

The modified Picard iteration and its applications

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The Picard iteration

- A process to solve differential equations

Consider the following initial value problem:

$$\frac{dy}{dx} = f(x, y) \quad y(x_0) = y_0$$

The general solution:

$$y(x) = y_0 + \int_{x_0}^x f(X, y) dX$$

The Picard iterative process:

$$y_{n+1}(x) = y_0 + \int_{x_0}^x f(X, y_n) dX \quad y(x_0) = y_0$$

If “f” is analytical, the sequence of Taylor polynomials also converge to “y”

The modified Picard iteration

(The Parker-Sochacki method)

- The Picard iteration is not so practical - “...It is impossible to compute explicitly more than a few members of the sequence, so the limit function can only be found in rare cases...”
- G. E Parker & J. S. Sochacki - Implementing the Picard iteration
 - ⇒ generate or approximate the Taylor series solutions to any ODE, that has a polynomial generator

The Parker-Sochacki method

- Equations of motions:

$$\frac{dx}{dt} = v$$

$$\frac{dv}{dt} = f(x, t)$$

- Assume that x and v can be expressed as a truncated Taylor series in time:

$$x := x_0 + x_1 t + x_2 t^2 + \dots + x_n t^n = \sum_{j=0}^n x_j t^j$$

$$v := v_0 + v_1 t + v_2 t^2 + \dots + v_n t^n = \sum_{j=0}^n v_j t^j$$

$$x_1 + 2x_2 t + \dots + n x_n t^{n-1} = v_0 + v_1 t + v_2 t^2 + \dots + v_n t^n$$

$$\rightarrow x_{n+1} = v_n / (n+1)$$

$$v_{n+1} = k_n(x) / (n+1)$$

The N-body problem

- Equations of motions:

$$\frac{dx_{il}}{dt} = v_{il}$$

$$\frac{dv_{il}}{dt} = \sum_{k=1}^N G m_k \frac{(x_{ik} - x_{il})}{\rho_{lk}^3}$$

$$\rho_{lk} = \left(\sum_{i=1}^3 [x_{il} - x_{ik}]^2 \right)^{1/2}$$

- The separation in the denominator makes it analytically unsolvable
- The key of the modified Picard iteration: to replace these factors with a polynomial approximation: $\xi_{lk} := 1/\rho_{lk}$

- The modified equations:

$$\frac{dx_{il}}{dt} = v_{il}$$

$$\frac{dv_{il}}{dt} = \sum_{k=1}^N G m_k (x_{ik} - x_{il}) \xi_{lk}^3$$

$$\frac{d\xi_{lk}}{dt} = -\xi_{lk}^3 \left[\sum_{i=1}^3 (x_{il} - x_{ik})(v_{il} - v_{ik}) \right]$$

The N-body problem

$$\begin{aligned}\frac{dx_{il}}{dt} &= v_{il} \\ \frac{dv_{il}}{dt} &= \sum_{k=1}^N G m_k (x_{ik} - x_{il}) \xi_{lk}^3 \\ \frac{d\xi_{lk}}{dt} &= -\xi_{lk}^3 \left[\sum_{i=1}^3 (x_{il} - x_{ik})(v_{il} - v_{ik}) \right]\end{aligned}$$

- The method:
- Assuming that

$$x := \sum_{j=0}^p x_p dt^p; \quad v := \sum_{j=0}^p v_p dt^p; \quad \xi := \sum_{j=0}^p \xi_p dt^p$$

⇒

$$x_{i,l,p+1} = v_{ilp} / (p + 1)$$

$$v_{i,l,p+1} = \left[\sum_{k=1}^N G m_k \sum_{j=0}^p (x_{ikj} - x_{ilj}) (\xi_{lk}^3)_{p-j} \right] / (p + 1)$$

$$\xi_{l,k,p+1} = \left[\sum_{j=0}^p (\xi_{lk}^3)_j \alpha_{l,k,p-j} \right] / (p + 1)$$

$$\text{ahol} \quad \alpha_{lkp} = \sum_{j=0}^p \sum_{i=1}^3 (x_{ilj} - x_{ikj}) (v_{i,l,p-j} - v_{i,k,p-j})$$

$$(\xi_{lk}^3)_p = \sum_{j=0}^p (\xi_{lk}^2)_j \xi_{l,k,p-j} \quad (\xi_{lk}^2)_p = \sum_{j=0}^p \xi_{lkj} \xi_{l,k,p-j}$$

The motions of satellites - The Gravitational potential

- Equations of motions of satellites:

$$\underline{\ddot{\mathbf{r}}} = \frac{d\underline{\mathbf{U}}}{d\underline{\mathbf{r}}}$$

$$U = GM \sum_{l=0}^{\infty} \sum_{m=0}^l \left[\frac{\bar{r}_{\oplus}^l}{r^{l+1}} \cdot P_l^m(\sin\varphi) \cdot (C_{lm}\cos(m\lambda) + S_{lm}\sin(m\lambda)) \right]$$

$$C_{lm} = \frac{1}{M\bar{r}_{\oplus}^l} \frac{2(l-m)!}{\delta_m(l+m)!} \int_V r'^l P_l^m(\sin\varphi') \cos(m\lambda') \rho d\tau,$$

$$S_{lm} = \frac{1}{M\bar{r}_{\oplus}^l} \frac{2(l-m)!}{\delta_m(l+m)!} \int_V r'^l P_l^m(\sin\varphi') \sin(m\lambda') \rho d\tau,$$

- Changing the coordinate system:

$$U = \sum_{l=0}^{\infty} \sum_{m=0}^l \sum_{k=0}^m \left[\frac{\beta_{lmk}}{(\mathbf{x}^2 + \mathbf{y}^2 + \mathbf{z}^2)^{\frac{l+1}{2}}} \cdot P_l^m \left(\frac{z}{(\mathbf{x}^2 + \mathbf{y}^2 + \mathbf{z}^2)^{1/2}} \right) \cdot \frac{\mathbf{x}^k \mathbf{y}^{m-k}}{(\mathbf{x}^2 + \mathbf{y}^2)^{m/2}} \right]$$

$$\beta_{lmk} = GM\bar{r}_{\oplus}^l \binom{m}{k} \left[C_{lm}\cos\left(\frac{1}{2}(m-k) \cdot \pi\right) + S_{lm}\sin\left(\frac{1}{2}(m-k) \cdot \pi\right) \right]$$

The Gravitational potential

- Equations of motions:

$$\frac{dx}{dt} = v_x \quad \frac{dy}{dt} = v_y \quad \frac{dz}{dt} = v_z$$

$$\begin{aligned} \frac{dv_x}{dt} = & \sum_{l=0}^{\infty} \sum_{m=0}^l \sum_{k=0}^m \frac{\beta_{lmk} \cdot x^{k-1} y^{m-k}}{(x^2 + y^2)^{\frac{m+2}{2}} (x^2 + y^2 + z^2)^{\frac{l+2}{2}}} \cdot \\ & \cdot \left[P_l^m \left(\frac{z}{(x^2 + y^2 + z^2)^{1/2}} \right) (x^2 + y^2 + z^2)^{1/2} \cdot [ky^2 + x^2(k - m - l - 1)] + P_{l+1}^m \left(\frac{z}{(x^2 + y^2 + z^2)^{1/2}} \right) x^2 z(l - m + 1) \right] \end{aligned}$$

$$\begin{aligned} \frac{dv_y}{dt} = & \sum_{l=0}^{\infty} \sum_{m=0}^l \sum_{k=0}^m \frac{\beta_{lmk} \cdot x^k y^{m-k-1}}{(x^2 + y^2)^{\frac{m+2}{2}} (x^2 + y^2 + z^2)^{\frac{l+2}{2}}} \cdot \\ & \cdot \left[-P_l^m \left(\frac{z}{(x^2 + y^2 + z^2)^{1/2}} \right) (x^2 + y^2 + z^2)^{1/2} \cdot [y^2(k + l + 1) + x^2(k - m)] + P_{l+1}^m \left(\frac{z}{(x^2 + y^2 + z^2)^{1/2}} \right) y^2 z(l - m + 1) \right] \end{aligned}$$

$$\frac{dv_z}{dt} = \sum_{l=0}^{\infty} \sum_{m=0}^l \sum_{k=0}^m \frac{\beta_{lmk} \cdot x^k y^{m-k}}{(x^2 + y^2)^{\frac{m+2}{2}} (x^2 + y^2 + z^2)^{\frac{l+2}{2}}} \cdot \left[P_{l+1}^m \left(\frac{z}{(x^2 + y^2 + z^2)^{1/2}} \right) (x^2 + y^2)(l - m + 1) \right]$$

- Using additional variables:

$$r := (x^2 + y^2 + z^2)^{1/2} \quad s := 1/r; \quad u := (x^2 + y^2)^{1/2} \quad w := 1/u;$$

The Gravitational potential

- Equations of motions:

$$\begin{aligned} \frac{dx}{dt} &= v_x \\ \frac{dy}{dt} &= v_y \\ \frac{dz}{dt} &= v_z \\ \frac{dv_x}{dt} &= \sum_{l=0}^{\infty} \sum_{m=0}^l \sum_{k=0}^m \beta_{lmk} x^{k-1} y^{(m-k)} w^{(m+2)} s^{(l+2)} [P_l^m(s \cdot z) r \{ky^2 + (k - m - l - 1)x^2\} + P_{l+1}^m(s \cdot z) x^2 z(l - m + 1)] \\ \frac{dv_y}{dt} &= \sum_{l=0}^{\infty} \sum_{m=0}^l \sum_{k=0}^m \beta_{lmk} x^k y^{(m-k-1)} w^{(m+2)} s^{(l+2)} [-P_l^m(s \cdot z) r \{(k + l + 1)y^2 + (k - m)x^2\} + P_{l+1}^m(s \cdot z) y^2 z(l - m + 1)] \\ \frac{dv_z}{dt} &= \sum_{l=0}^{\infty} \sum_{m=0}^l \sum_{k=0}^m \beta_{lmk} x^k y^{(m-k)} w^{(m+2)} s^{(l+2)} [P_{l+1}^m(s \cdot z) (x^2 + y^2)(l - m + 1)] \\ \frac{dr}{dt} &= s(xv_x + yv_y + zv_z) \\ \frac{ds}{dt} &= -s^3(xv_x + yv_y + zv_z) \\ \frac{dw}{dt} &= -w^3(xv_x + yv_y) \end{aligned}$$

The Gravitational potential

- (1) $\mathbf{x}_{n+1} = \mathbf{v}_{x,n}/(\mathbf{n} + 1) \rightarrow (x^2)_n = \sum_{j=0}^n (x^2)_j x_{n-j}, \dots, (x^{(k-1)})_n = \sum_{j=0}^n (x^{(k-2)})_j x_{n-j}, (x^k)_n = \sum_{j=0}^n (x^{(k-1)})_j x_{n-j}$
- (2) $\mathbf{y}_{n+1} = \mathbf{v}_{y,n}/(\mathbf{n} + 1) \rightarrow (y^2)_n = \sum_{j=0}^n (y^2)_j y_{n-j}, \dots, (y^{(m-k-1)})_n = \sum_{j=0}^n (y^{(m-k-2)})_j y_{n-j}, (y^{(m-k)})_n = \sum_{j=0}^n (y^{(m-k-1)})_j y_{n-j}$
- (3) $\mathbf{z}_{n+1} = \mathbf{v}_{z,n}/(\mathbf{n} + 1)$
- (7) $\mathbf{r}_{n+1} = \left[\sum_{j=0}^n \mathbf{s}_j (\mathbf{r}^*)_{n-j} \right] / (\mathbf{n} + 1) \Leftrightarrow (r^*)_n = \sum_{j=0}^n (x_j v_{x,n-j} + y_j v_{y,n-j} + z_j v_{z,n-j})$
- (8) $\mathbf{s}_{n+1} = - \left[\sum_{j=0}^n (\mathbf{s}^3)_j (\mathbf{r}^*)_{n-j} \right] / (\mathbf{n} + 1) \Leftrightarrow (s^2)_n = \sum_{j=0}^n s_j s_{n-j}, (s^3)_n = \sum_{j=0}^n (s^2)_j s_{n-j}, \dots, (s^{(l+2)})_n = \sum_{j=0}^n (s^{(l+1)})_j s_{n-j}$
- (9) $\mathbf{w}_{n+1} = - \left[\sum_{j=0}^n (\mathbf{w}^3)_j (\mathbf{w}^*)_{n-j} \right] / (\mathbf{n} + 1) \Leftrightarrow (w^*)_n = \sum_{j=0}^n (x_j v_{x,n-j} + y_j v_{y,n-j})$
 $(w^2)_n = \sum_{j=0}^n w_j w_{n-j}, (w^3)_n = \sum_{j=0}^n (w^2)_j w_{n-j}, \dots, (w^{(m+2)})_n = \sum_{j=0}^n (w^{(m+1)})_j w_{n-j}$
- (4) $\mathbf{v}_{x,n+1} \cdot (\mathbf{n} + 1) = \sum_{l=0}^{\infty} \sum_{m=0}^l \sum_{k=0}^m \beta_{lmk} \sum_{j=0}^n (\mathbf{a}_{lmk})_j [(k - m - l - 1)(\mathbf{e}_{lm})_{n-j} + k(\mathbf{f}_{lm})_{n-j} + (l - m + 1)(\mathbf{g}_{lm})_{n-j}]$
- (5) $\mathbf{v}_{y,n+1} \cdot (\mathbf{n} + 1) = \sum_{l=0}^{\infty} \sum_{m=0}^l \sum_{k=0}^m \beta_{lmk} \sum_{j=0}^n (\mathbf{b}_{lmk})_j [(m - k)(\mathbf{e}_{lm})_{n-j} - (k + l + 1)(\mathbf{f}_{lm})_{n-j} + (l - m + 1)(\mathbf{h}_{lm})_{n-j}]$
- (6) $\mathbf{v}_{z,n+1} \cdot (\mathbf{n} + 1) = \sum_{l=0}^{\infty} \sum_{m=0}^l \sum_{k=0}^m \beta_{lmk} (l - m + 1) \sum_{j=0}^n (\mathbf{c}_{lmk})_j [(\mathbf{G}_{lm})_{n-j} + (\mathbf{H}_{lm})_{n-j}] (l - m + 1)$

↑

$\mathbf{a}_{lmk} = A_{km} D_{lm} \quad \mathbf{b}_{lmk} = B_{km} D_{lm} \quad \mathbf{c}_{lmk} = C_{km} D_{lm} \quad \mathbf{f}_{lm} = r F_{lm} \quad \mathbf{e}_{lm} = r E_{lm} \quad \mathbf{g}_{lm} = z G_{lm} \quad \mathbf{h}_{lm} = z H_{lm}$

$A_{km} = x^{k-1} y^{m-k} \quad B_{km} = x^k y^{m-k-1} \quad C_{km} = x^k y^{m-k} \quad D_{lm} = w^{m+2} s^{l+2}$

$E_{lm} = P_l^m(s \cdot z) x^2 \quad F_{lm} = P_l^m(s \cdot z) y^2 \quad G_{lm} = P_{l+1}^m(s \cdot z) x^2 \quad H_{lm} = P_{l+1}^m(s \cdot z) y^2$

The motions of satellites:

Additional forces:

Additional terms in the equations of motions

- Inertial Forces:

Centrifugal force, Coriolis force, Euler force

- N-body problem (+ Moon, +Sun)

- Air Friction force