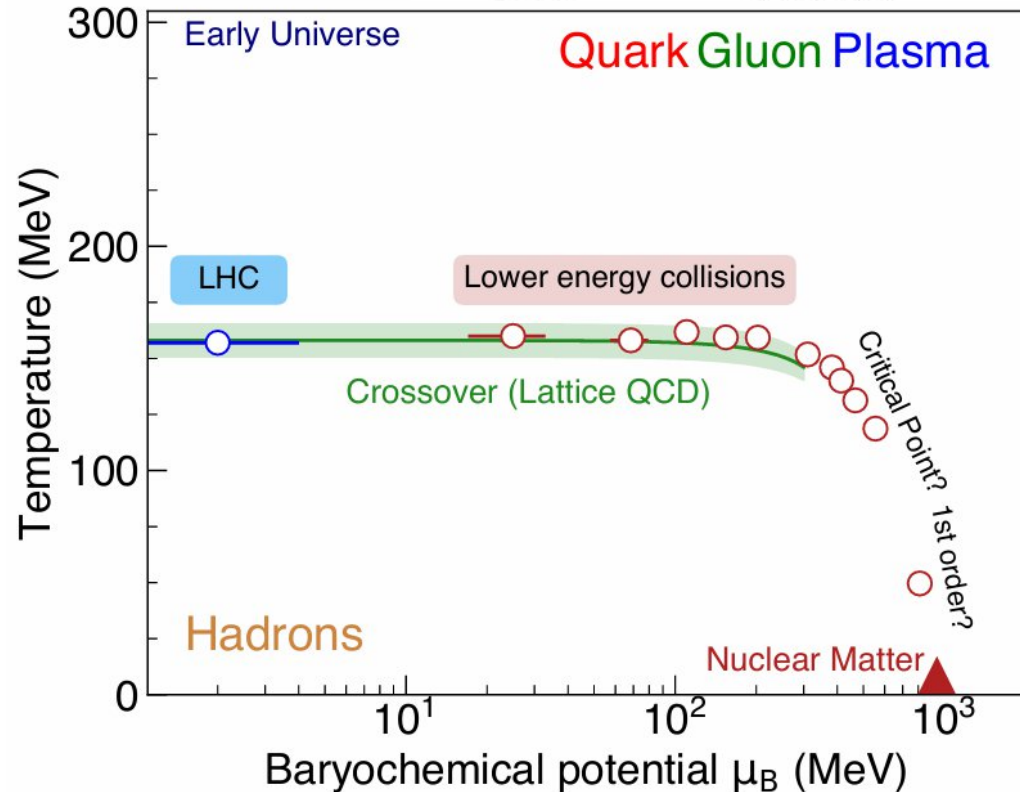
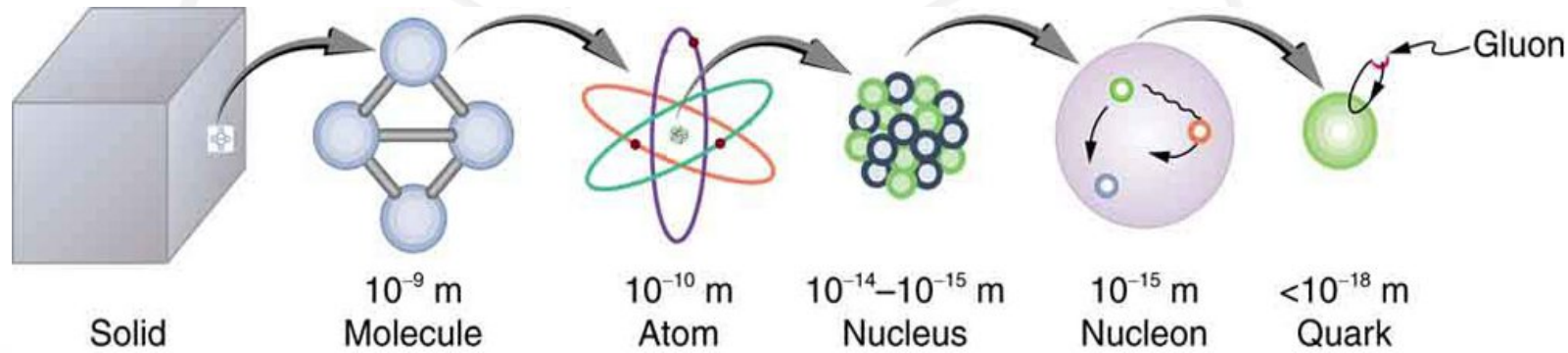
**Suraj Prasad**

✉ Suraj.prasad@cern.ch

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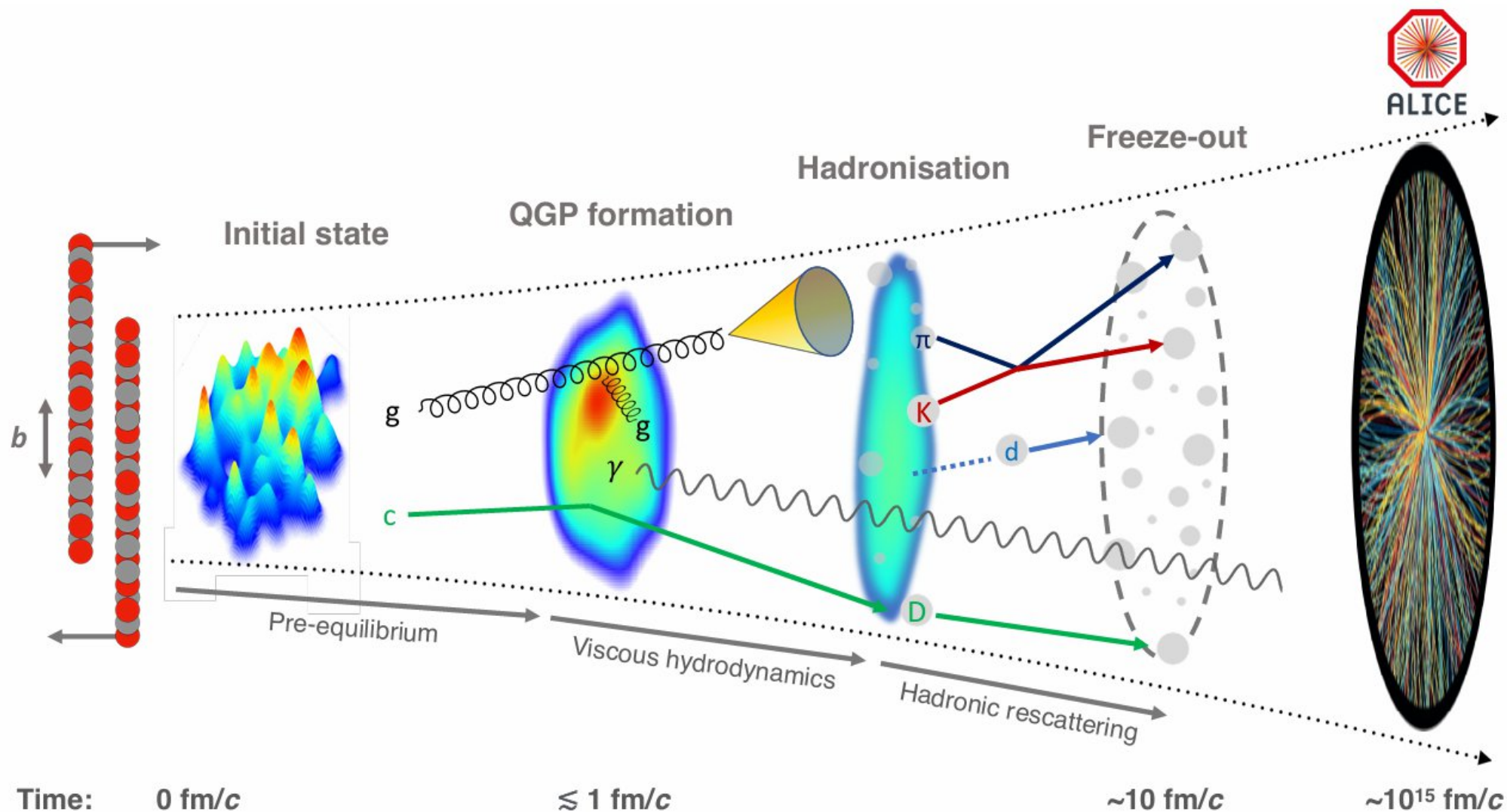
08 August 2025

# QCD Phase diagram



- The quarks are bound through gluons inside hadrons, such as , protons and neutrons. No isolated parton can be observed → Color Confinement
- If the matter is heated up to incredible temperatures ( $\sim 10^5$  times the temperature in core of the Sun) a “soup” of partons can be formed where quarks and gluons can move freely
- A similar state can be achieved by increasing the baryochemical potential

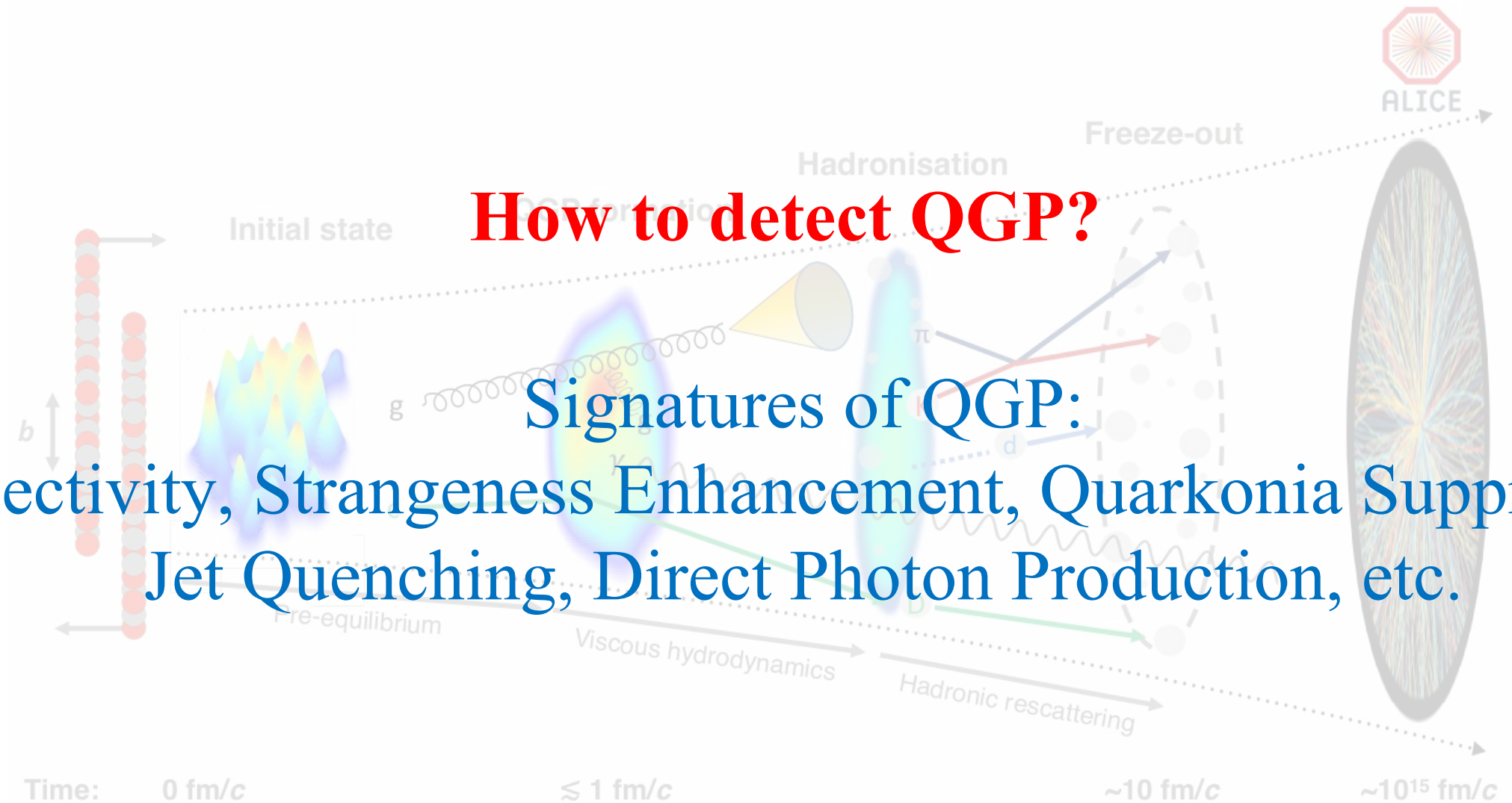
# Heavy-Ion Collisions and Quark Gluon Plasma



S. Acharya et al. (ALICE), Eur. Phys. J. C 84, 813 (2024)

**How to detect QGP?**

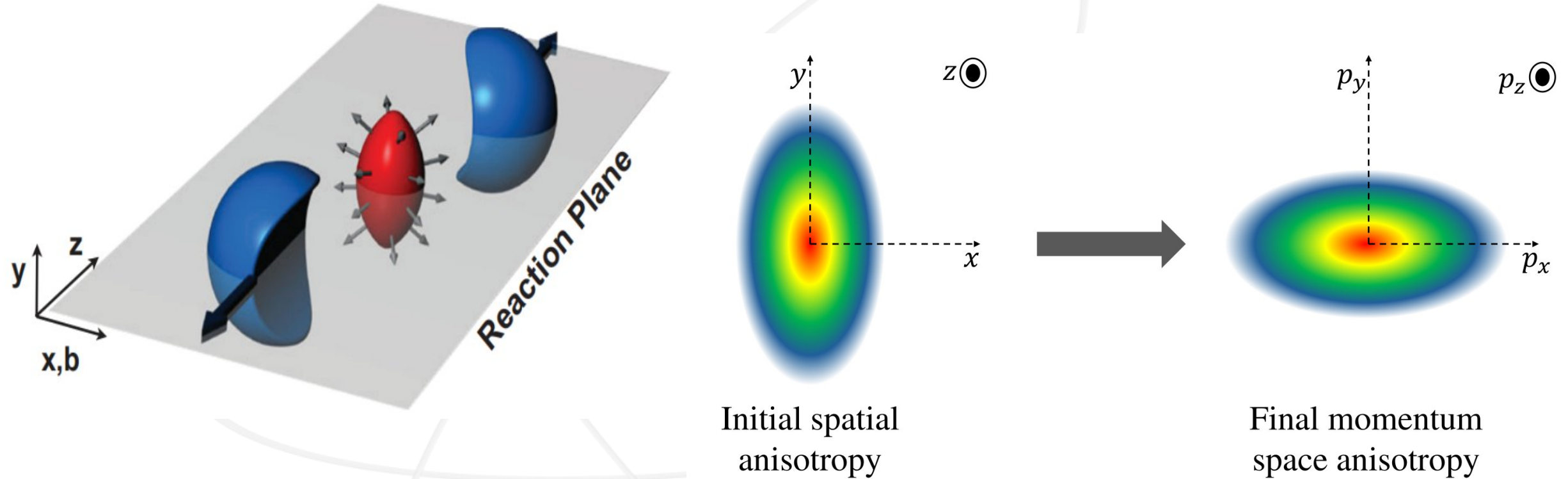
**Signatures of QGP:**  
Collectivity, Strangeness Enhancement, Quarkonia Suppression,  
Jet Quenching, Direct Photon Production, etc.



S. Acharya et al. (ALICE), Eur. Phys. J. C 84, 813 (2024)



# Anisotropic Flow

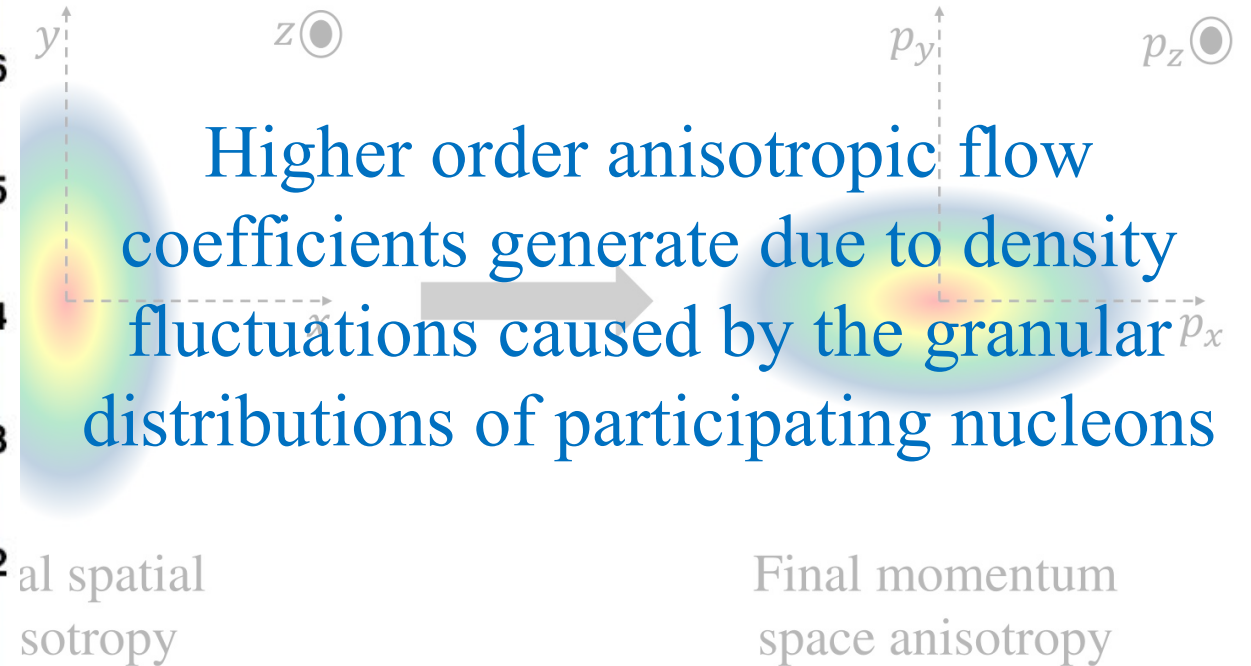
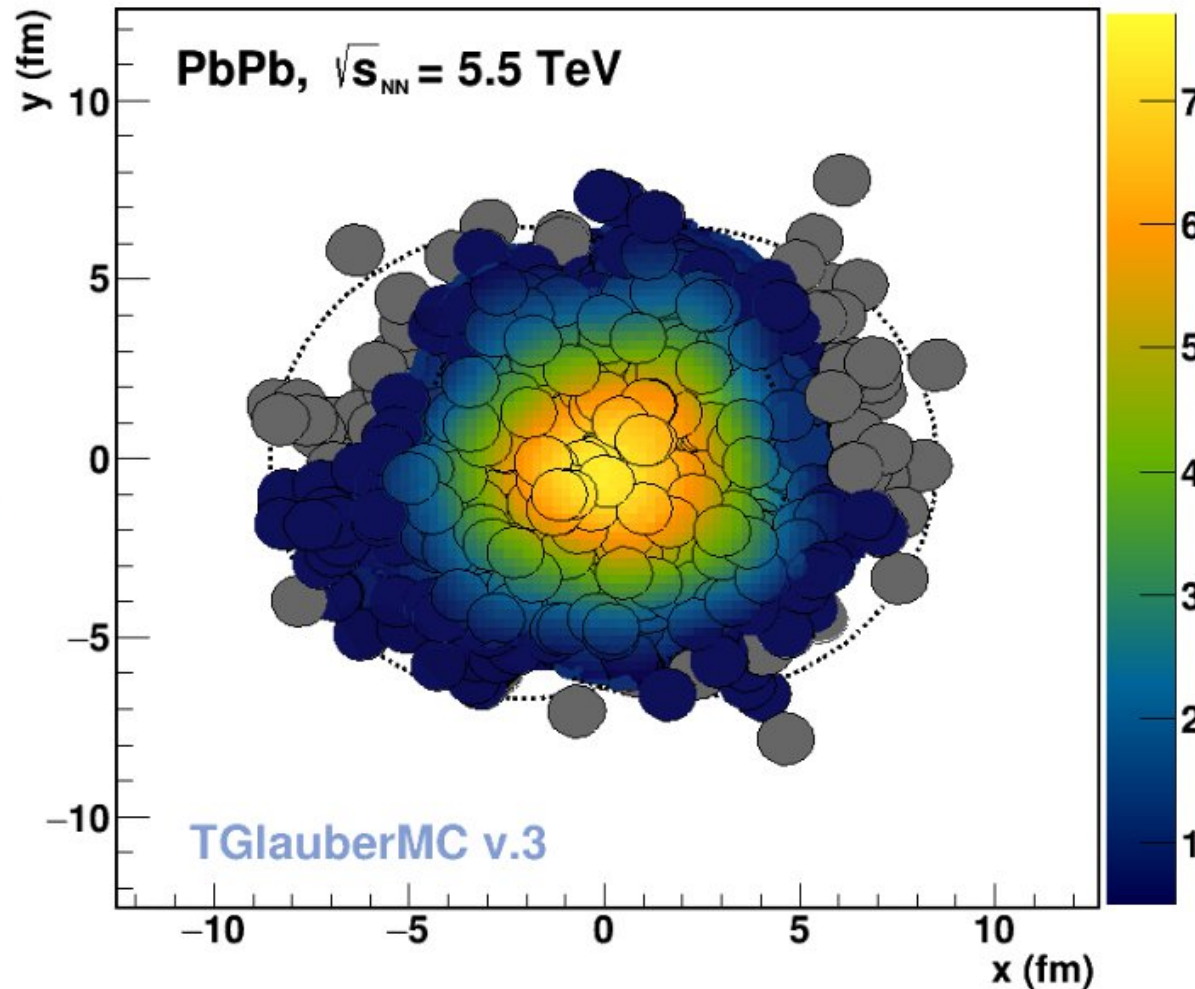


- Anisotropic flow: hydrodynamic response to spatial deformation of the initial participant distribution

Annual Review of Nuclear and Particle Science 71:315-44 (2021)

N. Mallick, R. Sahoo, S. Tripathy, and A. Ortiz, J. Phys. G 48, 045104 (2021)  
ALICE Collaboration, JHEP 1506, 190 (2015)

# Anisotropic Flow



- Anisotropic flow: hydrodynamic response to spatial deformation of the initial participant distribution

# Anisotropic Flow

$$E \frac{d^3 N}{dp^3} = \frac{d^2 N}{2\pi p_T dp_T dy} \left( 1 + 2 \sum_{n=1}^{\infty} v_n \cos[n(\phi - \psi_n)] \right) = \frac{d^2 N}{2\pi p_T dp_T dy} \left( 1 + \underbrace{2v_1 \cos(\phi - \psi_1)}_{\text{Directed flow}} + \underbrace{2v_2 \cos[2(\phi - \psi_2)]}_{\text{Elliptic flow}} + \underbrace{2v_3 \cos[3(\phi - \psi_3)]}_{\text{Triangular flow}} + \dots \right)$$

- Anisotropic flow coefficients are excellent probes for transport properties of the medium
- Higher order coefficients are more sensitive to system response/transport properties of the medium

Elliptic and triangular flow interfere with measurements of higher order harmonics

$$\begin{aligned} v_4 e^{i4\Phi_4} &= a_0 \epsilon_4 e^{i4\psi_4} + a_1 (\epsilon_2 e^{i2\psi_2})^2 + \dots \\ &= c_0 e^{i4\psi_4} + c_1 (v_2 e^{i2\Phi_2})^2 + \dots, \end{aligned}$$

$$\begin{aligned} v_5 e^{i5\Phi_5} &= a_0 \epsilon_5 e^{i5\psi_5} + a_1 \epsilon_2 e^{i2\psi_2} \epsilon_3 e^{i3\psi_3} + \dots \\ &= c_0 e^{i5\psi_5} + c_1 v_2 v_3 e^{i(2\Phi_2 + 3\Phi_3)} + \dots, \end{aligned}$$

Large fluctuations of flow tend to interfere with physical interpretation of observed phenomena

**Event Shape  
Analysis**

# Multiparticle Q-cumulant method

## Event plane and average method

$$v_n = \langle \cos[n(\phi - \psi_n)] \rangle \quad (\text{Affected by non-flow and requires } \psi_n)$$

## Multiparticle Q-cumulant method

- Flow vector  $Q_n = \sum_{j=1}^M e^{in\phi_j}$

- The 2- and 4-particle cummulants are:

$$\langle 2 \rangle = \frac{|Q_n|^2 - M}{M(M-1)},$$

$$\langle 4 \rangle = \frac{|Q_n|^4 + |Q_{2n}|^2 - 2 \cdot \text{Re}[Q_{2n} Q_n^* Q_n^*]}{M(M-1)(M-2)(M-3)} - 2 \frac{2(M-2) \cdot |Q_n|^2 - M(M-3)}{M(M-1)(M-2)},$$

$$c_n\{2\} = \langle \langle 2 \rangle \rangle,$$

$$c_n\{4\} = \langle \langle 4 \rangle \rangle - 2 \cdot \langle \langle 2 \rangle \rangle^2$$

$$v_n\{2\} = \sqrt{c_n\{2\}},$$
$$v_n\{4\} = \sqrt[4]{-c_n\{4\}}.$$



# Multiparticle Q-cumulant method

## Suppressing the non-flow contribution:

- Kinematical cut: 2 sub-events, A&B are introduced, with a rapidity gap:

$$\langle 2 \rangle_{\Delta\eta} = \frac{Q_n^A \cdot Q_n^{B*}}{M_A \cdot M_B},$$

$$\longrightarrow v_n\{2, |\Delta\eta|\}(p_T) = \frac{d_n\{2, |\Delta\eta|\}}{\sqrt{c_n\{2, |\Delta\eta|\}}}$$

- Differential flow cumulants:

$$d_n\{2\} = \langle\langle 2' \rangle\rangle,$$

$$d_n\{4\} = \langle\langle 4' \rangle\rangle - 2\langle\langle 2' \rangle\rangle\langle\langle 2 \rangle\rangle$$

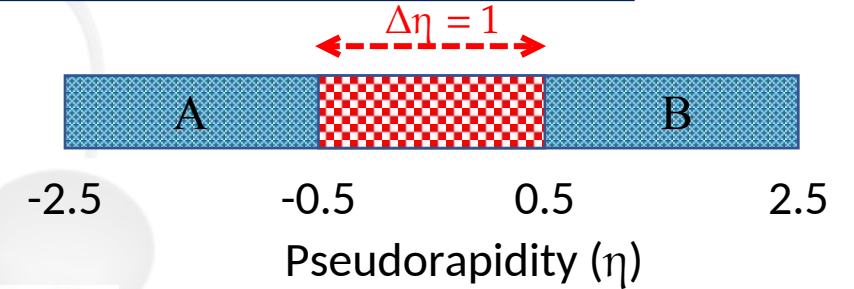
$$\longrightarrow d_n\{2, |\Delta\eta|\} = \langle\langle 2' \rangle\rangle_{\Delta\eta}$$

- Mean and the fluctuations of the flow & ratio:

$$\langle v_n \rangle = \sqrt{\frac{v_n^2\{2, |\Delta\eta|\} + v_n^2\{4\}}{2}}$$

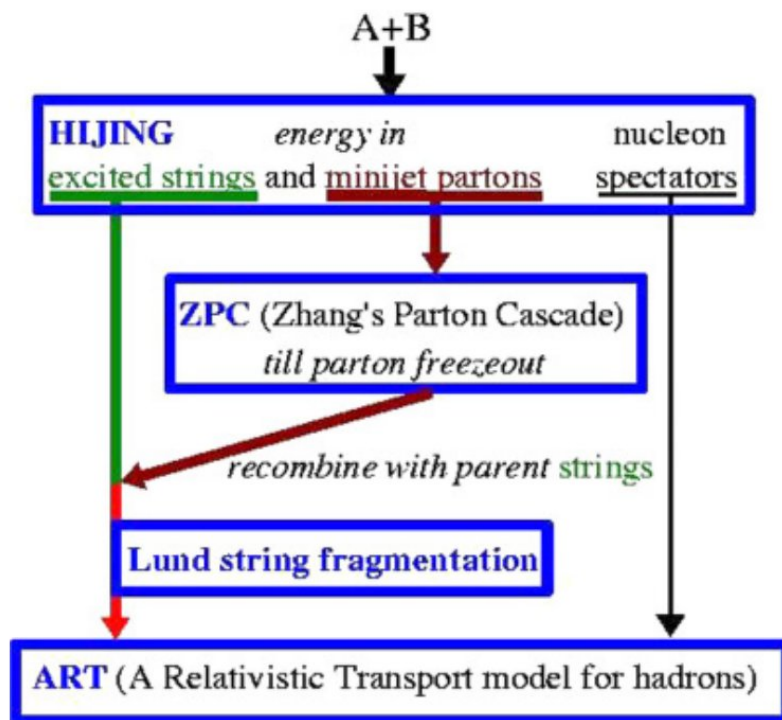
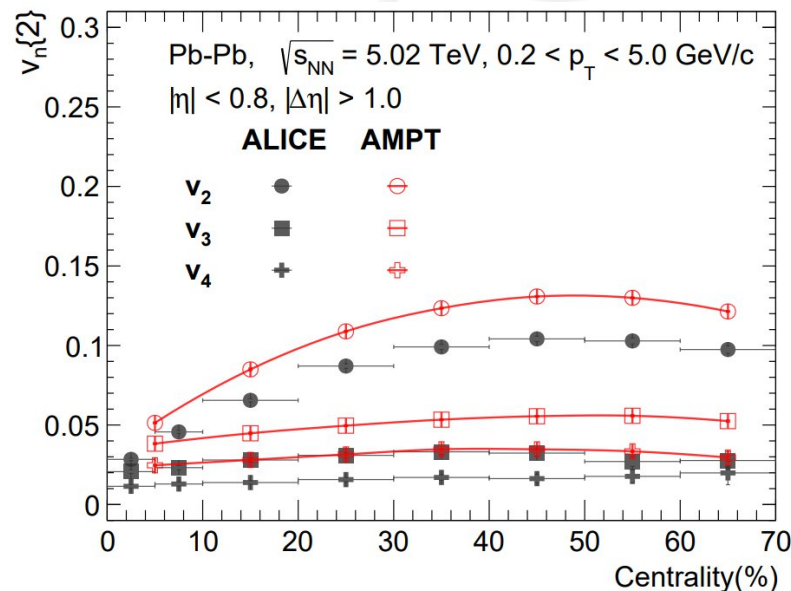
$$\sigma_{v_n} = \sqrt{\frac{v_n^2\{2, |\Delta\eta|\} - v_n^2\{4\}}{2}}$$

$$F(v_n) = \frac{\sigma_{v_n}}{\langle v_n \rangle}$$

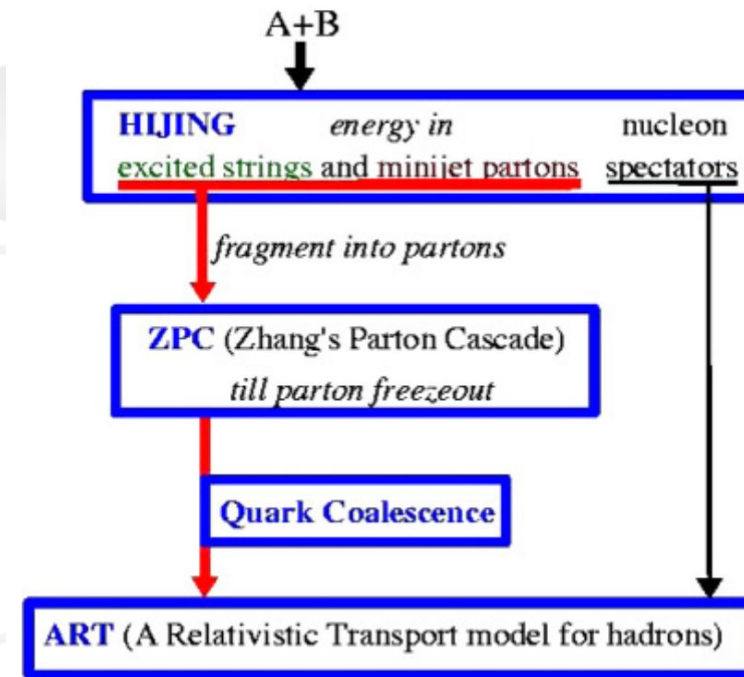


# A Multi-Phase Transport Model (AMPT)

- Initialisation of collisions (HIJING)
- Parton Transport (ZPC)
- Hadronisation (Lund string fragmentation / Quark coalescence)
- Hadron Transport (ART)



Default Version



String Melting Version

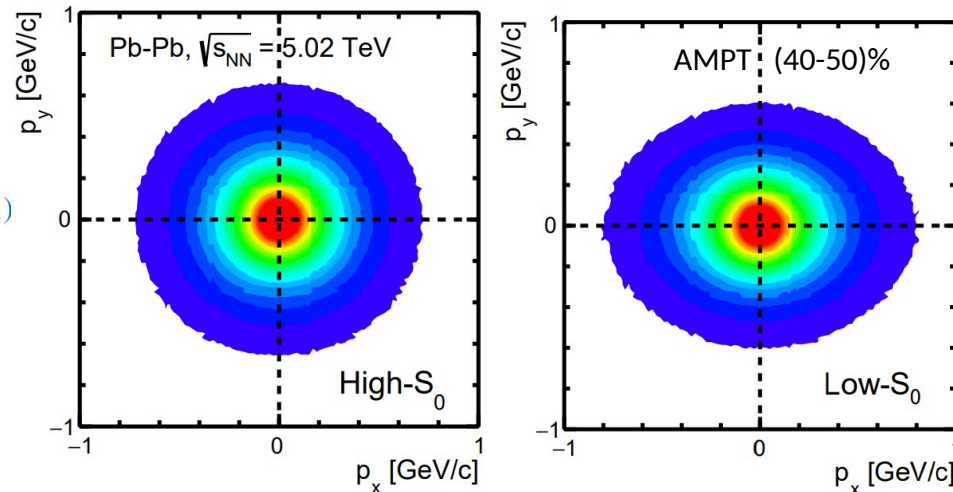
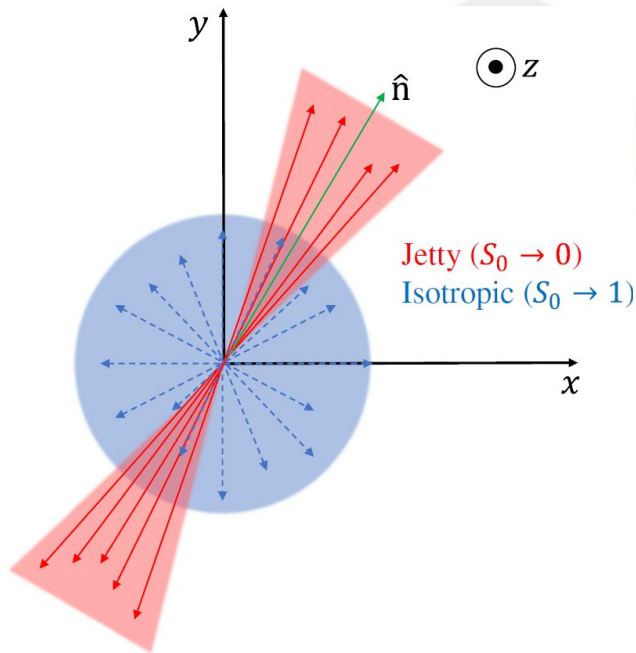
S. Prasad, A. Menon K. R., R. Sahoo, N. Mallick, Phys. Lett. B 868 (2025) 139753  
 Zi-Wei Lin, Che Ming Ko, Bao-An Li, Bin Zhang, and Subrata Pal, Phys. Rev. C 72, 064901 (2005)

# Transverse Spherocity

- Event shape observable: separates events based on geometrical shapes

$$S_0 = \frac{\pi^2}{4} \min_{\hat{n}} \left( \frac{\sum_i |\vec{p}_{T_i} \times \hat{n}|^2}{\sum_i |\vec{p}_{T_i}|^2} \right)$$

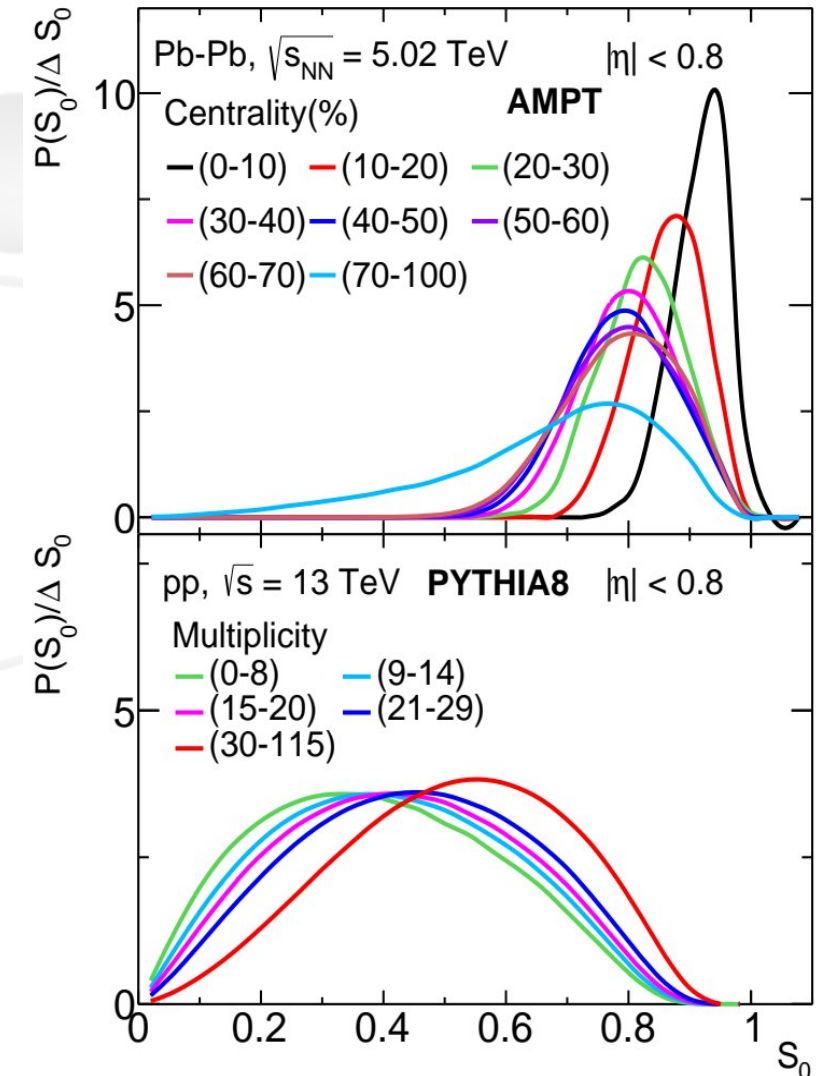
- “Isotropic” limit (Soft-QCD processes)
- “Pencil-like” or “jetty” limit (Hard-QCD processes)



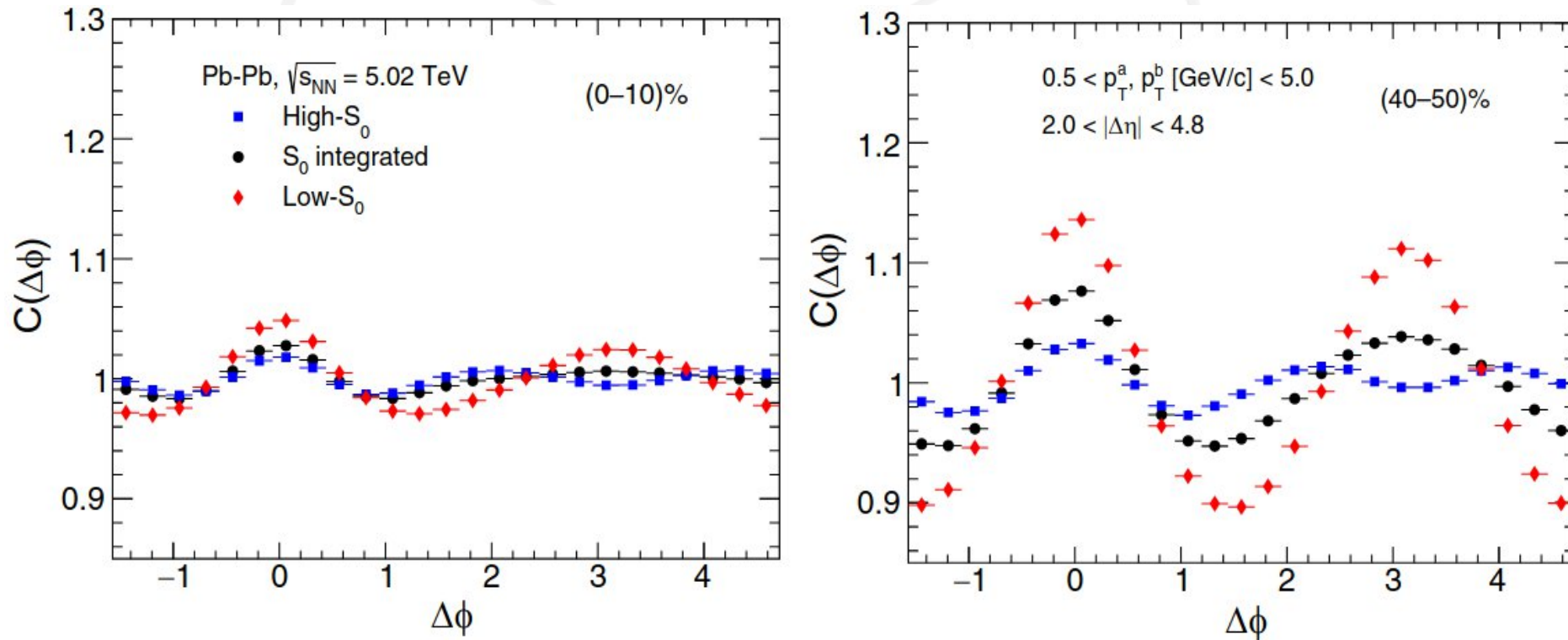
N. Mallick, R. Sahoo, S. Tripathy, A. Ortiz, J. Phys. G 48, 045104 (2021)

S. Prasad, N. Mallick, S. Tripathy, and R. Sahoo, Phys. Rev. D 107, 074011 (2023)

S. Prasad, N. Mallick, D. Behera, R. Sahoo and S. Tripathy, Sci. Rep. 12, 3917 (2022)



# Two-Particle Correlation



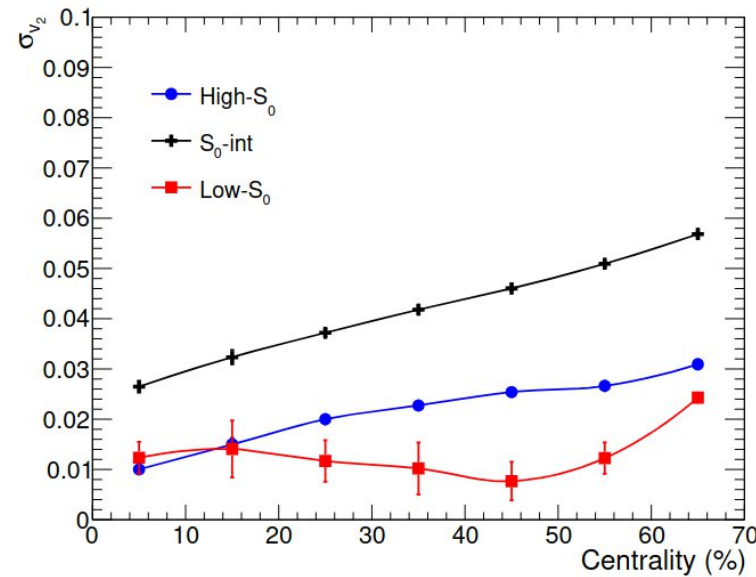
S. Prasad, N. Mallick, S. Tripathy, and R. Sahoo, Phys. Rev. D 107, 074011 (2023)

$$C(\Delta\phi) = \frac{dN_{\text{pairs}}}{d\Delta\phi} \propto \left[ 1 + 2 \sum_{n=1}^{\infty} v_{n,n}(p_T^a, p_T^b) \cos n\Delta\phi \right]$$

$$C(\Delta\phi) = \frac{dN_{\text{pairs}}}{d\Delta\phi} = A \times \frac{S(\Delta\phi)}{B(\Delta\phi)} = A \times \frac{\int S(\Delta\eta, \Delta\phi) d\Delta\eta}{\int B(\Delta\eta, \Delta\phi) d\Delta\eta}$$

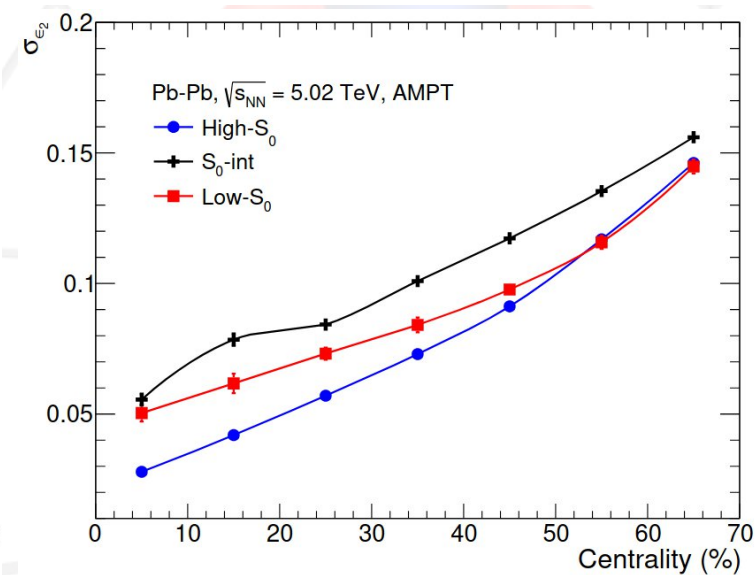
- Correlation function is stronger for (40-50)% centrality than (0-10)% centrality → More flow in midcentral collisions
- Clear transverse spherocity dependence; 3 peak structure in high- $S_0$  events → Large triangular flow
- Low- $S_0$  events show larger two peak structure → Large elliptic flow





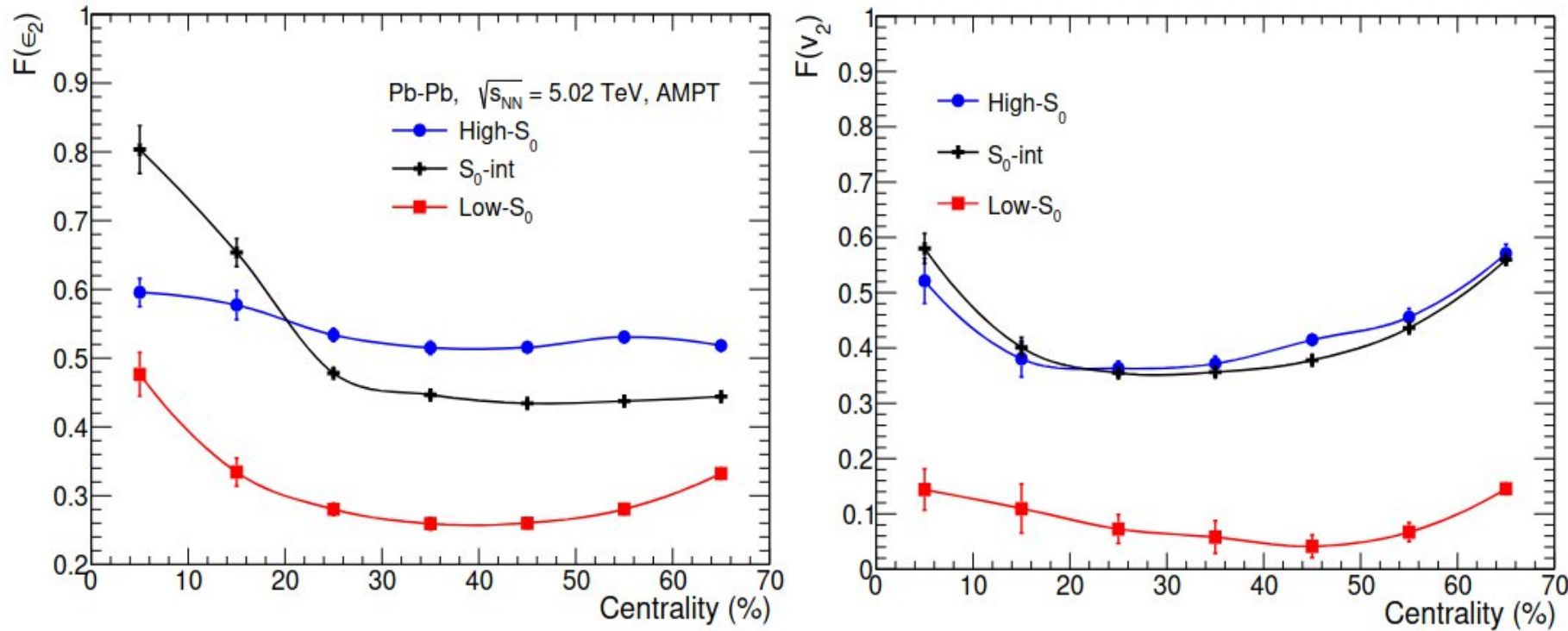
$$\epsilon_n = \frac{\sqrt{\langle r^n \cos(n\phi_{\text{part}}) \rangle^2 + \langle r^n \sin(n\phi_{\text{part}}) \rangle^2}}{\langle r^n \rangle}$$

- Flow estimated using multi-particle Q-cumulant method
- Rise of  $v_2$  from central to mid-central collisions followed by a decrease towards peripheral collisions  $\longrightarrow$  Contribution from initial stage
- Clear  $v_2$  dependence on  $S_0 \longrightarrow$  Low  $S_0$  events have large  $v_2$  and vice versa
- $\sigma_{v_2}$  is large for  $S_0$ -int events and peripheral collisions
- $\sigma_{v_2}$  is smallest for low- $S_0$  events
- $\sigma_{\epsilon_2}$  shows a different  $S_0$  dependence



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# Relative Fluctuations



$$\langle v_n \rangle = \sqrt{\frac{v_n^2\{2, |\Delta\eta|\} + v_n^2\{4\}}{2}}$$

$$\sigma_{v_n} = \sqrt{\frac{v_n^2\{2, |\Delta\eta|\} - v_n^2\{4\}}{2}}$$

$$F(v_2) = \frac{\sigma_{v_2}}{\langle v_2 \rangle}$$

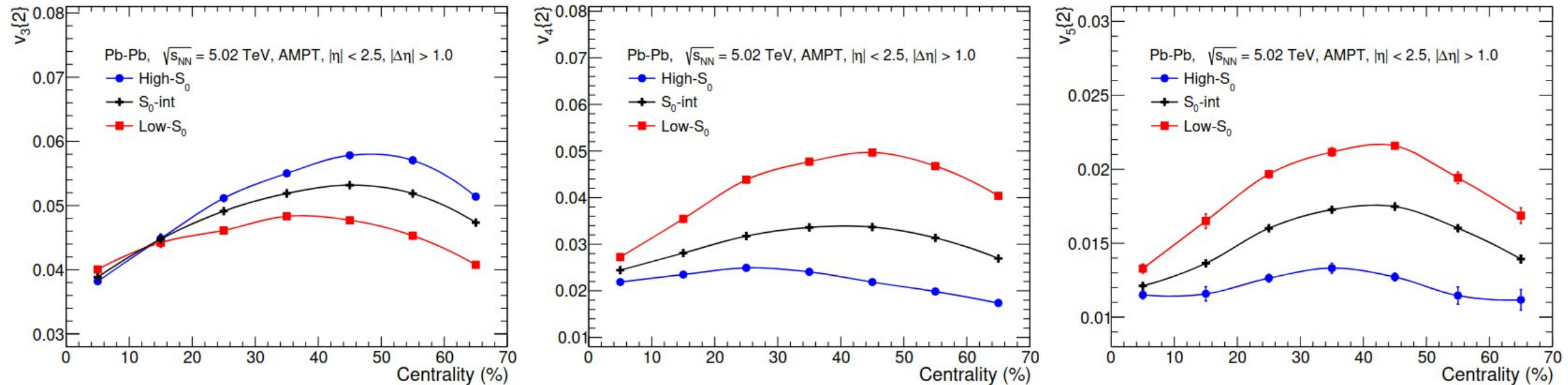
$$\sigma_{\epsilon_2} = \sqrt{\langle \epsilon_2^2 \rangle - \langle \epsilon_2 \rangle^2}$$

$$F(\epsilon_2) = \frac{\sigma_{\epsilon_2}}{\langle \epsilon_2 \rangle}$$

- Central collisions have large  $F(v_2)$  —> Decreases towards mid-central followed by a rise towards peripheral collisions
- In central collisions,  $S_0$ -int events have largest  $F(v_2)$  and  $F(\epsilon_2)$
- Low- $S_0$  events have smallest  $F(v_2)$  and  $F(\epsilon_2)$

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# Higher Order Flow Coefficients

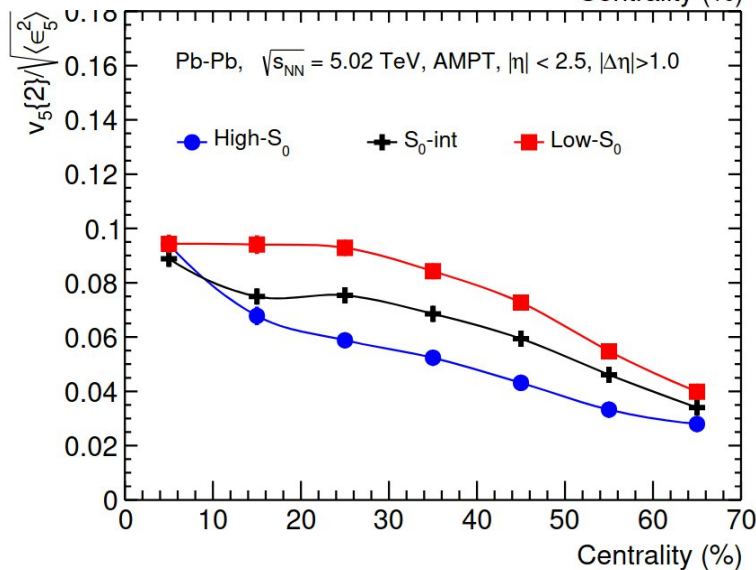
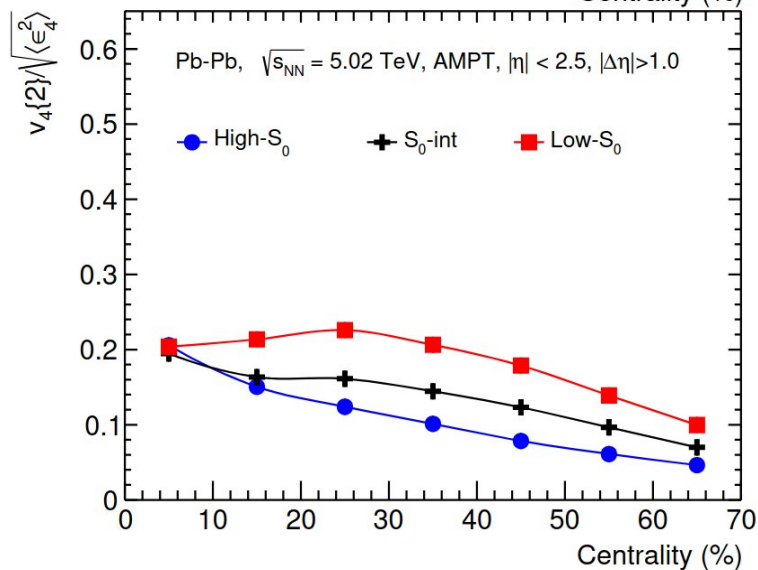
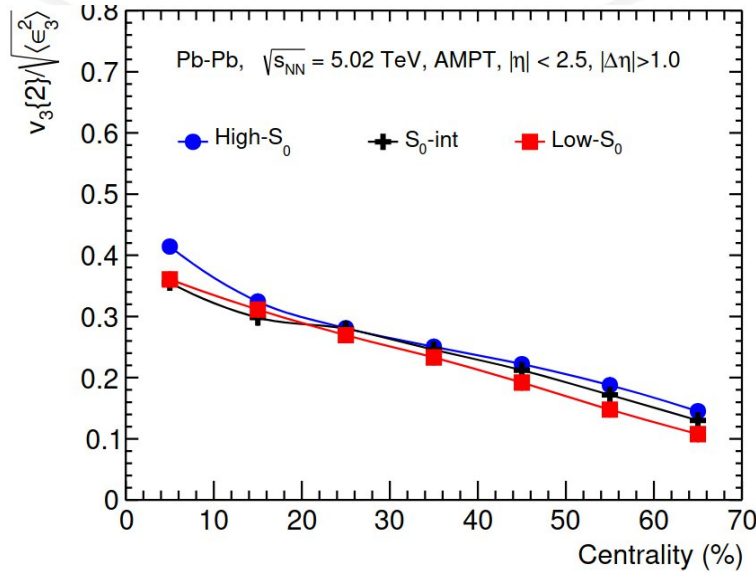
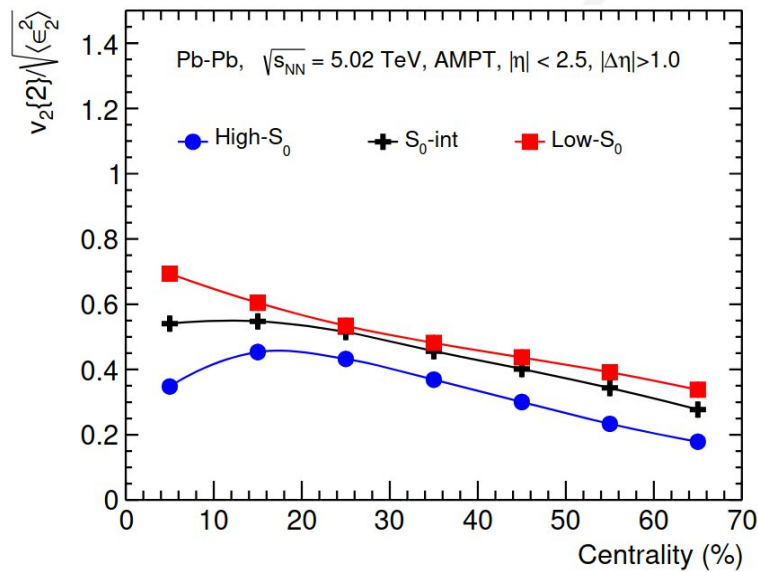


- Centrality dependence weakens for  $v_n$  when going from  $n=3$  to  $n=5$   $\rightarrow$  Stronger damping effects for  $v_5$  than  $v_4$  than  $v_3$
- $v_3$  shows an anti-correlation with  $S_0$  selection  $\rightarrow$  Attributed to anti-correlation between  $v_2$  and  $v_3$
- $v_4$  and  $v_5$  has non-linear contributions from  $v_2$   $\rightarrow$  Reflected in their anti-correlation with  $S_0$  selection
- High- $S_0$  events have lowest contribution from  $v_2$  and Low- $S_0$  events have highest contribution from  $v_2$

S. Prasad, A. Menon K. R., R. Sahoo, N. Mallick, Phys. Lett. B 868 (2025) 139753



# Spherocity Dependence of $v_n\{2\}/\sqrt{\langle\epsilon_n^2\rangle}$



- System response to geometrical anisotropy is different from that arising due to fluctuations
- Isotropisation adversely affects the transformation of  $\epsilon_2$  to  $v_2$
- High- $S_0$  events show stronger damping effects for  $n=4$  and  $n=5$
- Low- $S_0$  events for  $n=4$  shows a small rise near (20-30)% centrality
- Low- $S_0$  events show saturation behaviour near central collisions for  $n=5$

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# Summary

- $S_0$  based event selection are capable of selecting events with larger or smaller  $v_2$  than in comparison to  $q_2$
- $S_0$  based event selection is better than  $q_2$  event in low-multiplicity environments  $\rightarrow$  Useful for O-O, Ne-Ne and pO collisions
- $v_2$  is anti-correlated to event selection based on  $S_0$ ,  $v_3$  has positive correlation
- $v_2$  fluctuations are large for high- $S_0$  events, smaller for low- $S_0$  events
- System response to geometrical anisotropy is different from that arising due to fluctuations
- Isotropisation adversely affects the transformation of  $\epsilon_2$  to  $v_2$
- High- $S_0$  events show stronger damping effects
- With  $S_0$ , one can select events with smaller values of  $v_2$   $\rightarrow$  useful to study higher order harmonics



**Thank You**  
**for Your Attention**

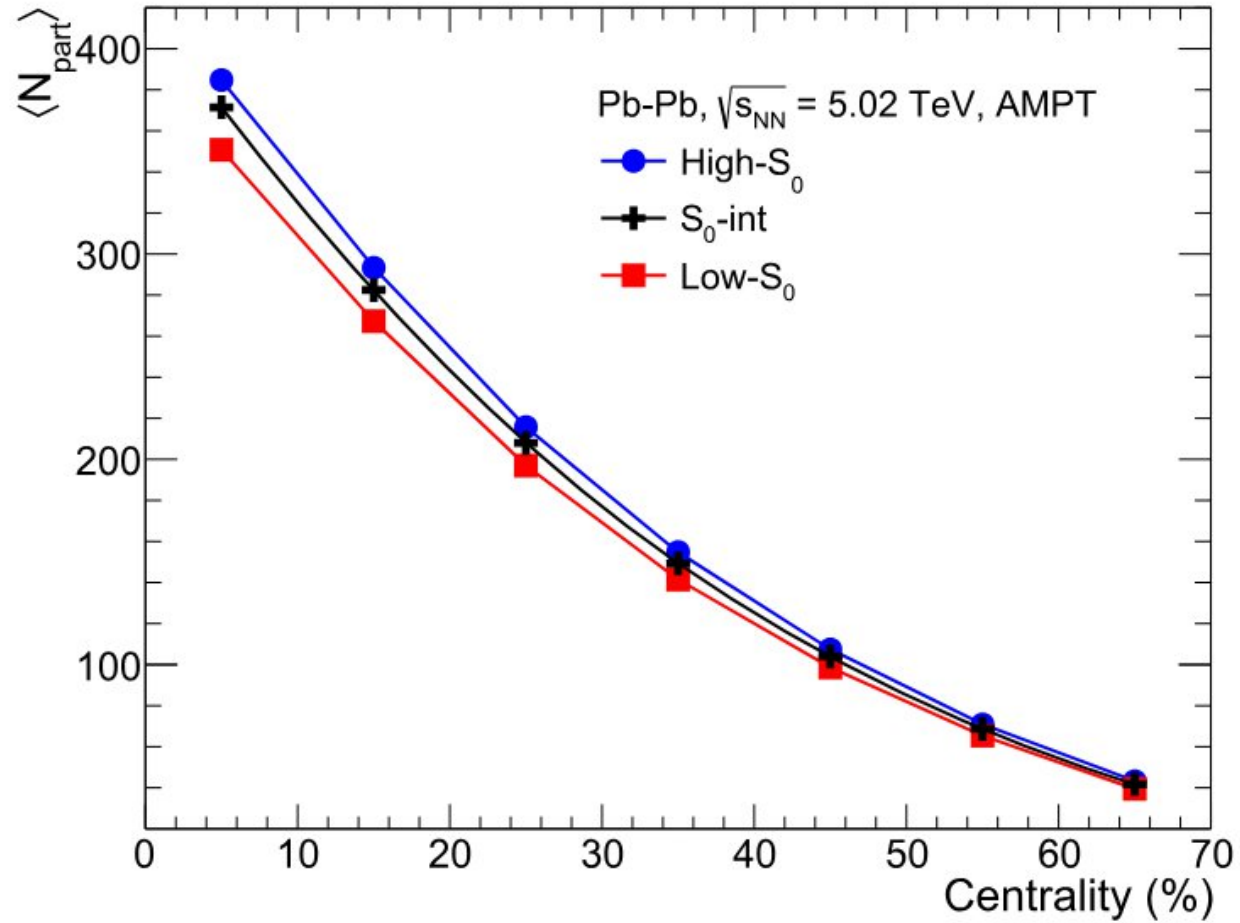


# Backup Slides





# $S_0$ versus $N_{\text{part}}$



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