

Annual meeting of Academia Europaea
Budapest
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Beauty and the Quantum

Takaharu Otsuka

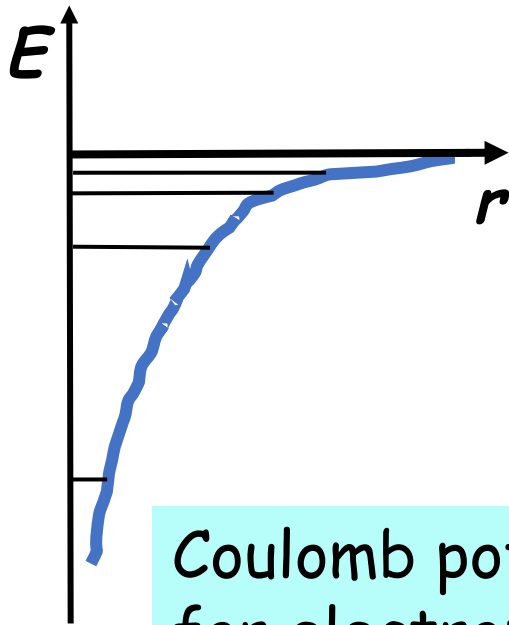


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Noritaka Shimizu (CNS, Tokyo), Takashi Abe (Tokyo)

This work has been supported by MEXT and JICFuS as a priority issue (Elucidation of the fundamental laws and evolution of the universe) to be tackled by using Post 'K' Computer

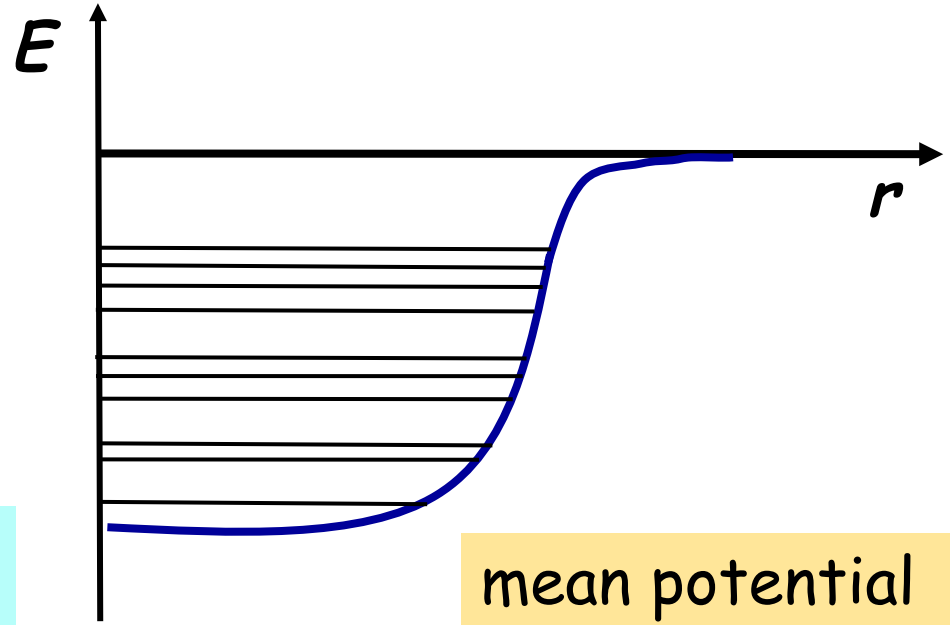
Single-particle states - starting point -

electrons in **atom**

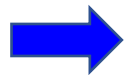


Coulomb potential
for electrons
by the nucleus

nucleons (protons and neutrons)
in atomic **nucleus**

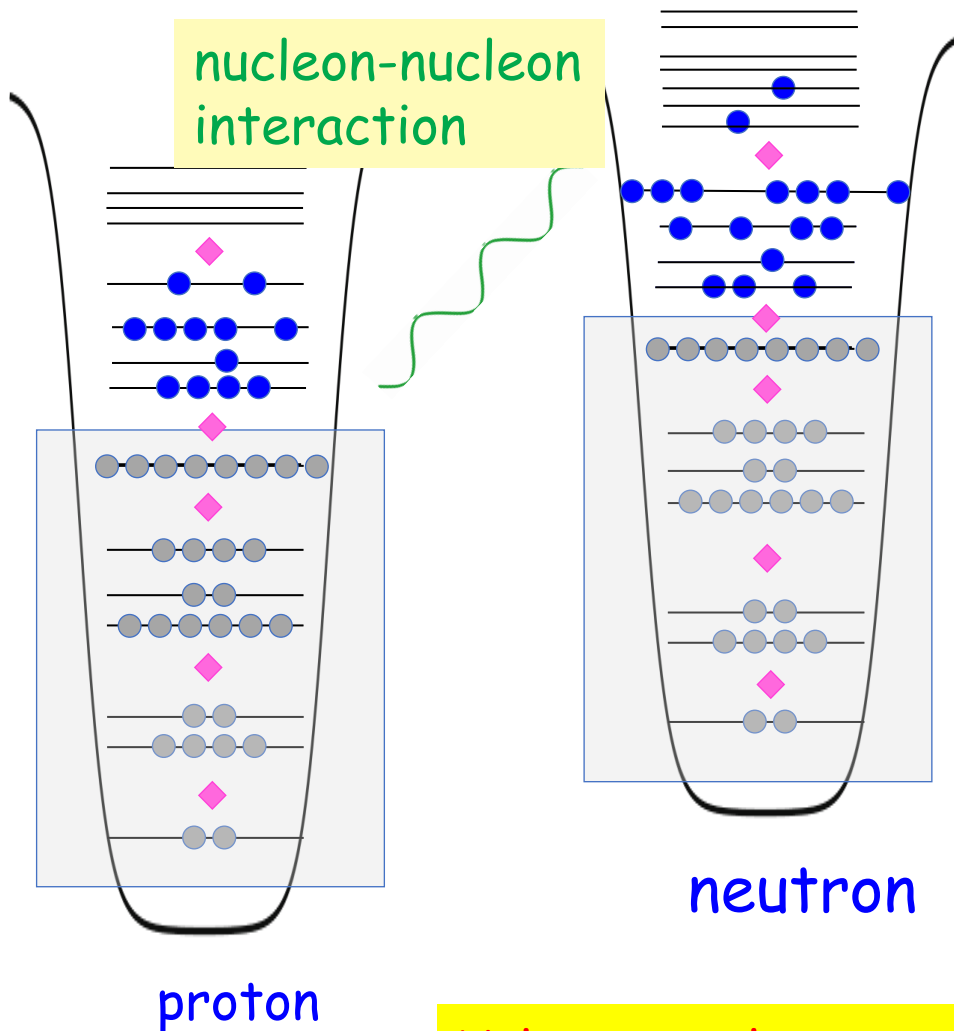


mean potential
for nucleons
by the nucleons



magic numbers, shell structure, etc.

shell structure and nucleon-nucleon interaction



Protons and neutrons are orbiting in the mean potential like a "vase"

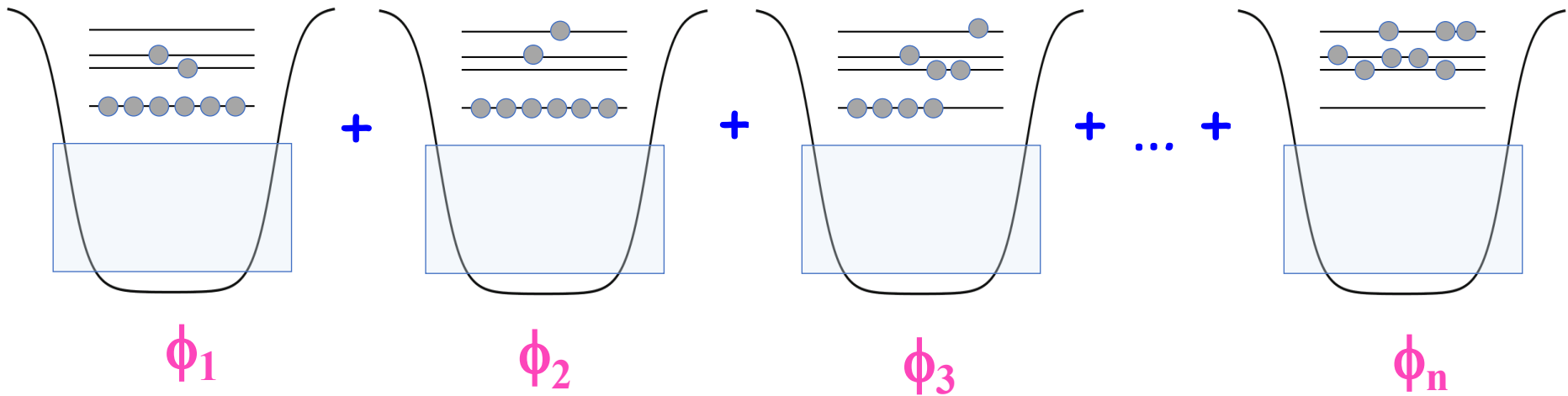
→ single-particle states

Lower orbits form the inert core (or closed shell) (shaded parts in the figure)

Upper orbits are only partially occupied (valence orbits and nucleons).

◆ shell gap

Valence nucleons are the major source of nuclear dynamics at low excitation energy.



The superposition of various configurations is fixed as an Eigenvalue problem

$$\mathbf{H} \Psi = \mathbf{E} \Psi$$

$$\Psi = c_1 \phi_1 + c_2 \phi_2 + c_3 \phi_3 + \dots$$

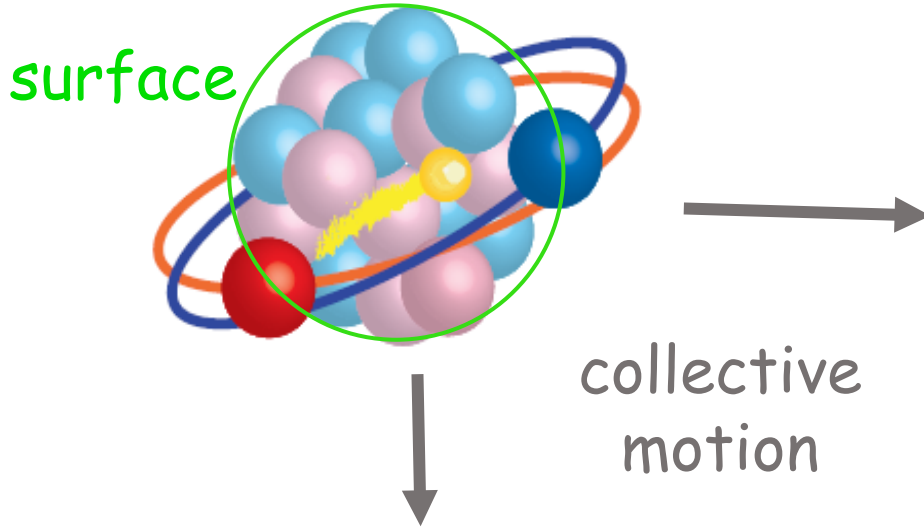
c_i probability amplitudes

Solved by Monte Carlo Shell Model (like CI calc.) on supercomputers, for instance, K computer in Kobe.

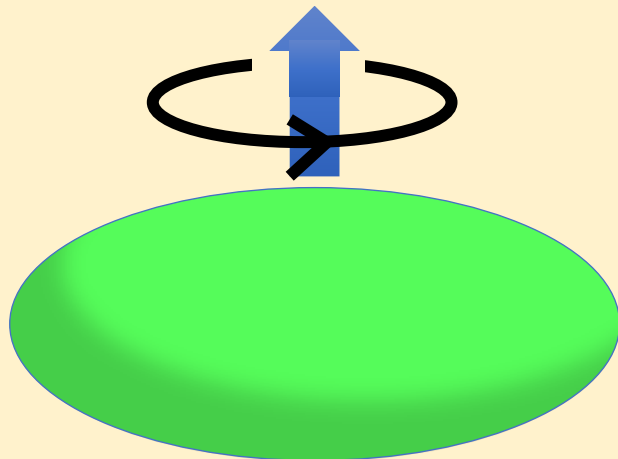
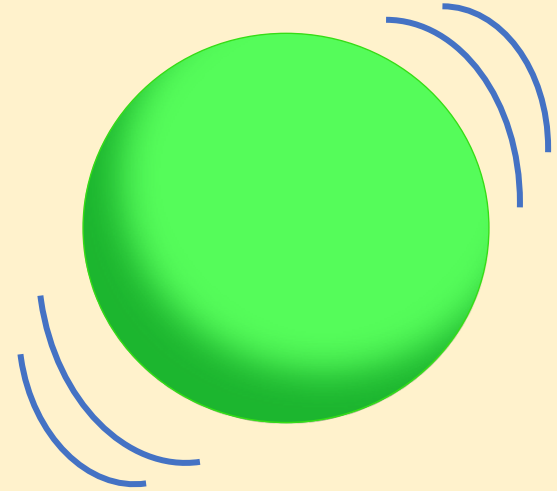


What properties appear from such correlations ?

Assembly of
protons and neutrons



Vibration between
sphere and ellipsoid



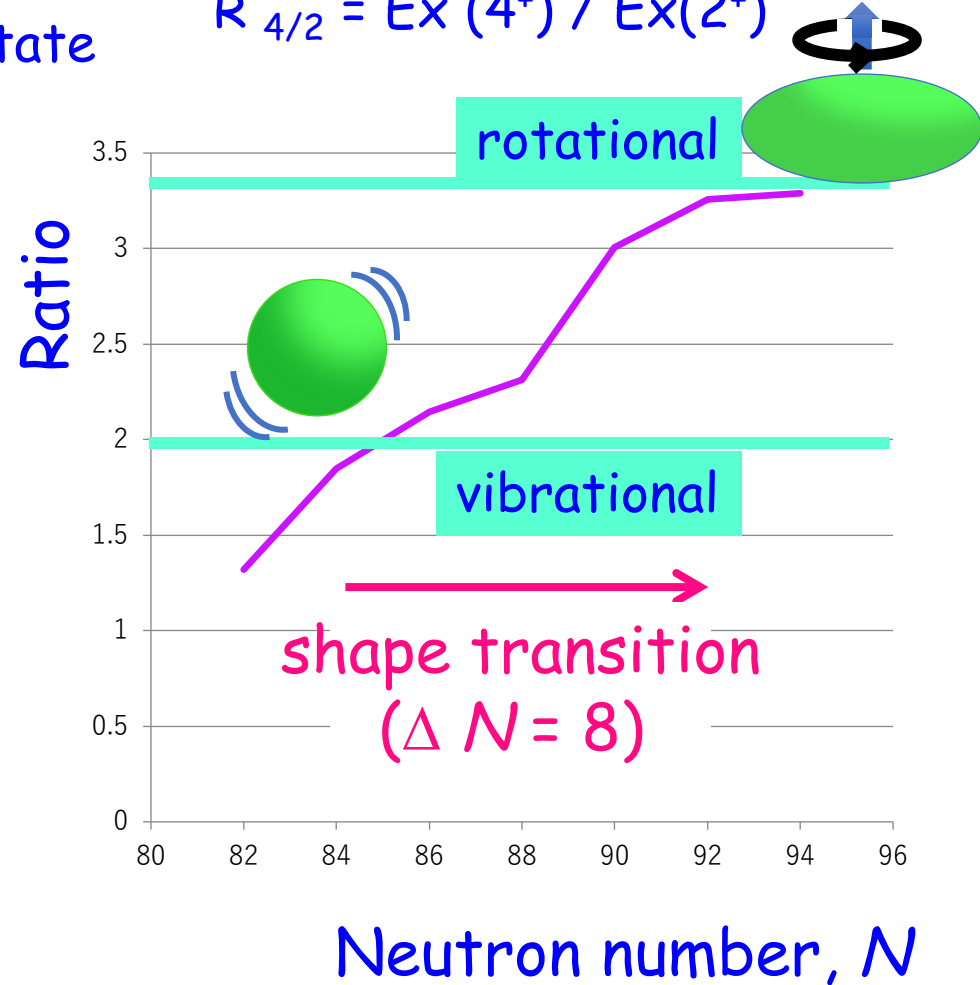
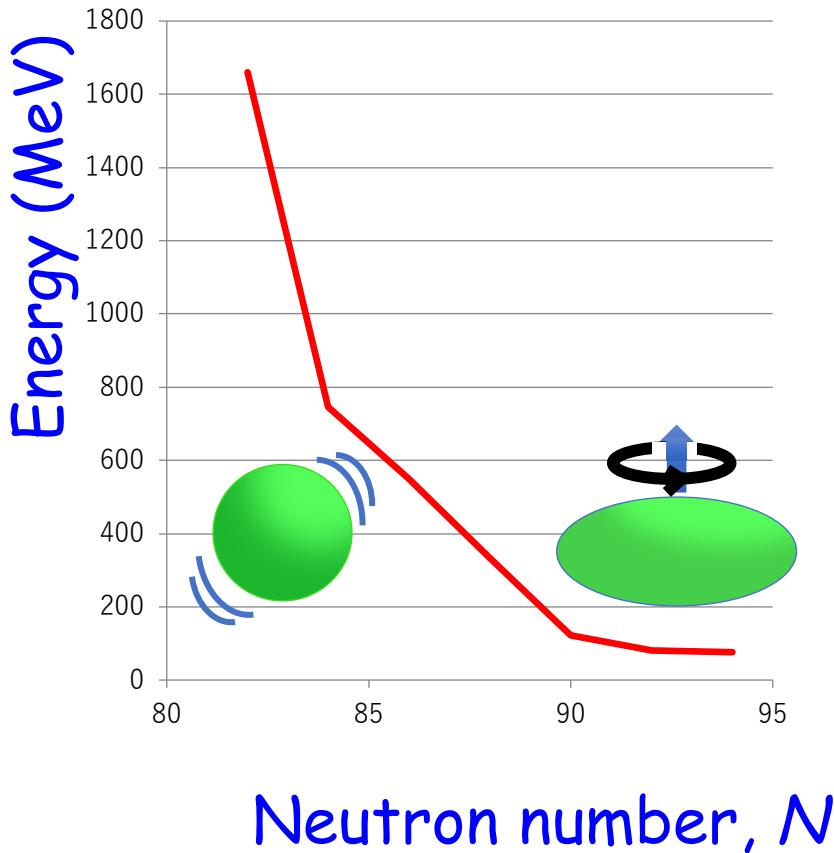
Rigid Ellipsoidal
Deformation
and its Rotation

What properties appear ? : Shape change as a function of N (or Z)

2^+ and 4^+ level properties of Sm ($Z=62$) isotopes

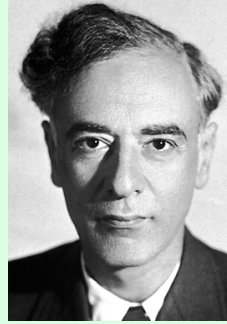
Ex (2^+) :
excitation energy of first 2^+ state

$$R_{4/2} = \text{Ex}(4^+) / \text{Ex}(2^+)$$



Atomic nucleus is a quantum Fermi liquid :

The nucleus is composed of almost *free nucleons* interacting weakly via residual forces in a (solid) (mean) potential like a *solid vase*.

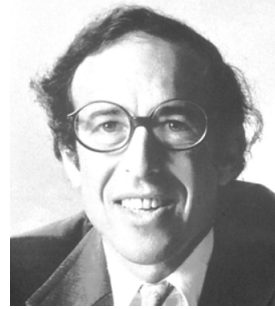


Landau

The shape of atomic nucleus can be deformed dynamically (vibration) or statically (rotation) as collective modes (motions).



A. Bohr

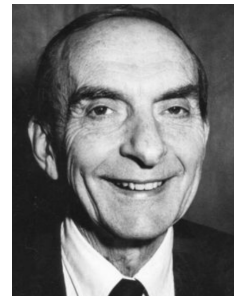


Mottelson

T. Schaefer, Fermi Liquid theory: A brief survey in memory of Gerald E. Brown, NPA 2014)

One of Gerry's main scientific pursuits was to understand the nuclear few and many-body problem in terms of microscopic theories based on the measured two and three-nucleon forces. One of the challenges of this program is to understand how the observed single-particle aspects of finite nuclei, in particular shell structure and the presence of excited levels which carry the quantum numbers of single particle states, can be reconciled with the strong nucleon-nucleon force, and **how single particle states can coexist with collective modes**. A natural framework for addressing these questions is the **Landau theory of Fermi liquids**. **Landau Fermi liquid theory**

G.E. Brown



Quest for Quantum Phase Transition: Shapes of Zr isotopes by Monte Carlo Shell Model

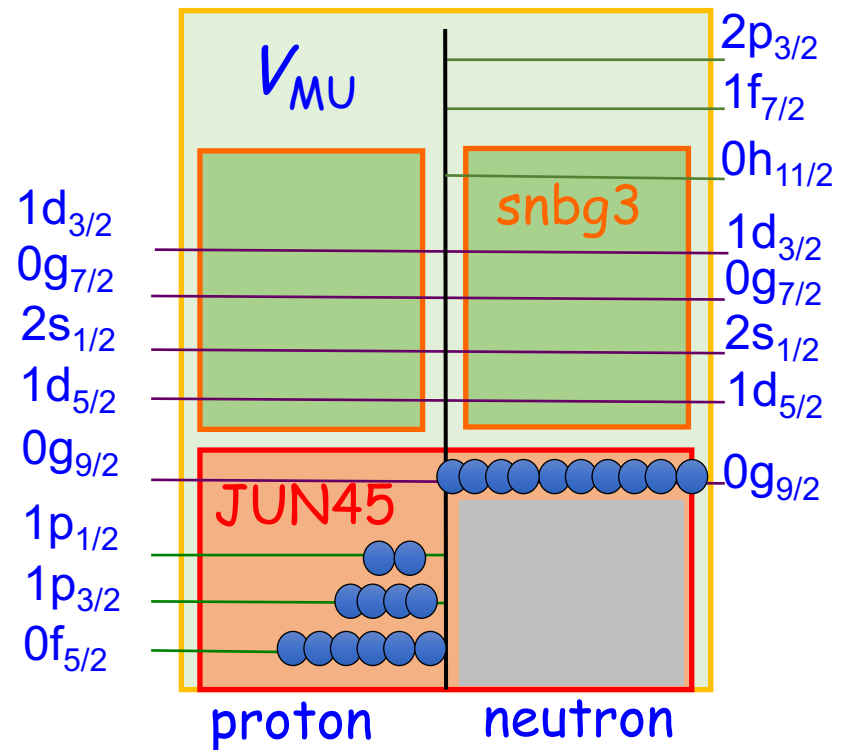
- Effective interaction:
JUN45 + **snbg3** + V_{MU}

known effective interactions

+ minor fit for a part of
T=1 TBME's

Nucleons are excited fully
within this model space
(no truncation)

We performed **Monte Carlo Shell Model (MCSM)** calculations, where the largest case corresponds to the diagonalization of 3.7×10^{23} **dimension** matrix.



^{56}Ni

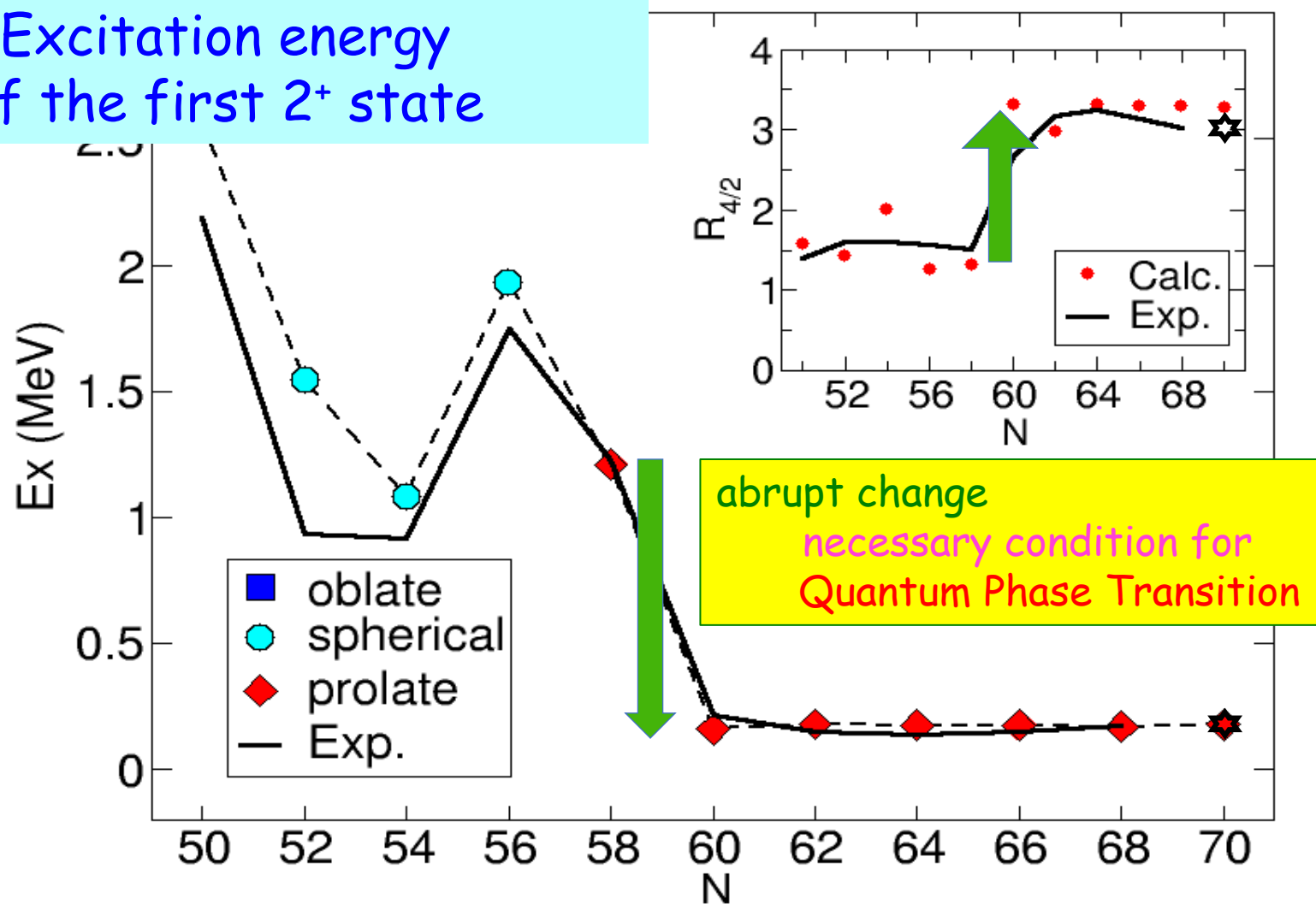
Togashi, Tsunoda, TO *et al.* PRL
117, 172502 (2016)



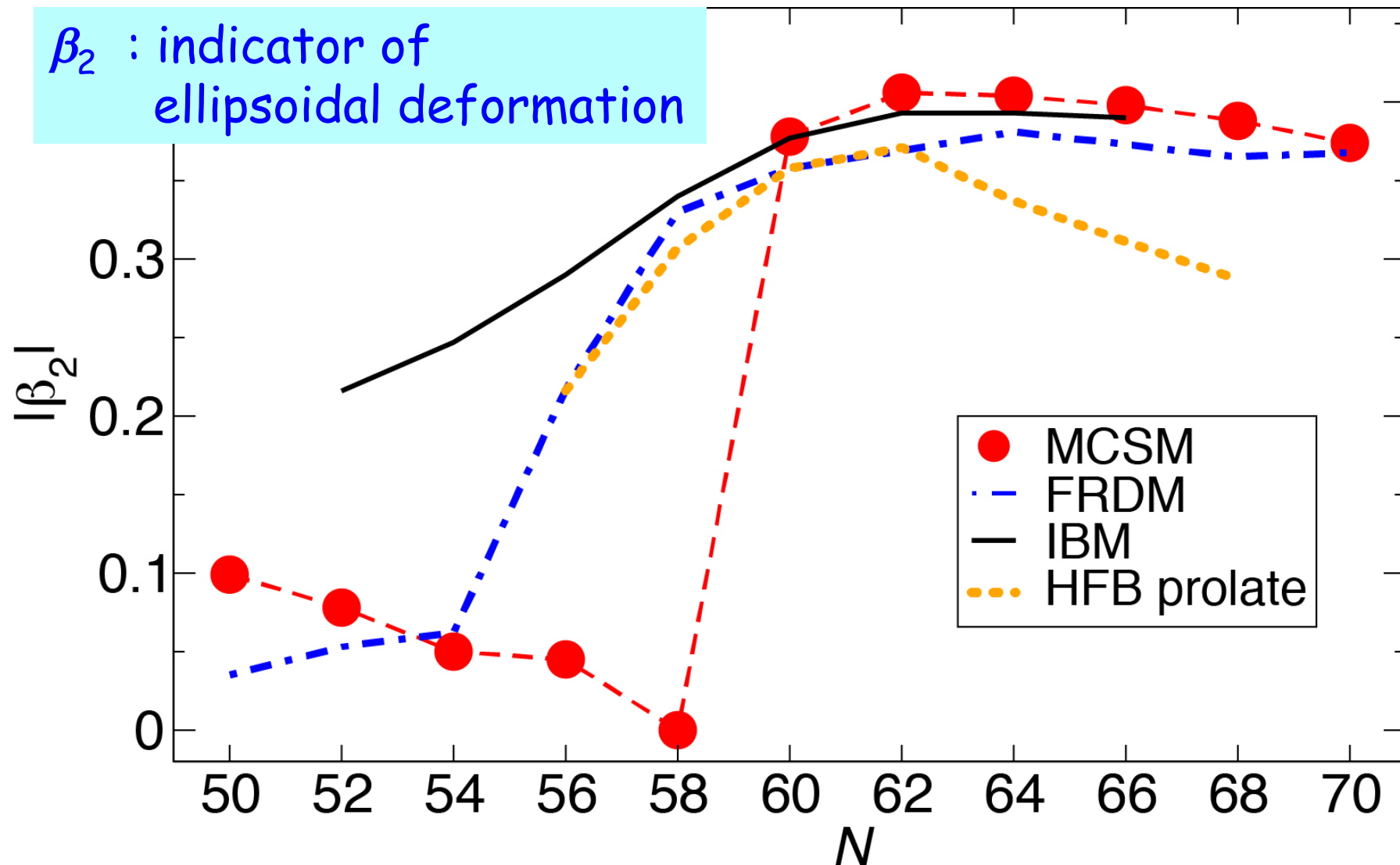
Quantum Phase Transition in the Shape of Zr isotopes

Tomoaki Togashi,¹ Yusuke Tsunoda,¹ Takaharu Otsuka,^{1,2,3,4} and Noritaka Shimizu¹

Excitation energy
of the first 2^+ state



β_2 : indicator of
ellipsoidal deformation



FRDM: S. Moeller et al. *At. Data Nucl. Data Tables* 59, 185 (1995).

IBM: M. Boyukata et al. *J. Phys. G* 37, 105102 (2010).

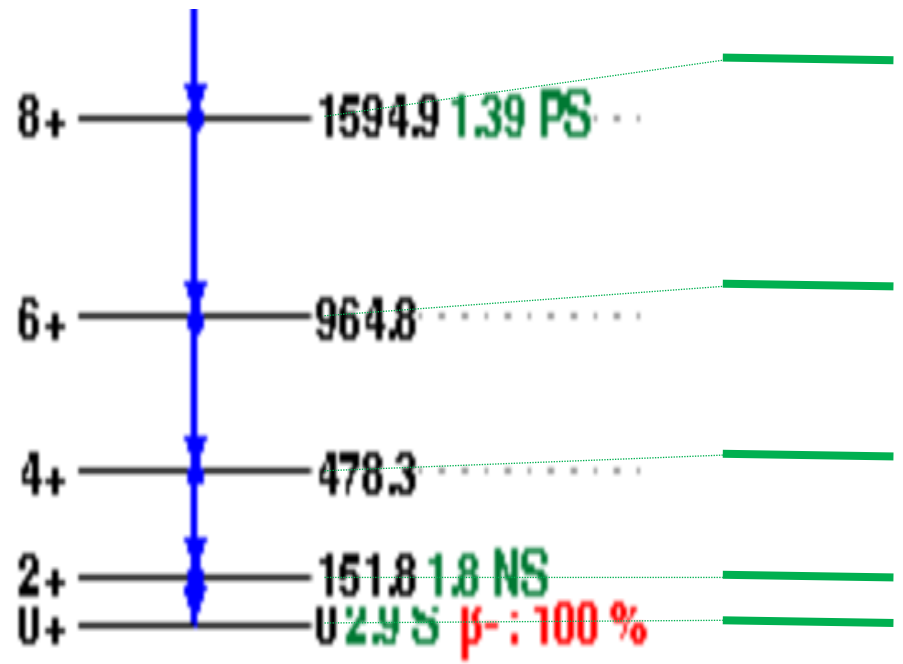
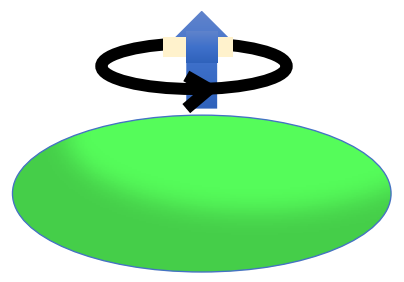
HFB: R. Rodriguez-Guzman et al. *Phys. Lett. B* 691, 202 (2010).¹⁰

Beautiful regularities arise.

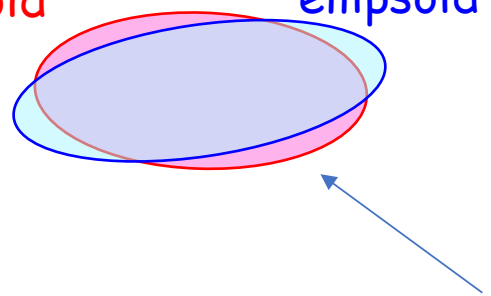
Example :
Rotational band in ^{102}Zr

exp.

Quantized rigid rotor
 $E_x(J) \propto J(J+1)$



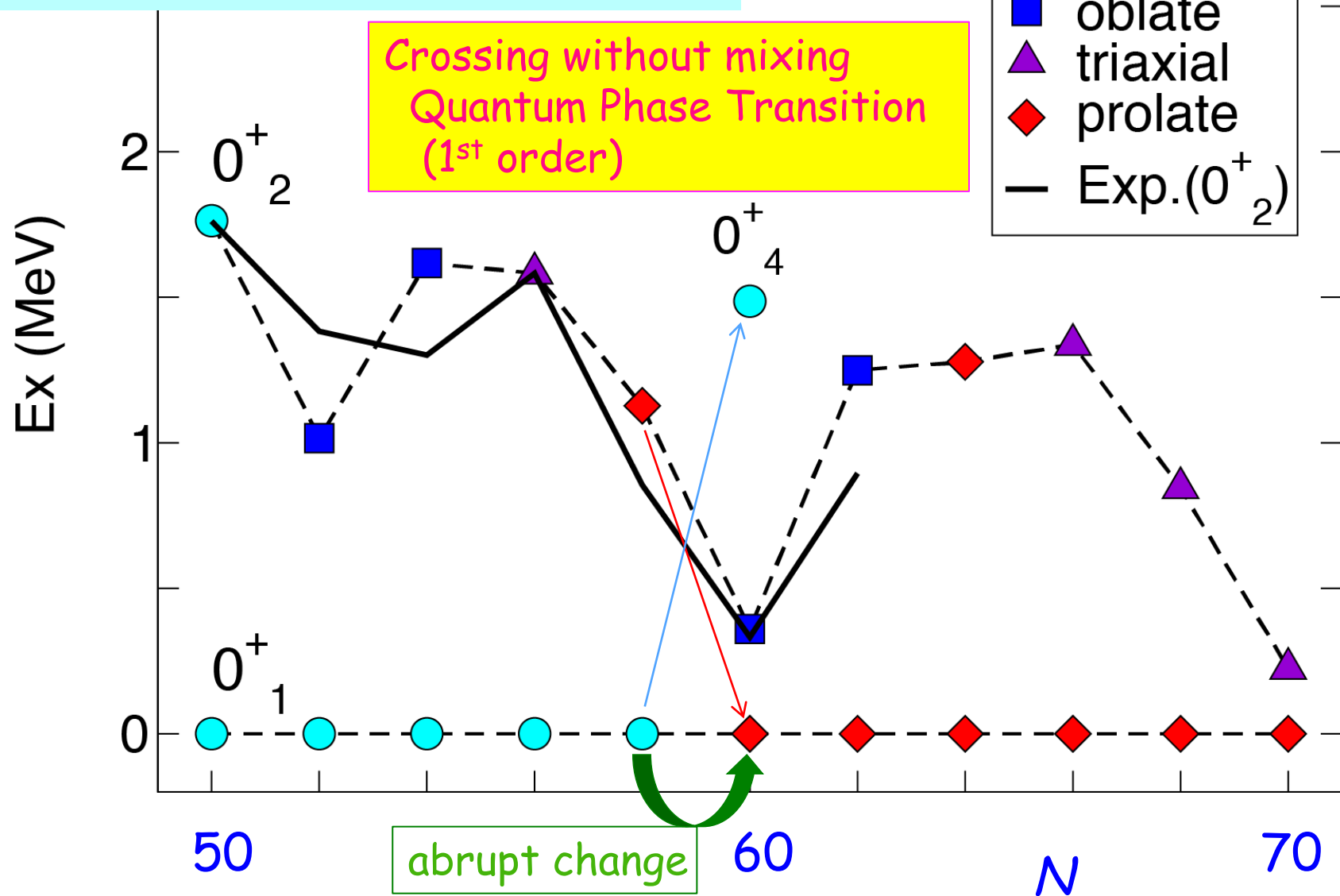
neutron ellipsoid
proton ellipsoid

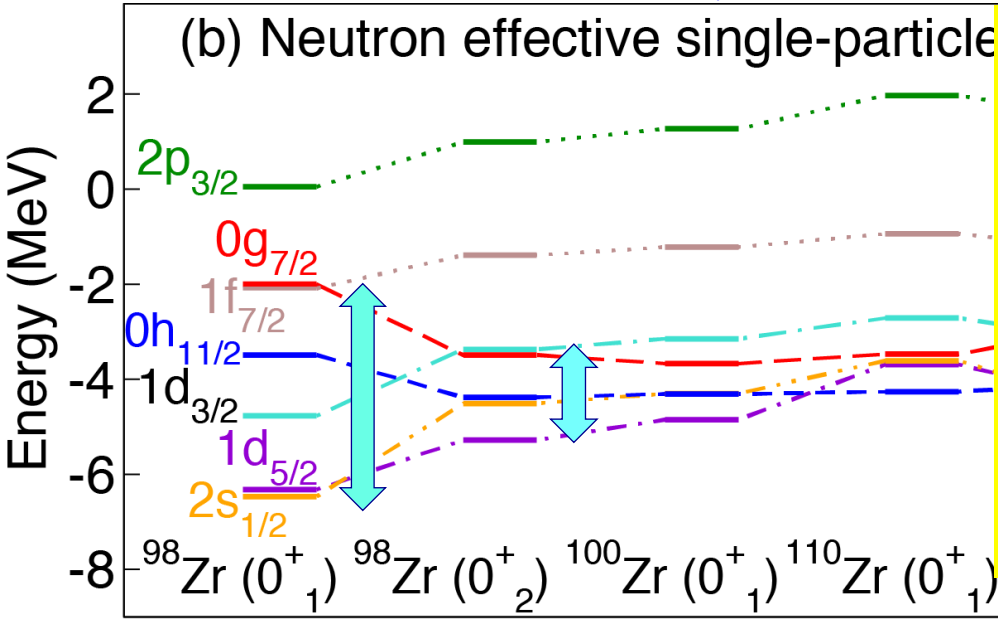
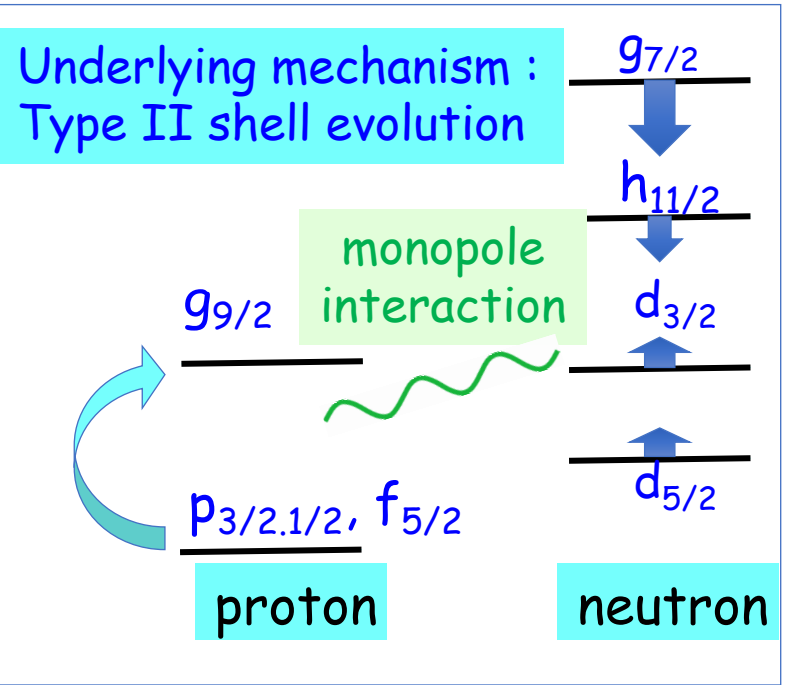
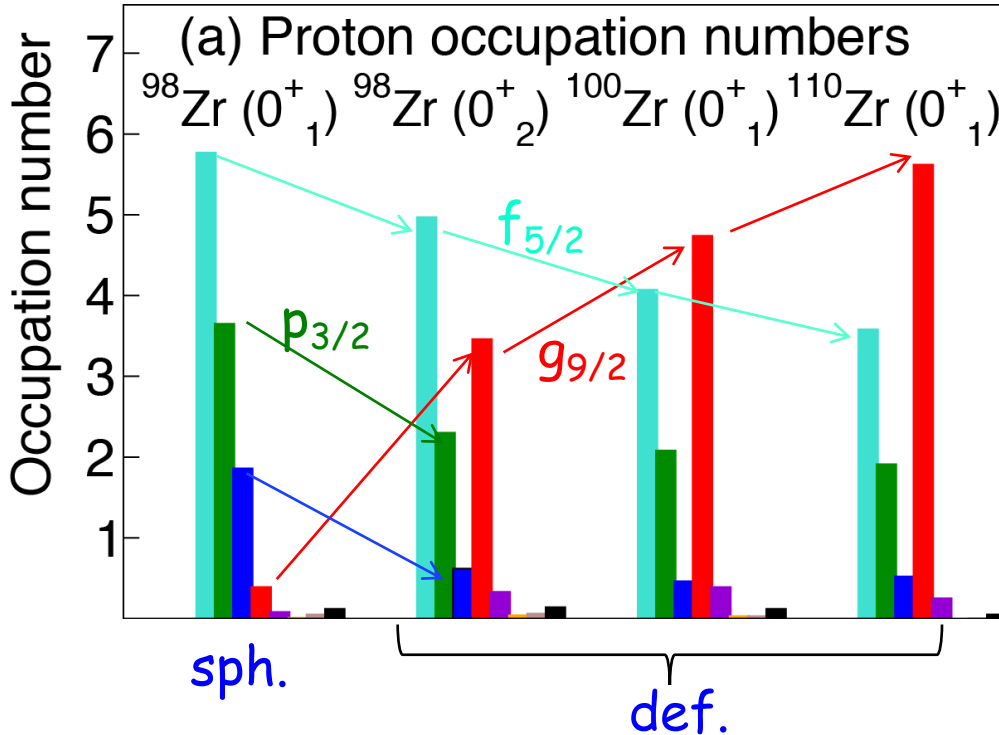


Origin : proton-neutron quadrupole-quadrupole interaction

This is not the whole story !

Excitation energy of the first and second 0^+ states





Relevant neutron single-particle levels get closer as a combined effect of nuclear forces (tensor and central) and particular configurations.

The resistance power against deformation is then reduced.

Large difference in ESPEs and configurations \rightarrow crossing w/o mixing

This is one of the cases of

Quantum Self Organization

$$\text{deformation} = \frac{\text{quadrupole force}}{\text{resistance power}}$$

resistance power ← pairing force

↑ single-particle energies

Atomic nuclei can “organize” their *single-particle energies* by taking particular configurations of protons and neutrons *optimized for each eigenstate*, thanks to orbit-dependences of *monopole components of nuclear forces* (e.g., tensor force).

→ an enhancement of Jahn-Teller effect.

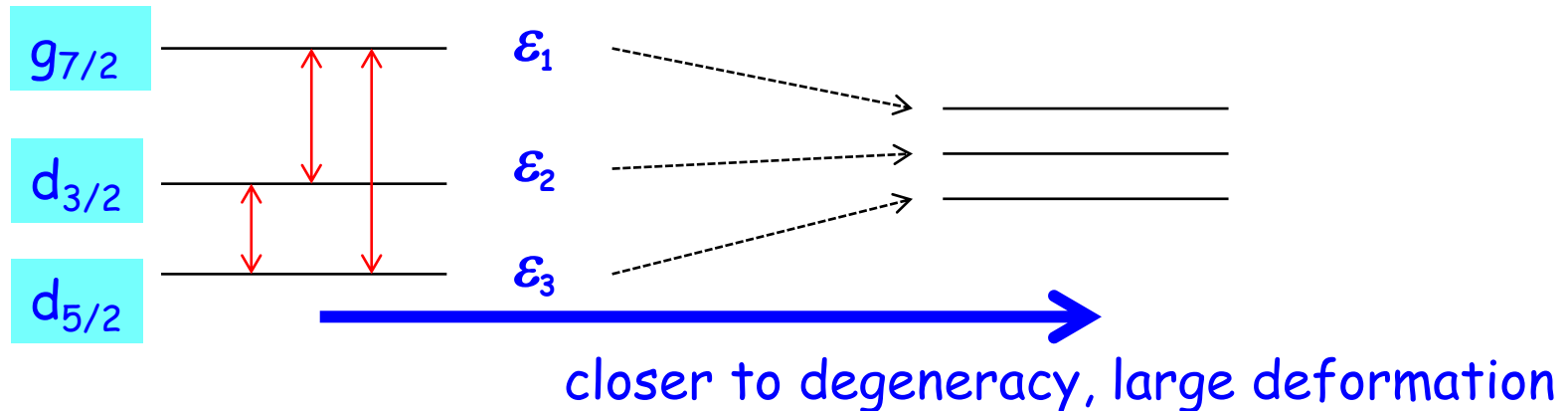
Reminder : Jahn -Teller effect for nuclear deformation

(Self-consistent) quadrupole deformed field $\propto Y_{2,0}(\theta,\phi)$
mixes the orbits below

$$\Psi (J_z=1/2) = c_1 |g_{7/2}; j_z=1/2\rangle + c_2 |d_{3/2}; j_z=1/2\rangle + c_3 |d_{5/2}; j_z=1/2\rangle$$

stronger mixing = larger quadrupole deformation

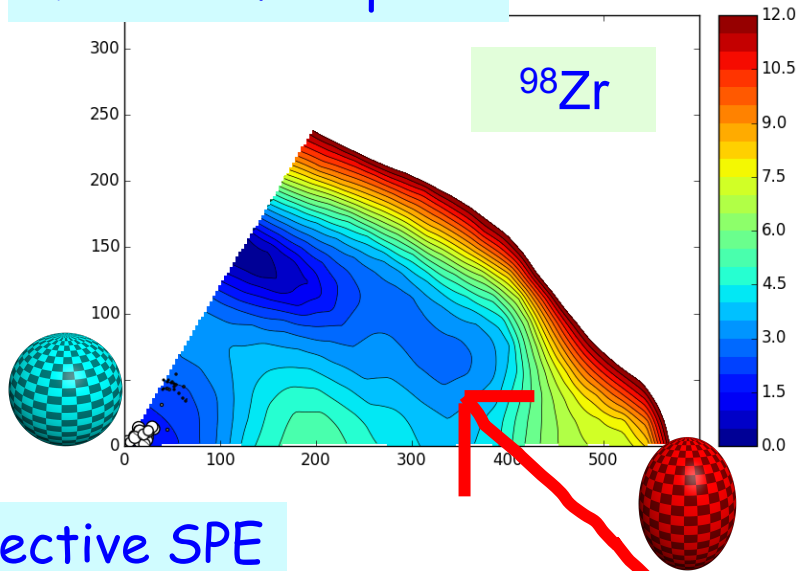
Mixing depends not only on the strength of the $Y_{2,0}(\theta,\phi)$ field, but also the spherical single-particle energies $\epsilon_1, \epsilon_2, \epsilon_3$, etc.



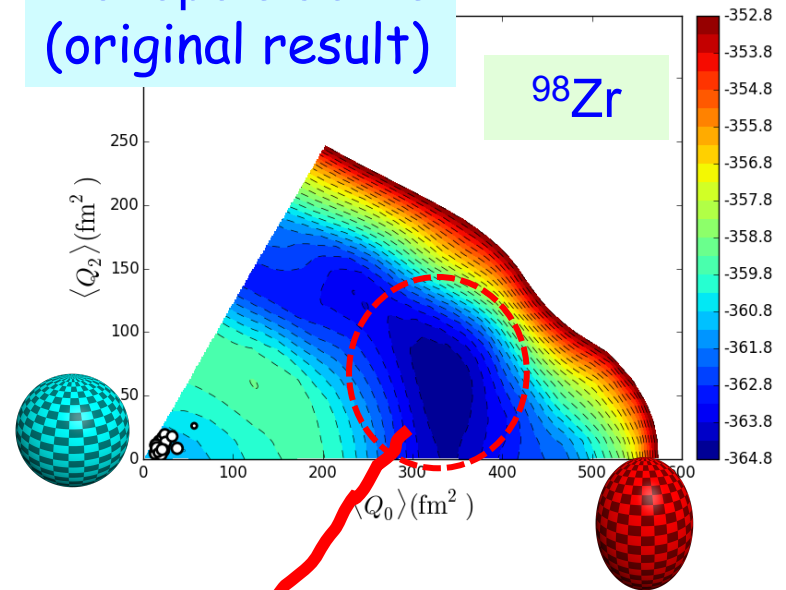
Note: single-particle states are enemies against collective modes if their splitting is too large

Anatomy of this effect : Freezing this monopole effect

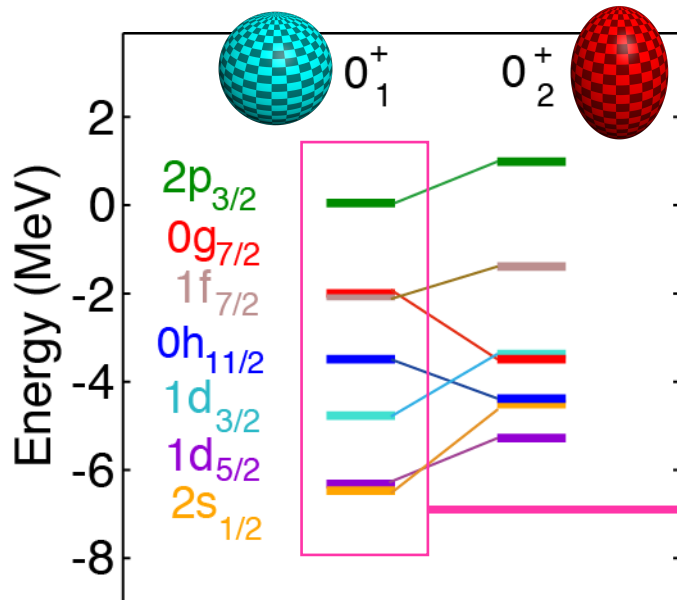
Frozen monopole



monopole active (original result)



Effective SPE



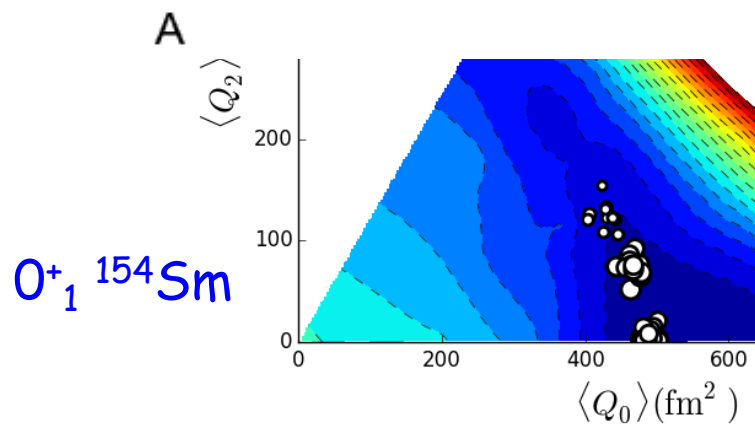
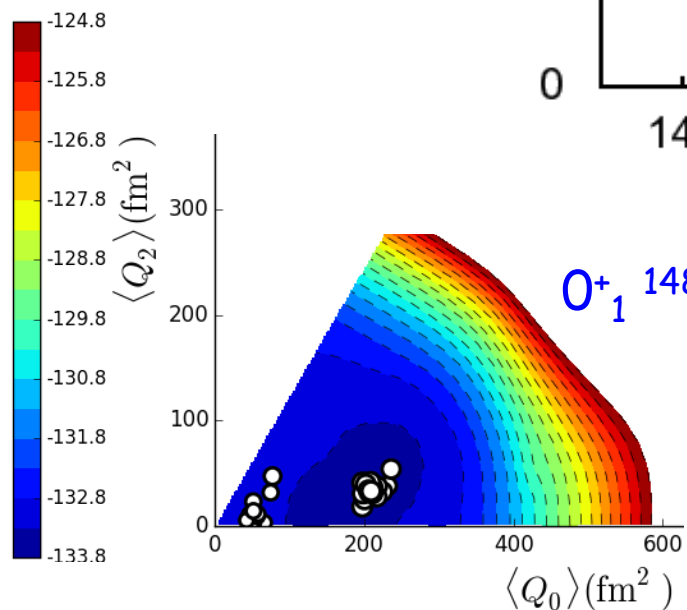
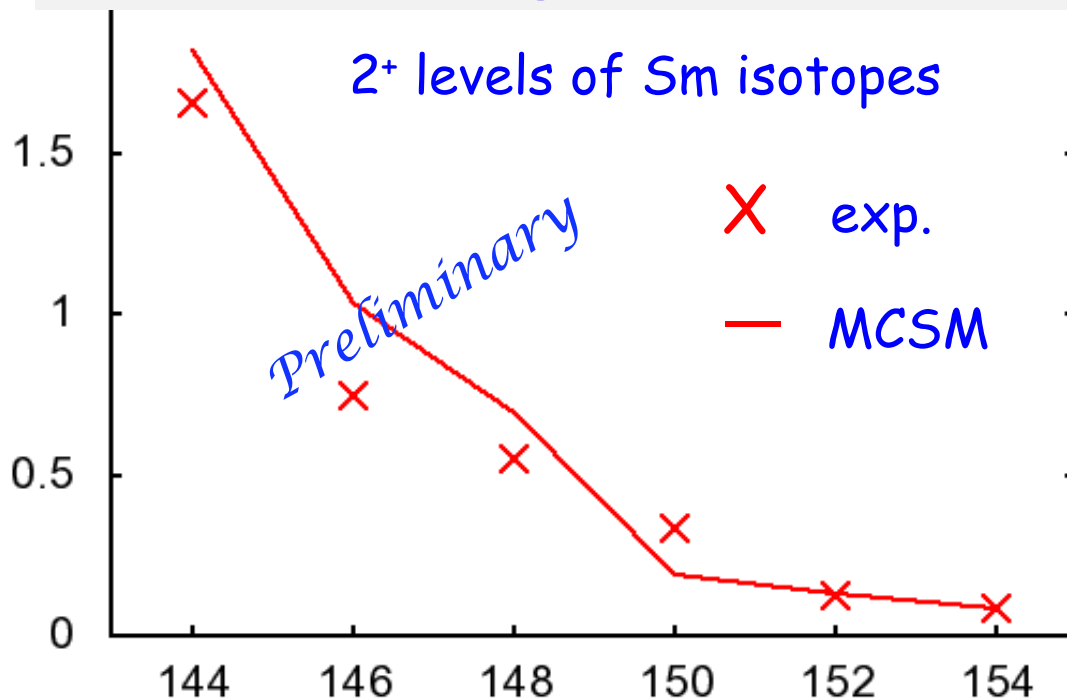
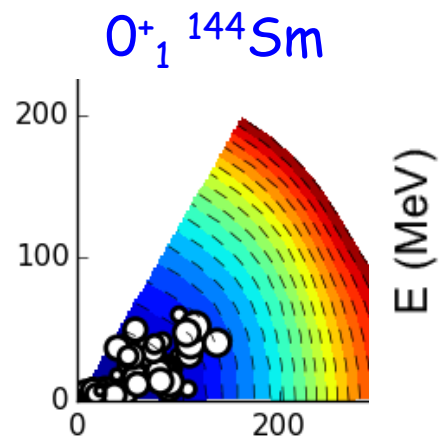
energy minimum for ellipsoid is gone

Use them as constant single-particle energies *independent of configurations, putting monopole int. aside*
 → Frozen monopole treatment

Sm isotopes

proton 8 orbits
neutron 10 orbits

Interaction VMU (gaussian central + $\pi+\rho$ tensor)



$$\text{deformation} = \frac{\text{quadrupole force}}{\text{resistance power}}$$

resistance power ← pairing force

↑
single-particle energies

Analogy to electric current,

$$\text{current} = \frac{\text{voltage}}{\text{resistance}}$$

Remarks

Naïve Fermi liquid picture (a la Landau) is revised, as atomic nuclei are not necessarily like simple solid vases containing almost free nucleons.

Nuclear forces are rich enough to optimize single-particle energies for each eigenstate (especially in the cases of collective-mode states), as referred to as **quantum self-organization**.

The **quantum self-organization** produces sizable effects with

- (i) **two quantum fluids** (protons and neutrons),
- (ii) **two major forces** : *e.g.*, quadrupole interaction to **drive collective mode**
monopole interaction to **control resistance**

Quantum phase transition, shape coexistence, various deformation, fission, ... are related to the **quantum self-organization**.

Thus, non-specific forces work coherently so that the beauty of the collective modes is achieved or enhanced : single-particle states are not always enemies but friends of collective modes.

Time-dependent version for reactions is of great interest (e.g., fission). A more intriguing topic is a possible relation to biological evolution.