Nobel prize for Eugene P. Wigner





Nobelprize in Physics 1963

"for his contributions to the theory of the atomic nucleus and the elementary particles, particularly through the discovery and application of fundamental symmetry principles"

Eugene P. Wigner *17.11.1902, Budapest 1.1.1995, Princeton

together with Maria Göppert-Maier and J. Hans D. Jensen

Symmetries and in-medium effects

breaking and restoration of chiral symmetry and in-medium modifications of hadrons

Volker Metag* II. Physikalisches Institut

> JUSTUS-LIEBIG-UNIVERSITAT GIESSEN

*and University of Bonn

<u>Outline</u>

- chiral symmetry breaking and restoration
- experimental approaches for studying in-medium properties of hadrons
- in-medium properties of the η ' meson
- summary and outlook

*funded by the DFG within SFB/TR16



Wigner 111 symposium Nov. 11-14. 2013, Budapest, Hungary



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 - invariant under $SU(3)_L \otimes SU(3)_R$ chiral transformation $\psi_{L,I}$

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• $q_L \Rightarrow q_R; \chi$ symmetry broken in presence of a chiral condensate

<0 $|q\bar{q}|$ 0> is order parameter for χ -symmetry

 χ -symmetry restored for <0|q \overline{q} |0> → 0

symmetry breaking in the hadronic sector

baryons: mesons: • for χ symmetry chiral partners $a_1:1^+1260$ $\sigma:0^{+}$ 600 $\frac{3}{2}$ 1520 1535 (hadronic states with same spin 290 490 but opposite parity) should be 470 600 1232 degenerate in mass: $m(J^{\pi}) = m(J^{-\pi})$ 22 22 experiment: $\Delta m \neq 0$ $\pi: 0^{-1}$ 135 $\rho:1^{-1}$ 770 $\frac{1/^{+}}{2}$ 938 vector scalar χ symmetry chiral nucleon mesons mesons symmetry

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H. Nagahiro et al., PRC 87 (2013) 045201 The NJL Model



mass as a result of symmetry breaking



Link between hadronic and QCD- descriptions

chiral condensate as function of baryon density ρ_B and temperature T



partial restoration of chiral symmetry $|\langle \bar{q}q \rangle_{\rho,T} | \rightarrow 0$ for ρ_B , T. S. Klimt et al., PLB 249 (1990) 386

parameter range (ρ_B ,T) reached in reactions with heavy-ion, γ , π , p-beams

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 $\frac{\text{QCD sum rules:}}{\frac{\text{hadronic side}}{\pi}} \frac{\text{QCD sum rules:}}{\frac{\text{QCD side}}{\frac{Q^2}{\pi}\int_0^\infty ds \frac{Im\Pi(s)}{s(s+Q^2)}} = -\frac{1}{8\pi^2} (1+\frac{\alpha_s}{\pi}) ln \frac{Q^2}{\Lambda^2} + \frac{m_q \langle \overline{q}q \rangle}{Q^4} + \frac{1}{24} \frac{\langle \frac{\alpha_s}{\pi} G^2 \rangle}{Q^4} - \frac{112}{81} \alpha_s \pi \frac{\langle \overline{q}q \rangle^2}{Q^6} + \dots$

no direct relation between in-medium properties of hadrons and QCD condensates QCD sum rules provide important constraints for hadronic models hadronic models needed for quantitative comparison to experiments 5

spectral function of the η ' meson

model predictions - based on NJL

V. Bernard und U.-G. Meißner, Phys. Rev. D 38 (1988) 1551



the mass of the η ' meson is almost independent of density

Y. Kwon, S.H. Lee, K. Morita, and G. Wolf, PRD 86 (2012) 034014

H. Nagahiro, M. Takizawa and S. Hirenzaki, Phys. Rev. C 74 (2006) 045203



even at normal nuclear density

 $U_A(I)$ breaking part of η' mass (\approx 460 MeV) lowered by 20%: $\Box \Rightarrow \Delta m_{\eta'}(\rho_0) \approx$ -90 MeV

discrepancy between theoretical predictions requires exprimental clarification !!

• $p p \rightarrow p p \eta'$ @ COSY (P. Mos

(P. Moskal et al., PLB 482 (2000) 356)

analysis of final state interactions gives estimate of η' - proton scattering length $|a_{\eta'P}| \approx 0.1$ fm weak η' - nucleon interaction

ultra-relativistic heavy-ion collisions @RHIC (Csörgo et al., PRL 105 (2010) 182301) two-pion Bose-Einstein correlation analysis of PHENIX and STAR data indicates massdrop of η ' meson in highly compressed and heated collision zone $\Delta m_{\eta'} \approx -200 \text{ MeV}$ wery strong η '- nucleon interaction

> conflicting experimental evidence !!! further experiments needed

photoproduction of η ' mesons off nuclei at ELSA

CBELSA/TAPS-detector system : $\approx 4\pi$ photon detector system

 $E_{Y} = 1.0-2.8 \text{ GeV}$



m = 958 MeV/c² σ = 11 MeV Δ m/m = 1.1% Experimental approaches to determine the meson-nucleus optical potential

meson-nucleus optical potential U(r) = V(r) + iW(r) Experimental approaches to determine the meson-nucleus optical potential

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$$\frac{\text{meson mass shift}}{\Gamma(r) = \Delta m(\rho_0) \cdot \frac{\rho(r)}{\rho_0}}$$

<u>line shape analysis:</u> direct determination of Δm

excitation function: provides information about the depth of V(r)

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meson-nucleus optical potential

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<u>meson-nucleus-bound states:</u> direct determination of E_{bin} (Δm) $\frac{\text{meson absorption}}{W(r) = -\Gamma_0/2 \cdot \frac{\rho(r)}{\rho_0}}$ $= -\frac{1}{2} \cdot \hbar c \cdot \rho(r) \cdot \sigma_{inel} \cdot \beta$

Transparency ratio measurement

$$\Gamma_A = \frac{\sigma_{\gamma A \to \eta' X}}{A \cdot \sigma_{\gamma N \to \eta' X}}$$

attenuation measurement of meson flux:

(D. Cabrera et al., NPA 733 (2004) 130)

$$T_{A} = \frac{\sigma_{\gamma A \to \eta' X}}{A \cdot \sigma_{\gamma N \to \eta' X}}$$

production probability per nucleon within the nucleus compared to production probability on the free nucleon;



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inelastic reactions remove ω , η' mesons, e.g. ω , $\eta' N \rightarrow \pi N$ shortening of ω , η' lifetime in the medium \Rightarrow increase in width

low density approximation:
$$\Gamma(\rho) = -\frac{Im\Pi(\rho)}{E} = \hbar c \cdot \rho \cdot \beta \cdot \sigma_{inel}; \ \Gamma(\rho) = \Gamma(\rho_0) \frac{\rho}{\rho_0}$$

information on imaginary part of meson-nucleus potential

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information on in-medium properties of mesons from measurement of their decay outside of the nucleus

The imaginary part of the η '-nucleus potential

 $E_{\gamma} = 1500 - 2200 \text{ MeV}$; photoproduction of η ' meson off ${}^{12}C$, ${}^{40}Ca$, ${}^{93}Nb$ and ${}^{208}Pb$

 $T_A^C = \frac{12 \cdot \sigma_{\gamma A \to \eta' X}}{A \cdot \sigma_{\gamma C \to \eta' X}} \text{ normalized to carbon}$

M. Nanova et al., PLB 710 (2012) 600



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comparison with other mesons



 η' interaction with nuclear matter much weaker than for η, ω mesons

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 $W_{\eta'}(\rho = \rho_0) = -\Gamma_0/2 = -(7.5-12.5) \text{ MeV}$

Determination of the real part of the η '-nucleus potential

J.Weil, U. Mosel and V. Metag, PLB 723 (2013) 120

E. Paryev, J. Phys. G: Nucl. Part. Phys. 40 (2013) 025201

measurement of the excitation function:

in case of dropping mass - higher meson yield for given \sqrt{s} because of increased phase space due to lowering of the production threshold



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 measurement of the momentum distribution: in case of dropping mass - when leaving the nucleus hadron has to become on-shell; mass generated at the expense of kinetic energy
 downward shift of momentum distribution



differential cross sections for η' photo production off carbon

M. Nanova et al., arXiv:1311.0122; accepted for publication in PLB



forward rise of cross section with increasing incident photon energy: t-channel production mechanism

excitation function for η' photoproduction off C

comparison of CBELSA/TAPS data with calculations by E. Paryev, J. Phys. G: Nucl. Part. Phys. 40 (2013) 025201 and priv. communication

decay mode: $\eta' \rightarrow \pi^0 \pi^0 \eta$ excitation function



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estimation of the real part of the η '-nucleus potential from the η ' excitation function

excitation function


0^μ, [mp] C data ο σ_{tot} σ_{diff} 1 $V(\rho = \rho_0) = 0 \text{ MeV}$ **V**(ρ=ρ₀) = -25 MeV **V**(ρ=ρ₀) = -50 MeV V(ρ=ρ₀) = -75 MeV = -100 MeV /(ρ=ρ₀) = -150 MeV $\sigma_{n'N}=11mb$ -1 10 $\mathbf{E}_{\mathbf{y}}$ thr 1000 2000 1500 2500 E_v[MeV]

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excitation function





 η ' momentum distribution off C

comparison of CBELSA/TAPS data with calculations by E. Paryev, J. Phys. G: Nucl. Part. Phys. 40 (2013) 025201 and priv. communication

momentum distribution



 η ' momentum distribution off C

comparison of CBELSA/TAPS data with calculations by E. Paryev, J. Phys. G: Nucl. Part. Phys. 40 (2013) 025201 and priv. communication



































consistent with predictions by:

S. Bass and A.W.Thomas, Acta Phys. Pol. B 41 (2010) 2239

Summary

I. Imaginary part of the η '- nucleus optical potential determined from transparency ratio measurements: $W(\rho = \rho_0) = -\Gamma_0/2 = -10\pm 2.5 \text{ MeV}$

2. Real part of the η '- nucleus optical potential determined from:

a. measurement of the excitation function of the η '-meson

$$V(\rho = \rho_0) = -40 \pm 6 \text{ MeV}$$

b. measurement of the momentum distribution of the η '-meson:

$$V(\rho = \rho_0) = -32 \pm 11$$
 MeV

$$U_{\eta'A}(\rho = \rho_0) = -(37 \pm 10(\text{stat}) \pm 10(\text{syst}) + i(10 \pm 2.5)) \text{ MeV}$$

M. Nanova et al., arXiv:1311.0122; to be published in PLB

First (indirect) observation of a mass drop of a peusdoscalar meson in nuclear matter at normal conditions ($\rho = \rho_0$; T=0)



3. $V >> W! \Rightarrow \eta'$ promising candidate for mesic states experiments searching for η' mesic states at FRS@GSI and BGO-OD@ELSA

search for η' mesic states at FRS@GSI and BGO-OD@ELSA

GSI ¹²C(p,d) η'X @ 2.5 GeV

K. Itahashi et al.:FRS@GSI

K. Itahashi et al., Prog. Theo. Phys. 128 (2012) 601

missing mass spectroscopy





γ+¹²C→η'⊗¹¹B+p @2.2-2.8 GeV **ELSA**

<u>semi-exclusive measurement:</u>

coincident detection of forward going proton and decays of η ' mesic states: η 'N $\rightarrow \eta$ N




differential cross sections for η ' photoproduction off ^{12}C

 $E_{\gamma} = 1250 - 2600 \text{ MeV}$ $\eta' \rightarrow \pi^0 \pi^0 \eta \rightarrow 6\gamma$ BR: 8.1%

sensitivity to different scenarios

E. Ya. Paryev, priv. communication



high sensitivity to different scenarios at threshold

strong mass shift not supported by data