



AWAKE Rb-Workshop

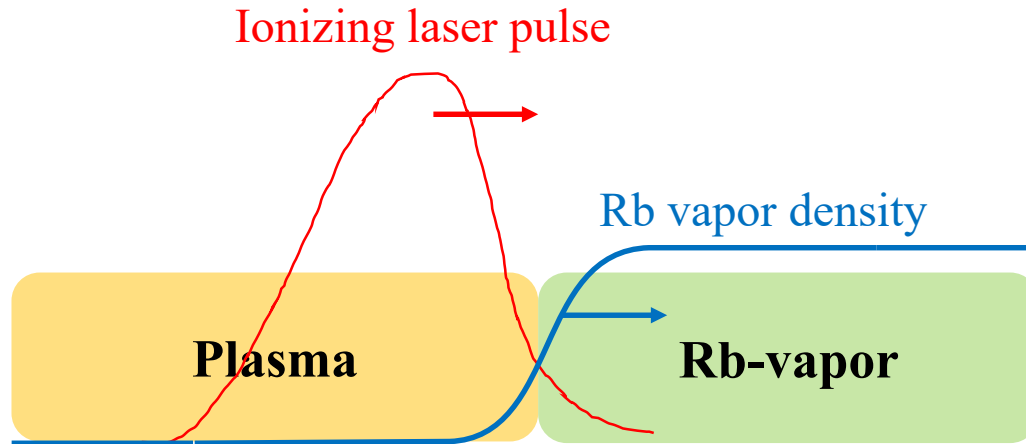
05-05-2017

Some vague proposals of Rb-vapor / plasma experiments

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- **Longitudinal Rb-cell diagnostic**
- **RF-wave reflection from plasma with moving ionization front boundary**
- **Self-focusing / guiding / trapping next to a resonant absorption line**

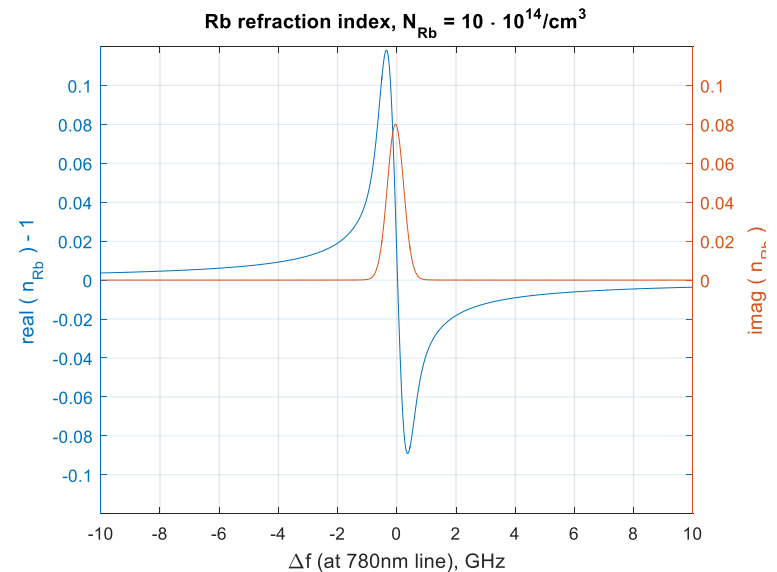
Rb vapour / plasma



phase velocity $\frac{V_{ph}}{c} = \frac{1}{\sqrt{1 - \left(\frac{f_p}{f_L}\right)^2}} \approx 1 + 2\left(\frac{f_p}{f_L}\right)^2$

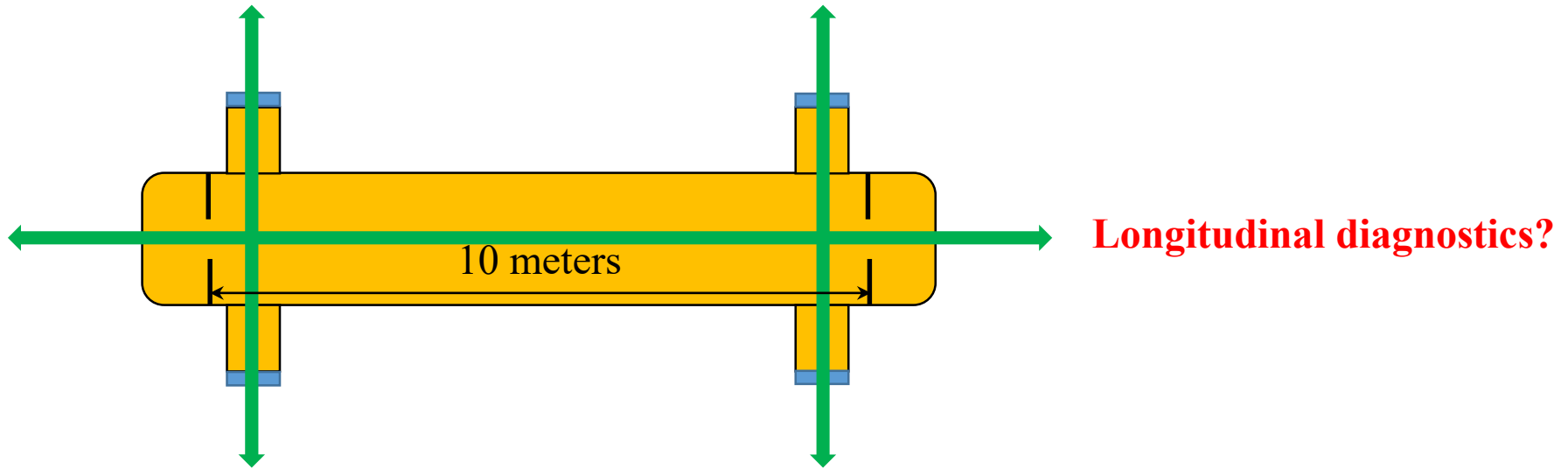
group velocity $\frac{V_{gr}}{c} = \sqrt{1 - \left(\frac{f_p}{f_L}\right)^2} \approx 1 - 2\left(\frac{f_p}{f_L}\right)^2$

$$f_p [GHz] \approx 90 \sqrt{N_{Rb} [10^{14} / cm^3]}$$

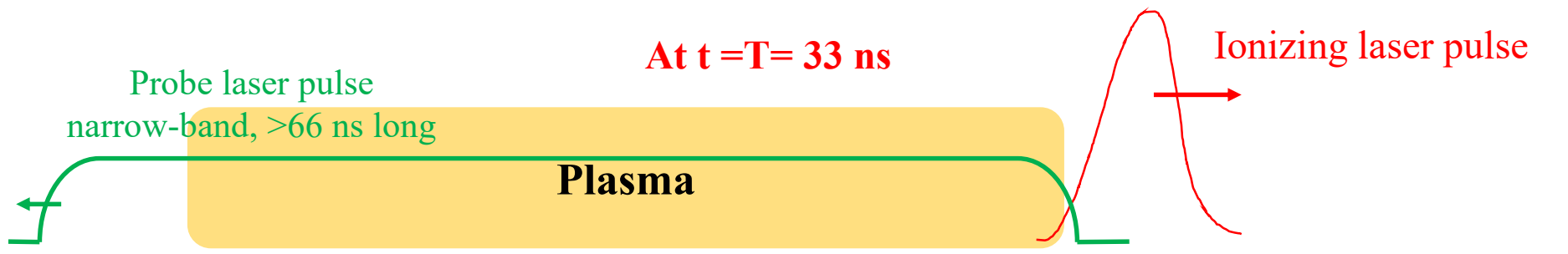
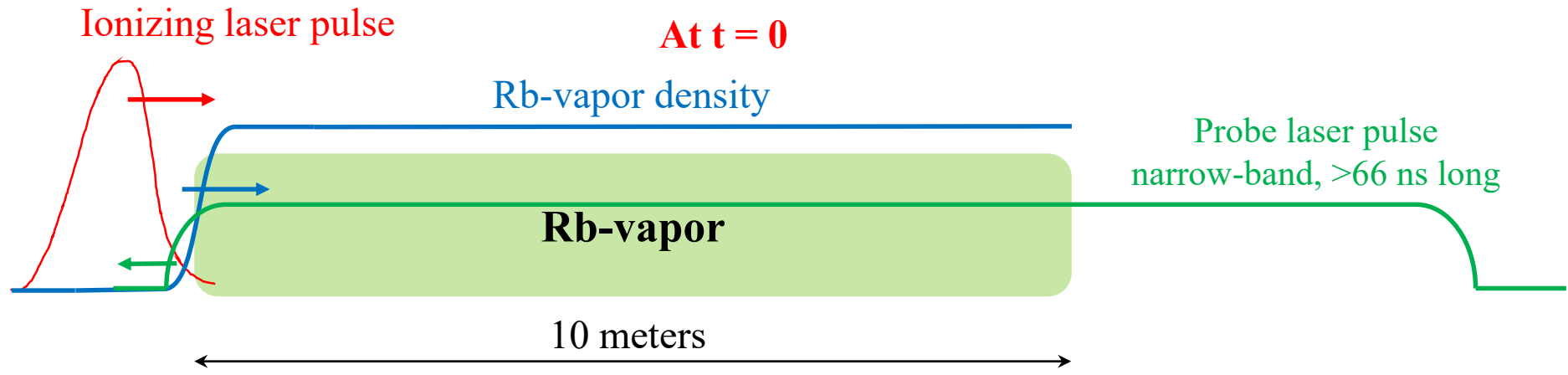


Rb vapour cell – possible probe directions

Two pairs of transverse view-ports:
Rb-density diagnostics, Schlieren imaging etc.



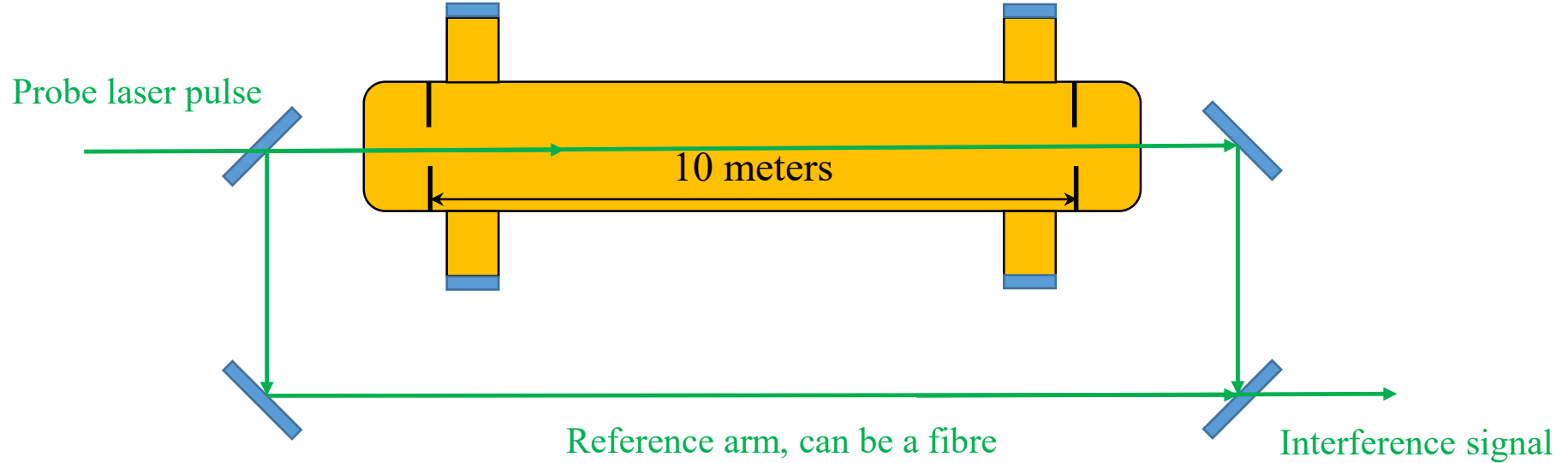
- Rb vapor cell → $N_{\text{Rb}} = 1 \div 10 \cdot 10^{14} / \text{cm}^3$
- Ti:Sa powerful laser → for ionization
- Narrow-band (<10MHz) tunable laser → for probing / diagnostics



Every slice of a long probing pulse (>66ns) is moving partially in Rb vapour and partially in plasma. Depending of the delay within the probing pulse every slice experience a different phase advance $\Delta\varphi(t)$

$$\Delta\varphi(t) = k_0 L(t / 2T + n_{Rb}(\omega)(1 - t / 2T))$$

Rb cell longitudinal diagnostic

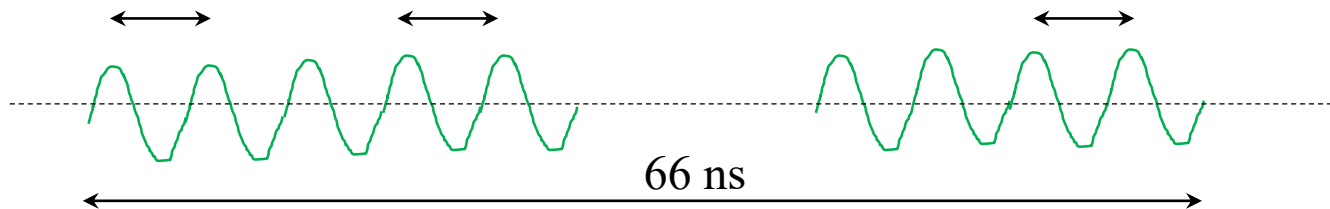


Interference signal is detected by a photo-diode on a fast scope

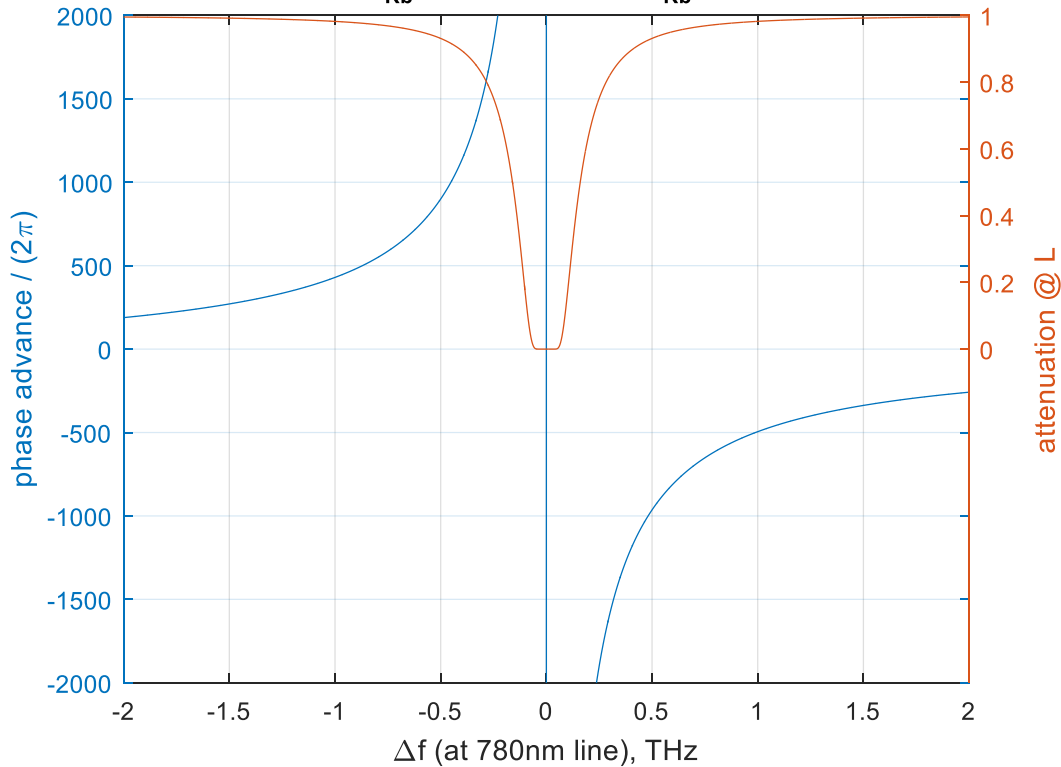
Interference signal is proportional to: $\sim \cos(k_0 L(t/2T + n_{Rb}(\omega)(1 - t/2T)))$

Number of fringes depends on detuning, fringes within 66ns has to be resolved by a fast scope.

Equidistance of the fringes tells about uniformity of $n_{Rb}(z)$



Rb cell, $N_{\text{Rb}} = 10 \cdot 10^{14} / \text{cm}^3$, $L_{\text{Rb}} = 1000 \text{ cm}$



Example:

Rb-density $N_{\text{Rb}} = 10^{15} / \text{cm}^3$, $L_{\text{Rb}} = 10 \text{ m}$

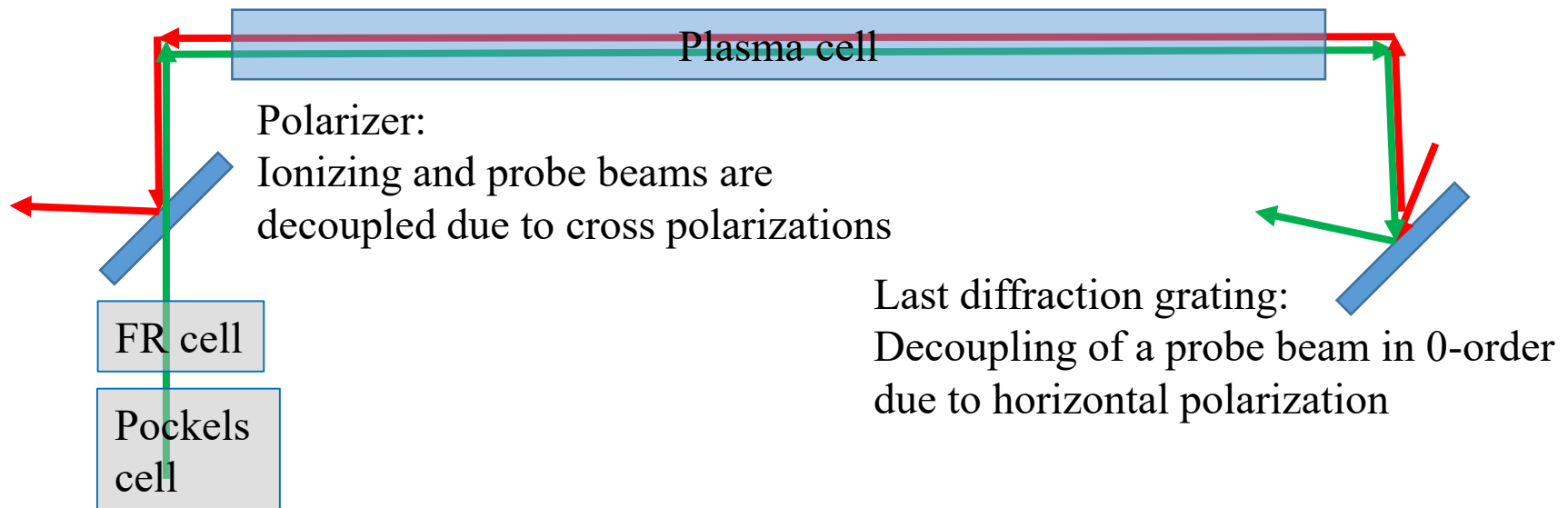
Low attenuation $\Delta f > 500 \text{ GHz}$ ($\sim 1 \text{ nm}$)

Phase advance $< 2\pi \cdot 800$

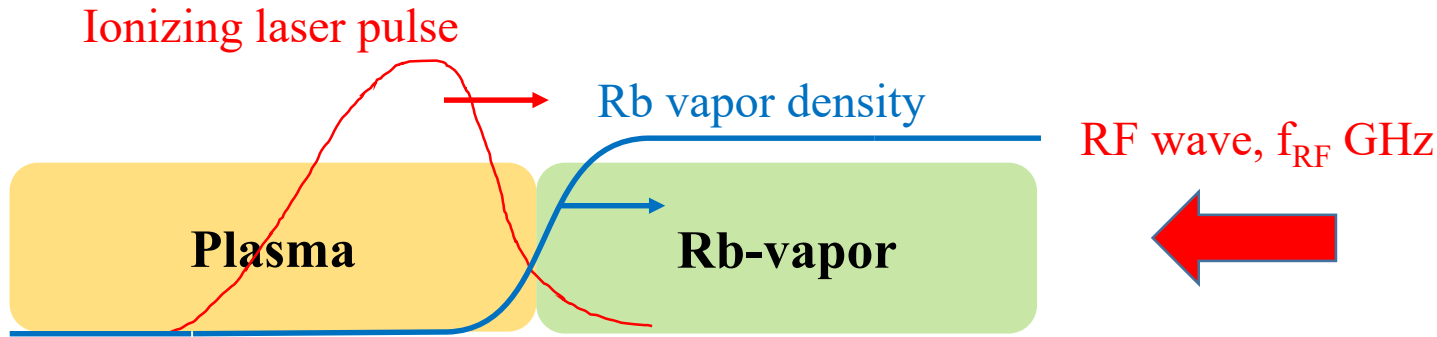
- Given typical resolution of a fast scope (20GHz/60Gsa) we can afford 200 interference fringes within 66ns, assuming 330ps/period
- So maximum phase advance is $\Delta\phi_{\text{max}} = 2\pi \cdot 200$, and thus detuning has to be $\sim 2 \text{ THz}$ ($\sim 4 \text{ nm}$)

Problems / solutions:

- Transverse overlap of ionizing and probe lasers
 - probe pulse mode degradation
- Coupling / decoupling of lasers
 - different polarizations
 - Bruster plates / polarizers
 - Zero-order grating reflections (inside a compressor)
 - Faraday isolators, Pockels cells etc.)



RF-wave reflection from an ionization front (a very old idea in general)



Ionization laser group velocity $\frac{V_{gr}}{c} = \sqrt{1 - \left(\frac{f_p}{f_L}\right)^2} \approx 1 - 2\left(\frac{f_p}{f_L}\right)^2$ $f_p [GHz] \approx 90\sqrt{N_{Rb} [10^{14} / cm^3]}$

Ionization front velocity equal (?) to group velocity of ionizing laser:

$f_p \sim 400 \text{ GHz}$ ($N_{Rb} = 2 \cdot 10^{15} / cm^3$), $f_L \sim 400 \text{ THz}$

$1 - V_{gr} / c \sim 2 \cdot 10^{-6} \quad \Leftrightarrow \quad \text{“equivalent Lorentz factor } \gamma \text{”} \sim 500$

What if seed an RF-wave counter propagating to an ionization front?

Boost to front frame: $f'_{RF} = f_{RF} \sqrt{\frac{1+\beta}{1-\beta}} = 1000 f_{RF}$, **plasma frequency stays the same ($N_e/m=inv$)!**

If $f'_{RF} < f_p$ (i.e. $f_{RF} < 0.4 \text{ GHz}$), then RF wave reflects from plasma boundary.

Boost back to lab frame: $f''_{RF} = f_{RF} \frac{1+\beta}{1-\beta} = 10^6 f_{RF} \sim 400 \text{ THz}$, in or example – **it is visible!**

RF-wave reflection from an ionization front.

Some practical limitations

Ionization front velocity:

$$1 - \frac{V_{gr}}{c} = 1 - \beta \approx 2 \left(\frac{f_p}{f_L} \right)^2$$

RF frequency in the moving front frame:
(must be less than f_p for high reflectivity)

$$f'_{RF} = f_{RF} \sqrt{\frac{2}{1-\beta}} = f_{RF} \left(\frac{f_L}{f_p} \right) \quad \text{must be } < f_p$$

Up-shifted RF frequency in lab frame:

$$f''_{RF} = f_{RF} \frac{2}{1-\beta} = f_{RF} \left(\frac{f_L}{f_p} \right)^2 < f_L$$

In reasonable conditions ($f'_{RF} < f_p$), up-shifted frequency always less than laser frequency!

Up-shifted frequency can be in near/far IR or THz range

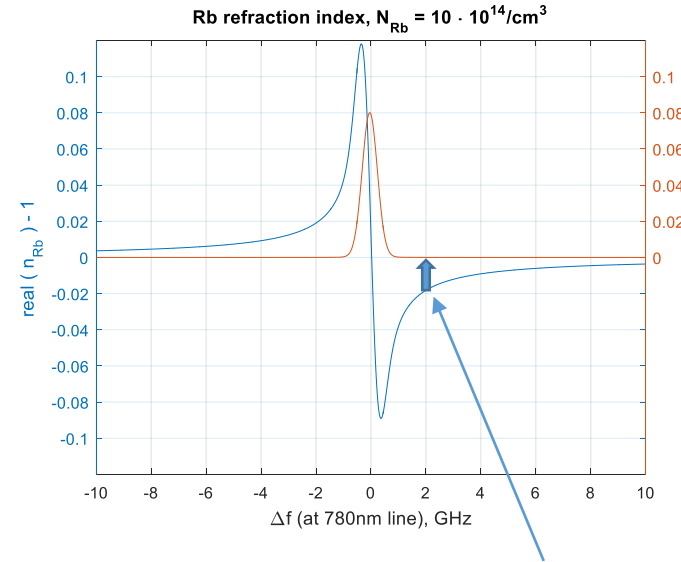
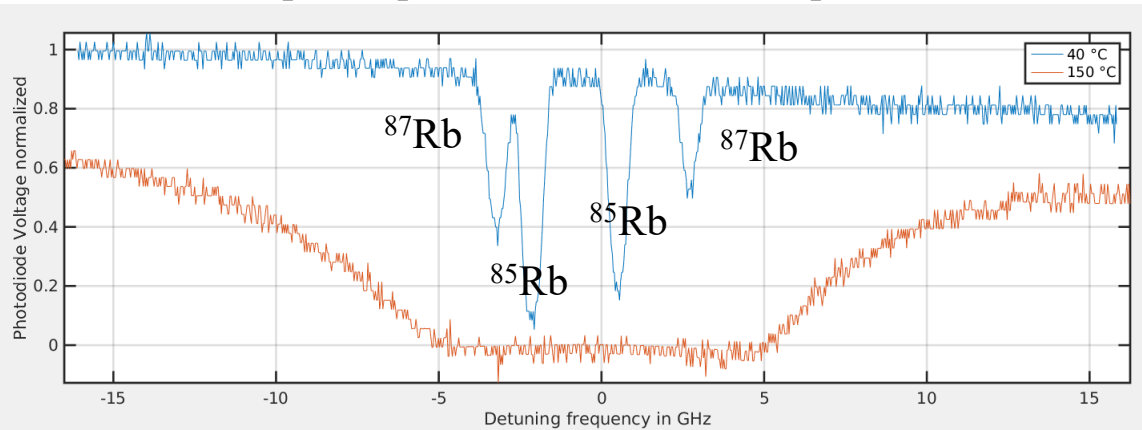
Sacrificing the efficiency, $f'_{RF} < 10f_p$, then $f''_{RF} < 10f_L$, it is potentially deep UV.

Wavelength of reflected visible light being measured by time resolving spectrometer gives information about $f_p(t = z / c)$

Self-focusing next to a resonant absorption line.

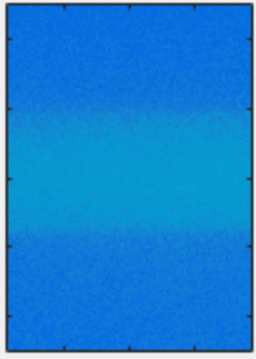
Known since 1966.

Absorption spectrum at different temperature



To the left of the line
- defocusing

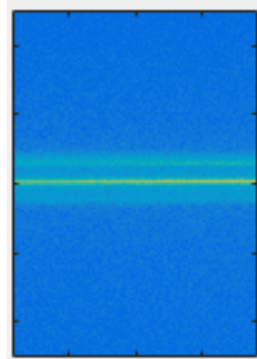
-1.9 GHz



Absolute values of detuning
are not very accurate

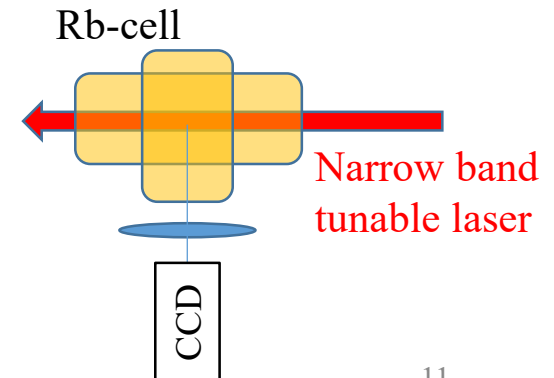
To the right of the line
- focusing

4.7 GHz



Courtesy of Anna Bachmann

Laser beam being properly tuned saturate an absorption and tends to self focus at positive frequency detuning



Equations are very similar to those for laser gain (Frantz-Nodvik)

$$N_0 = N_1 + N_2 = \text{const}$$

$$\frac{\partial N_1(x,t)}{\partial t} = -(N_1 - N_2)\sigma \frac{I(x,t)}{h\nu} + \frac{N_2}{\tau}$$

$$\frac{\partial N_2(x,t)}{\partial t} = +(N_1 - N_2)\sigma \frac{I(x,t)}{h\nu} - \frac{N_2}{\tau}$$

$$\frac{\partial I(x,t)}{\partial x} + \frac{1}{c} \frac{\partial I(x,t)}{\partial t} = -(N_1 - N_2)\sigma I(x,t)$$

$$A = \sigma(N_1 - N_2)$$

$$\begin{cases} \frac{\partial A(z,t)}{\partial t} = -2A \frac{I}{E_s} \\ \frac{\partial I(z,t)}{\partial z} = -AI \end{cases}$$

With a solution:

$$A(z,t) = \frac{A_0(z) \exp\left(\int_{-\infty}^z A_0(x) dx\right)}{\exp\left(\frac{2}{E_s} \int_{-\infty}^t I_0(x) dx\right) + \exp\left(\int_{-\infty}^z A_0(x) dx\right) - 1}$$

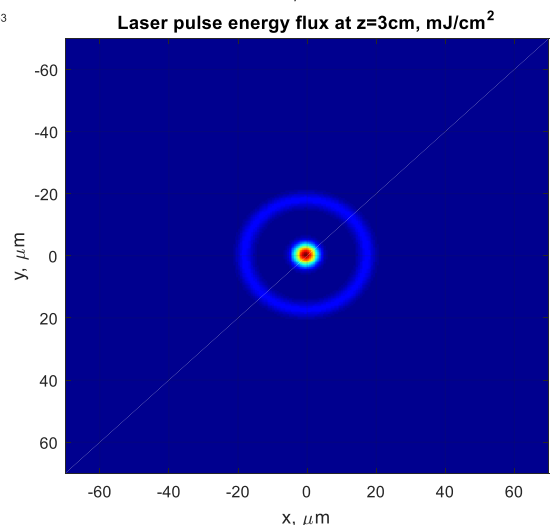
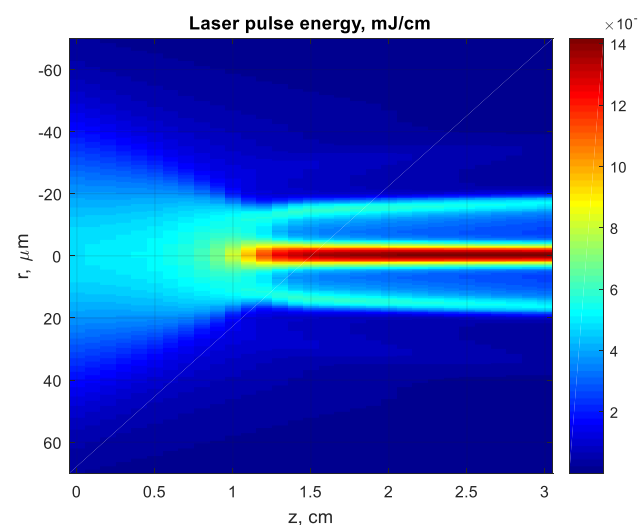
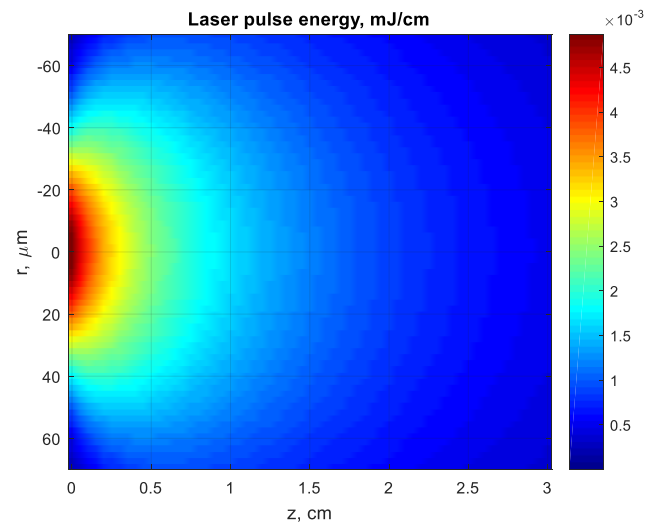
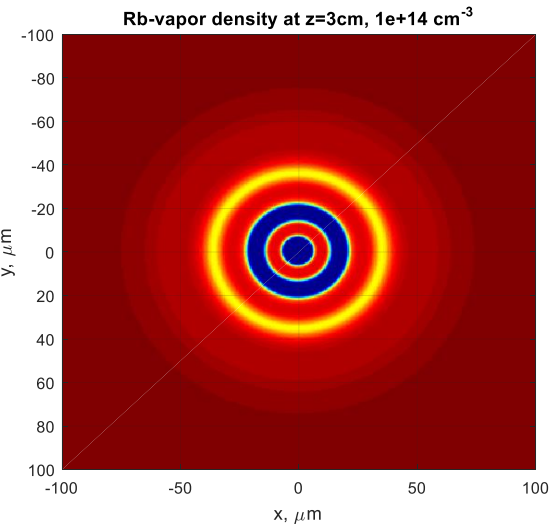
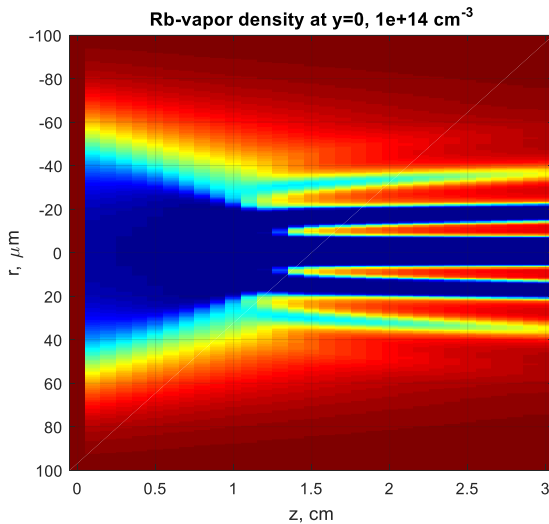
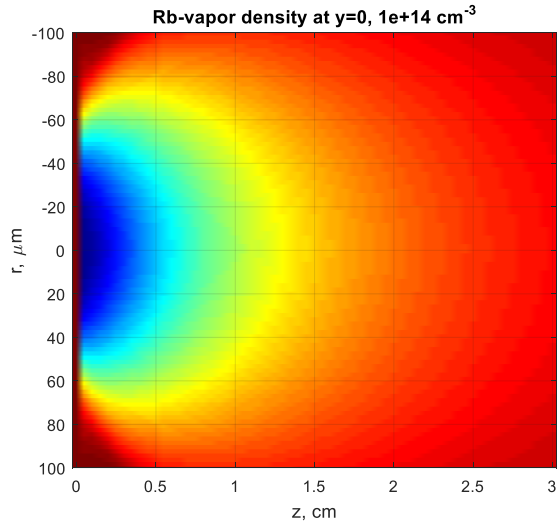
$$I(z,t) = \frac{I_0(t) \exp\left(\frac{2}{E_s} \int_{-\infty}^t I_0(x) dx\right)}{\exp\left(\frac{2}{E_s} \int_{-\infty}^t I_0(x) dx\right) + \exp\left(\int_{-\infty}^z A_0(x) dx\right) - 1}$$

Split-step algorithm, for each slice dz in z -direction:

- Calculate $\Delta N(x,y,z)$ and $I(x,y,t)$ according to Frantz-Nodvik
 - cross-section and saturation energy must be consistent with $\text{Im}\{\chi\}$ at given frequency
 - $\sigma = 2k_0 \text{Im}\{\chi\}$, $E_{sat} = h\nu / \sigma$
- Use ΔN to calculate the phase acquired on a z -slice: $U(x, y, t) = \sqrt{I(x, y, t)} e^{ik_0 dz \Delta N \text{Re}\{\chi\}/2}$
- Do the step propagation: $U(x, y, t) = \text{FFT}^{-1}\left(\text{FFT}(U) \cdot e^{idz(k_x^2+k_y^2)/2k_0}\right)$

$\Delta f = -5 \text{ GHz}$
defocusing

$\Delta f = +5 \text{ GHz}$
focusing





Thank you for your attention!