

Chiral asymmetry in nanomagnetism

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Outline of the talk

Spin Hamiltonians and Dzyaloshinskii-Moriya interaction (DMI)

Effects in magnetic nanostructures induced by DMI:

- *Magnetic pattern formation in ultrathin films*

Spin-spirals

Magnetic skyrmions & skyrmion lattices

- *Asymmetry of the spin-wave spectra*
- *Perpendicular exchange bias*

Introduction

Effects of spin-orbit coupling in nanomagnetism:

Magneto-crystalline anisotropy

Spin-reorientation phase transitions

Superparamagnetism

Magneto-optics: Faraday effect, Kerr effect

Spintronics:

Bychkov-Rashba effect

Transverse transport phenomena (AHE, SHE)

...

Classical spin Hamiltonian

$$H = \frac{1}{2} \sum_{i \neq j} \vec{e}_i \mathbf{J}_{ij} \vec{e}_j + \frac{1}{2} \sum_{i \neq j} \vec{e}_i \mathbf{M}_{ij} \vec{e}_j + \sum_i \vec{e}_i \mathbf{K}_i \vec{e}_i$$

exchange
interaction

magnetic dipole-dipole
interaction

on-site anisotropy

Tensorial exchange interaction

$$\mathbf{J}_{ij} = J_{ij} \mathbf{I} + \mathbf{J}_{ij}^S + \mathbf{J}_{ij}^A$$

$$J_{ij} = \frac{1}{3} \text{Tr} \mathbf{J}_{ij}$$

isotropic

$$\mathbf{J}_{ij}^S = \frac{1}{2} (\mathbf{J}_{ij} + \mathbf{J}_{ij}^t) - J_{ij} \mathbf{I}$$

anisotropic symmetric

$$\mathbf{J}_{ij}^A = \frac{1}{2} (\mathbf{J}_{ij} - \mathbf{J}_{ij}^t)$$

antisymmetric

relativistic (spin-orbit) effects

Dzyaloshinskii-Moriya interaction (DMI)

I. E. Dzyaloshinskii, Sov. Phys. JETP 5, 1259–1262 (1957)

T. Moriya, Phys. Rev. 120, 91–98 (1960)

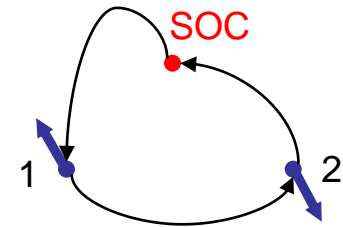
$$E_{DM} = \vec{e}_i \mathbf{J}_{ij}^A \vec{e}_j = \vec{D}_{ij} (\vec{e}_i \times \vec{e}_j)$$



$$D_{ij}^\alpha = \frac{1}{2} \epsilon_{\alpha\beta\gamma} J_{ij}^{\beta\gamma}$$

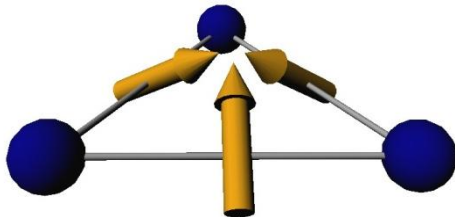
Asymmetric exchange (RKKY) interaction:
DMI scales with the strength of SOC

A. Fert and P. M. Levy, Phys. Rev. Lett. 44, 1538 (1980)



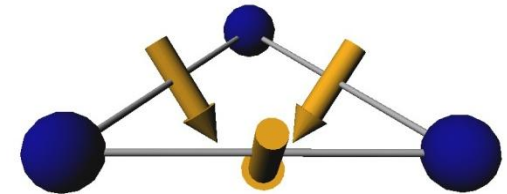
Equilateral Cr trimer on top of Au(111)

A. Antal *et al.*, Phys. Rev. B **77**, 174429 (2008)

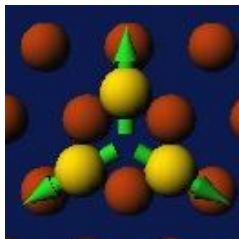


$$D_z > 0$$

DM vectors



$$D_z < 0$$

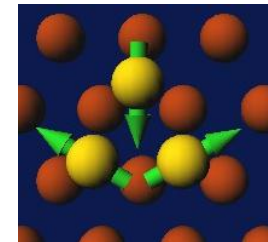


$$\kappa_z = -1$$

Chirality

$$\vec{\kappa} = \frac{2}{3\sqrt{3}} \sum_{(ij)} (\vec{\sigma}_i \times \vec{\sigma}_j)$$

$$E_{DM} = \frac{3\sqrt{3}}{2} D_z \kappa_z$$



$$\kappa_z = 1$$

Ab initio methods to calculate tensorial exchange interactions & anisotropy parameters

- Relativistic torque method
 - A.I. Liechtenstein *et al.*, JMMM 67, 65 (1987)
 - L. Udvardi *et al.*, PRB 68 104436 (2003)
 - H. Ebert and S. Mankovsky, PRB 79, 045209 (2009)

Based on zero temperature ordered (ferromagnetic) state

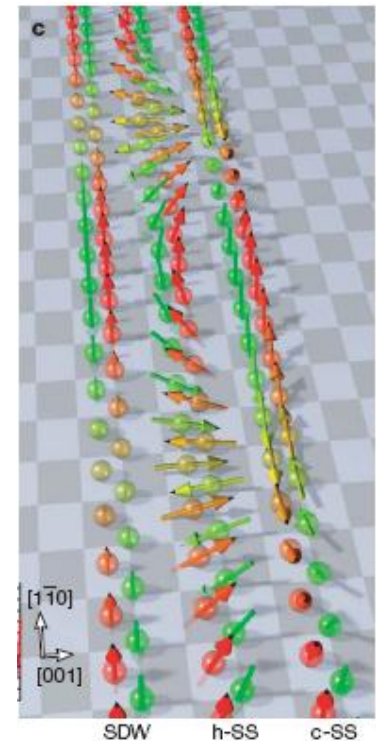
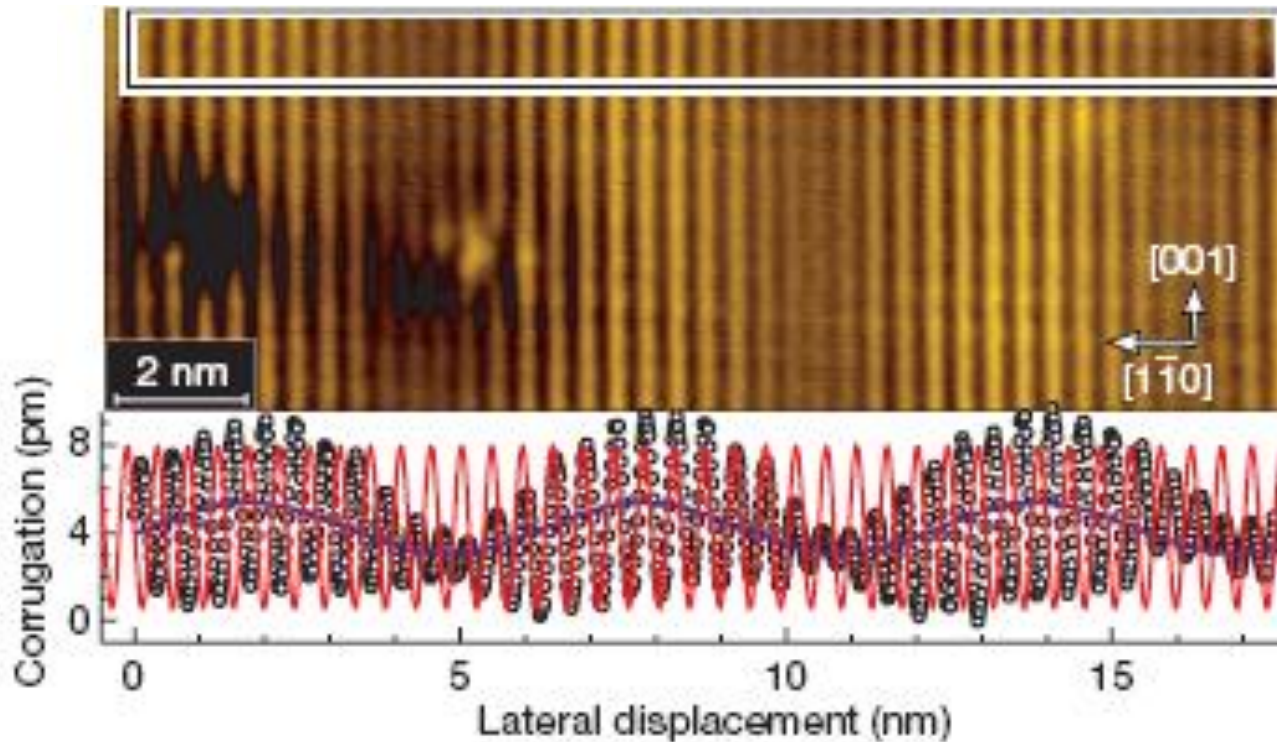
- Spin-cluster expansion
 - R. Drautz and M. Fähnle, PRB 69, 104404 (2004),
PRB 72, 212405 (2005)
 - L. Szunyogh *et al.*, PRB 83, 024401 (2011)
 - A. Deák *et al.*, PRB 84, 224413 (2011)

Based on high temperature paramagnetic (Disordered Local Moment) state

Effects of DM interactions in thin magnetic films: Formation of spin-spirals

Mn monolayer on W(110) M. Bode et al., Nature 447, 193 (2007)

Constant current SP-STM image



- row-by-row AF structure with a long-wavelength (12 nm) modulation
- cycloidal spin-spiral (CSS) → spins rotate around the (001) axis
- **ab initio calculations confirmed DMI as origin of CSS**

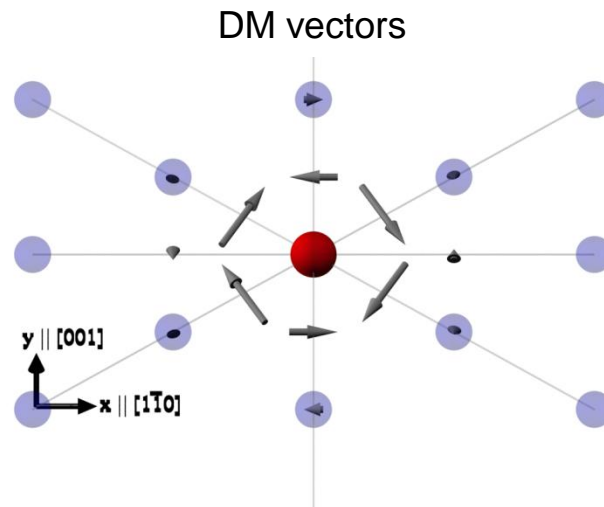
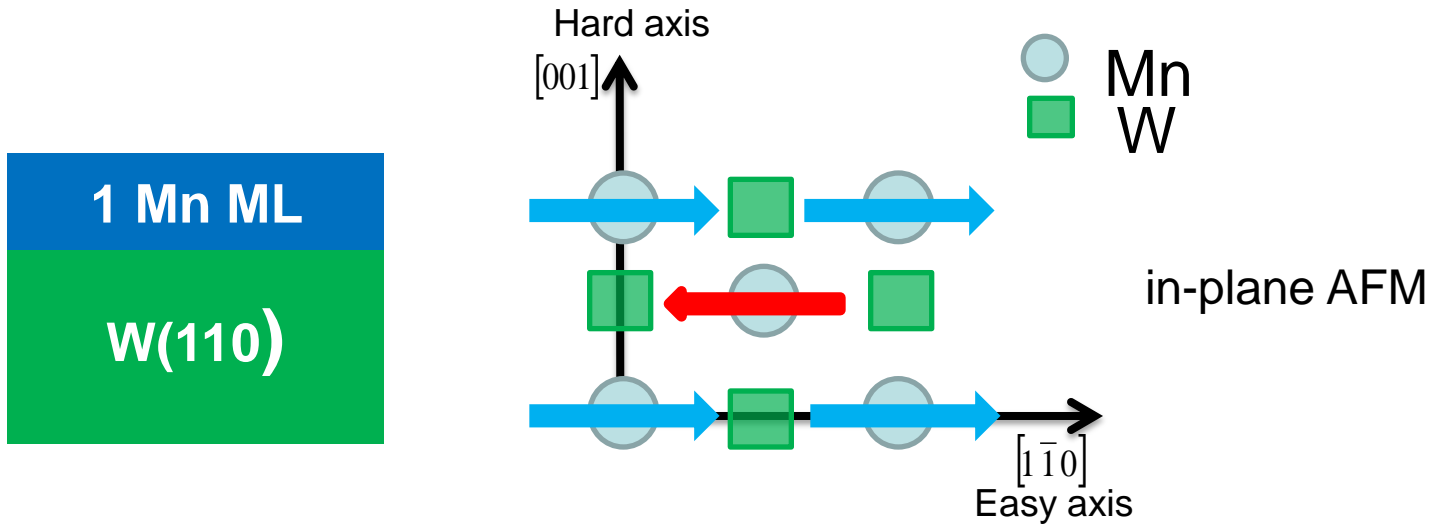
Formation of spin-spirals: spin-model for Mn/W(110)

L. Udvardi *et al.*, Physica B 403, 402-404 (2008)

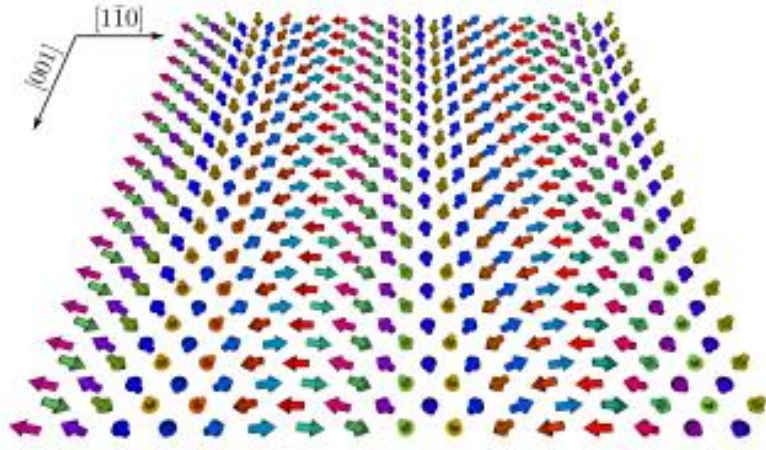
rel. torque method

G. Hasselberg *et al.* (2013) to be submitted

spin-cluster expansion



Formation of spin-spirals: spin-model for Mn/W(110)



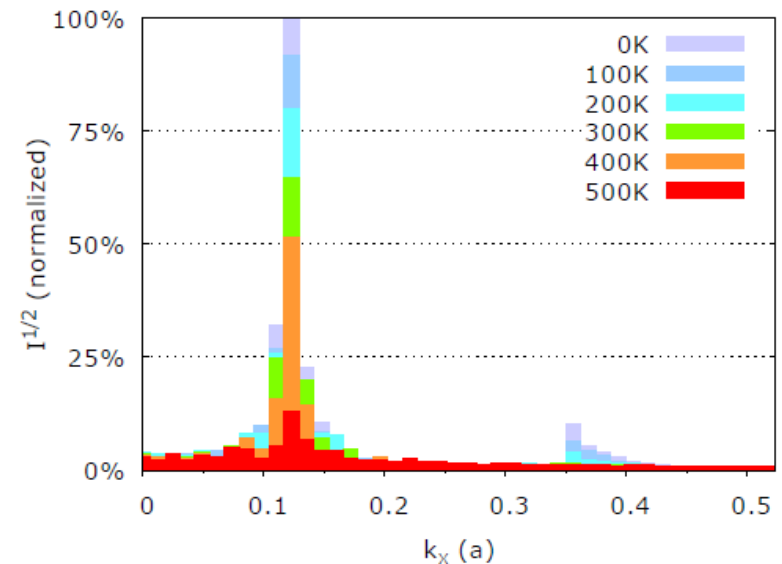
CSS ground state obtained by simulating a cooling-down process until $T=0$ K

Wave length $\lambda=12$ nm in good agreement with experiment.

Temperature dependence of λ

P. Sessi *et. al.*, Phys. Rev. Lett. 103, 167201 (2009)

Square root of the intensity of the Fourier transform of z-component of the magnetization plotted over k_x for different temperatures

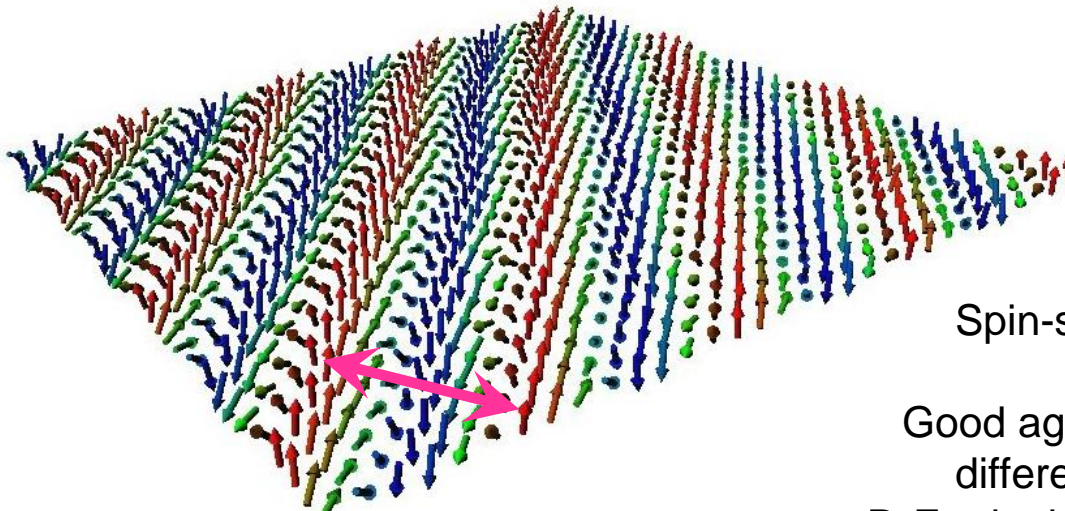
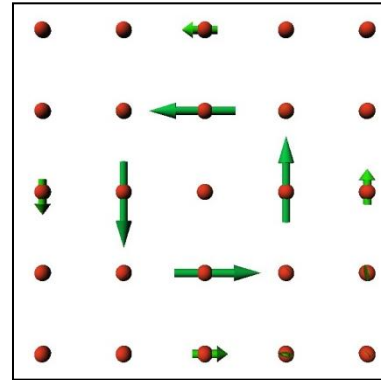


Formation of spin-spirals in Mn/W(100)

L. Udvardi *et al.*, Physica B 403, 402-404 (2008) rel. torque method

Calculated isotropic exchange interactions
and length of DM vectors (all data in mRyd)

NN	1	2	3
J_{ij}	-3.91	1.10	0.12
D_{ij}	0.57	0.04	0.24



Spin-spiral wavelength ~ 2.2 nm

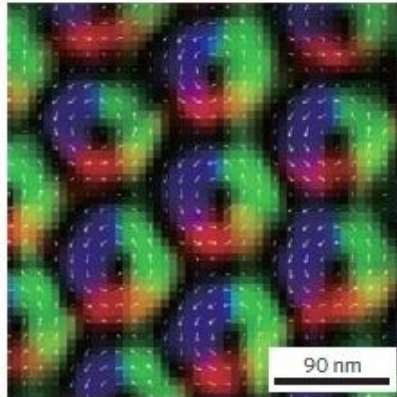
Good agreement with experiment and
different theoretical approach by
P. Ferriani *et al.*, PRL **101**, 027201 (2008)

Magnetic skyrmions & skyrmion lattices

Topologically protected whirling spin-configurations

$$S = \frac{1}{4\pi} \int \mathbf{n} \cdot \left(\frac{\partial \mathbf{n}}{\partial x} \times \frac{\partial \mathbf{n}}{\partial y} \right) dx dy \rightarrow 0, \pm 1$$

Bulk MnSi, FeCoSi, ...

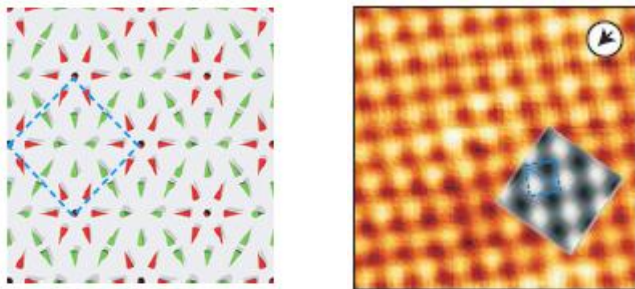


Surface skyrmions

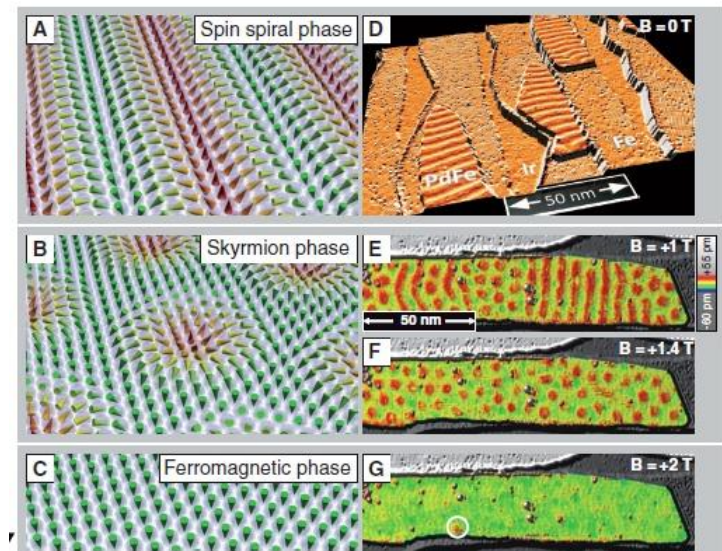


A. Fert *et al.*, Nature Nanotech. 8, 152 (2013):
promise for ultradense memory and logic devices

S. Heinze *et al.*, Nature Phys. 7, 713 (2011)
nanoskyrmions in Fe/Ir(1111)

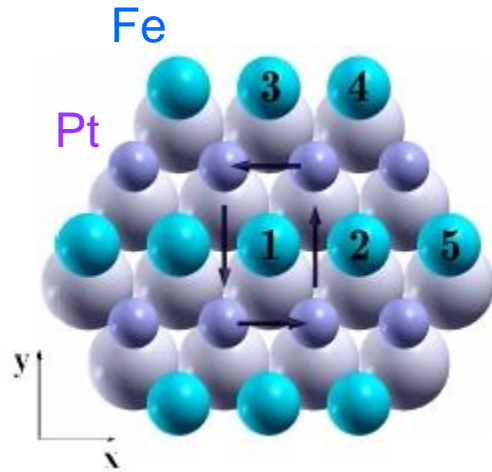


N. Romming *et al.*, Science 341, 636 (2013)
skyrmions in a PdFe bilayer on Ir(111)



Magnetic skyrmions & skyrmion lattices

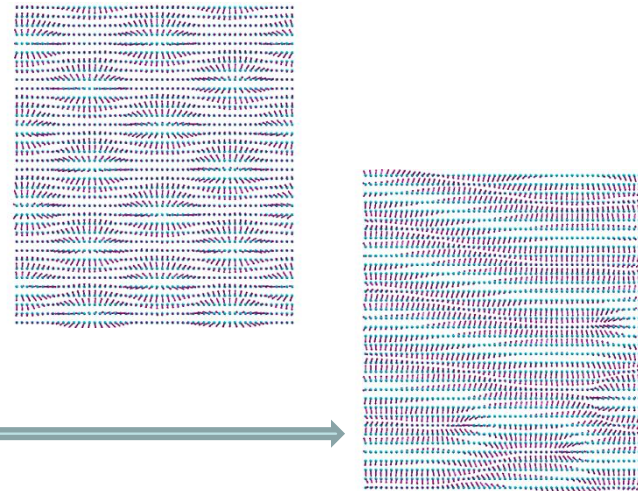
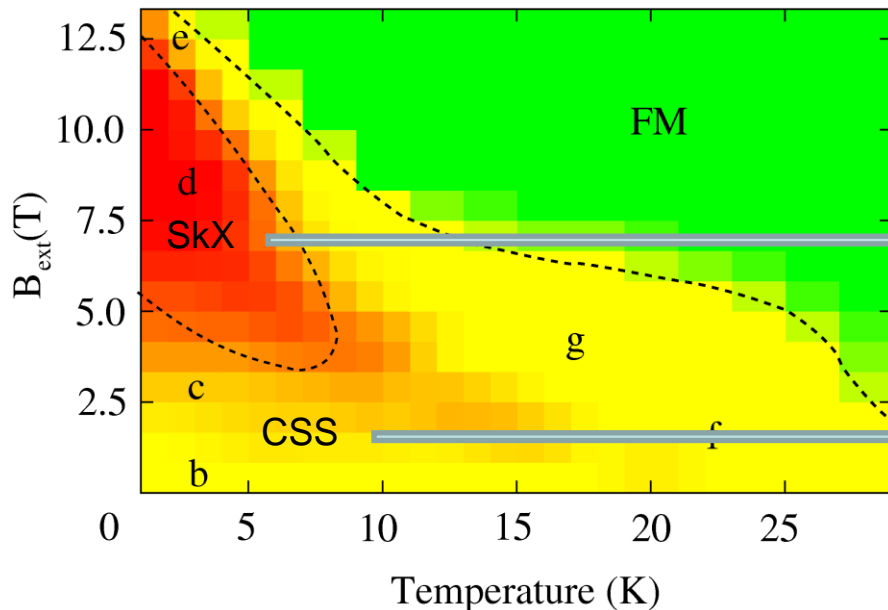
S. Polesya et al., arxiv:1310.5681 ab initio spin-model (RTM) & Monte-Carlo study of an FePt monolayer on Pt(111)



DM and isotropic exchange interactions

$i-j$	R_{ij}	D_{ij}^x	D_{ij}^y	D_{ij}^z	J_{ij}
1-2	0.707	0.00	2.44	0.39	23.90
1-3	1.225	-2.47	0.00	0.00	4.59
1-4	1.414	-0.31	0.50	-0.69	-0.56
1-5	1.414	0.00	0.39	-0.17	-1.07

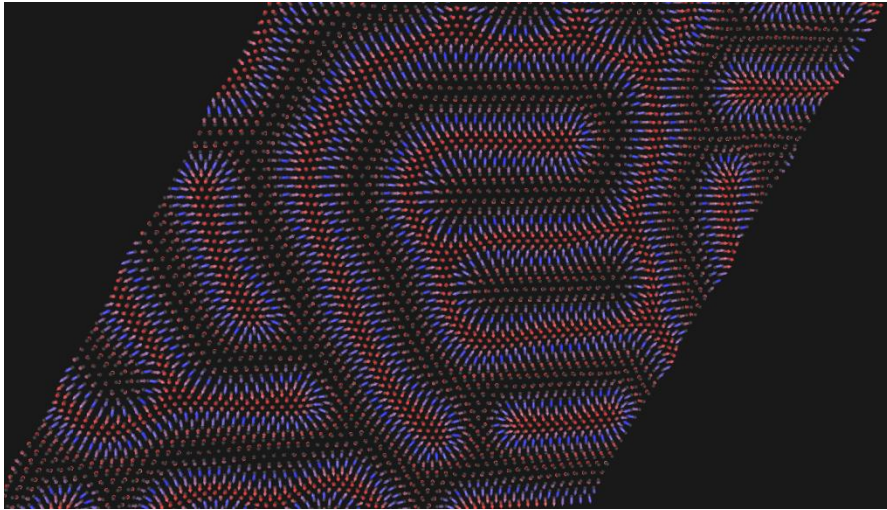
B-T phasediagram



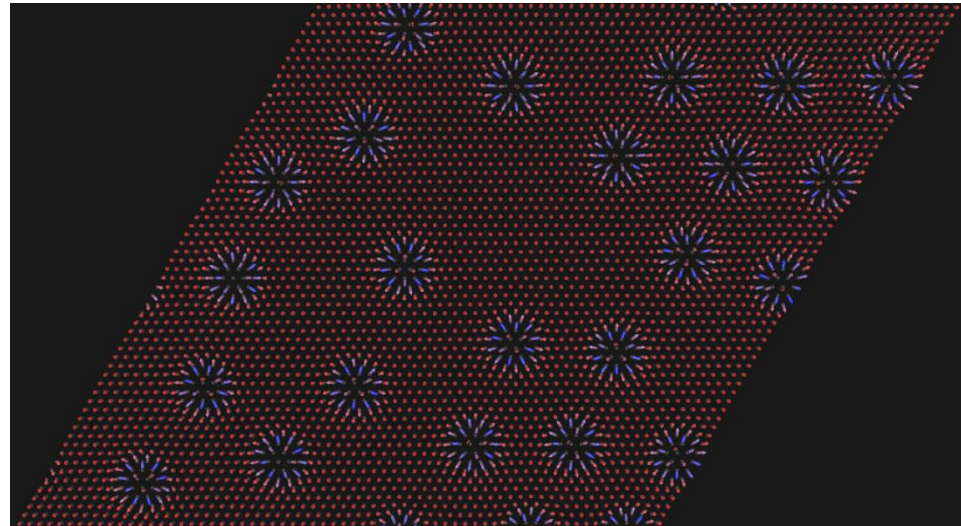
Magnetic skyrmions & skyrmion lattices

E. Simon et al. (unpublished) PdFe bilayer on Pt(111), SCE spin-model
($T = 0$ K)

$B = 0$ T \rightarrow CSS with $\lambda \approx 2.2$ nm

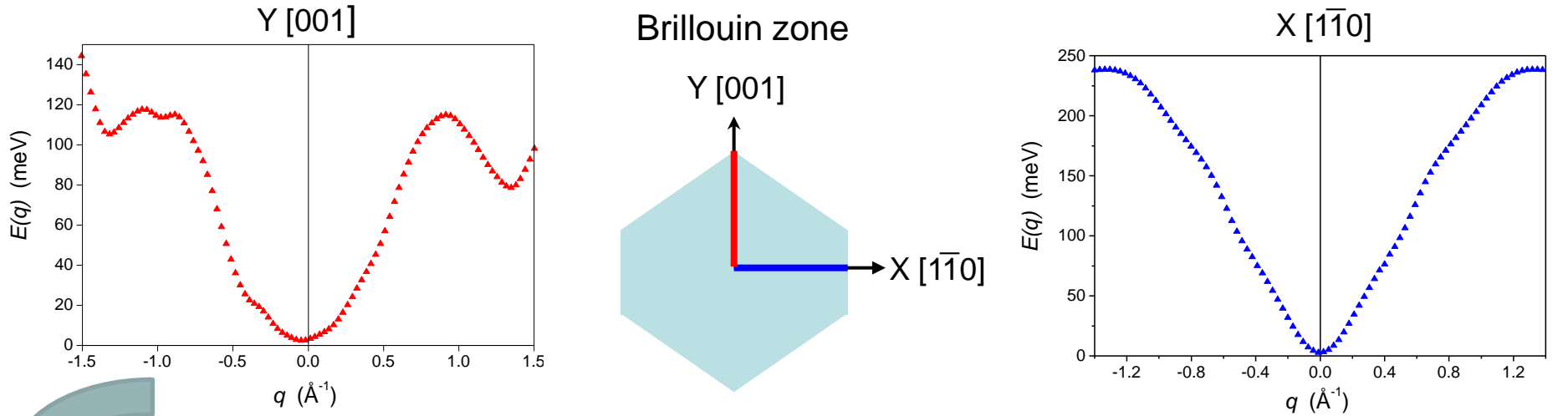


$B \approx 100$ T (???) \rightarrow SkX

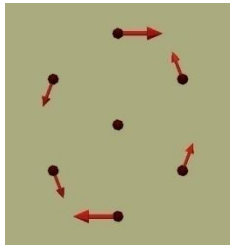
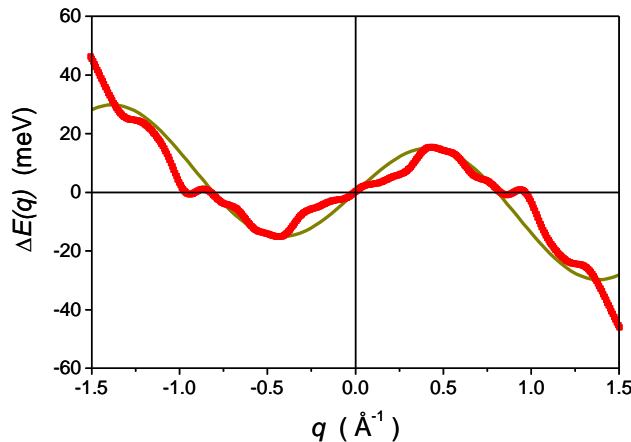


Chiral asymmetry of the spin-wave spectrum

Spin-waves in $\text{Fe}_1/\text{W}(110)$ L. Udvardi and L.Szunyogh, PRL 102, 207204 (2009)



Asymmetry



DM vectors

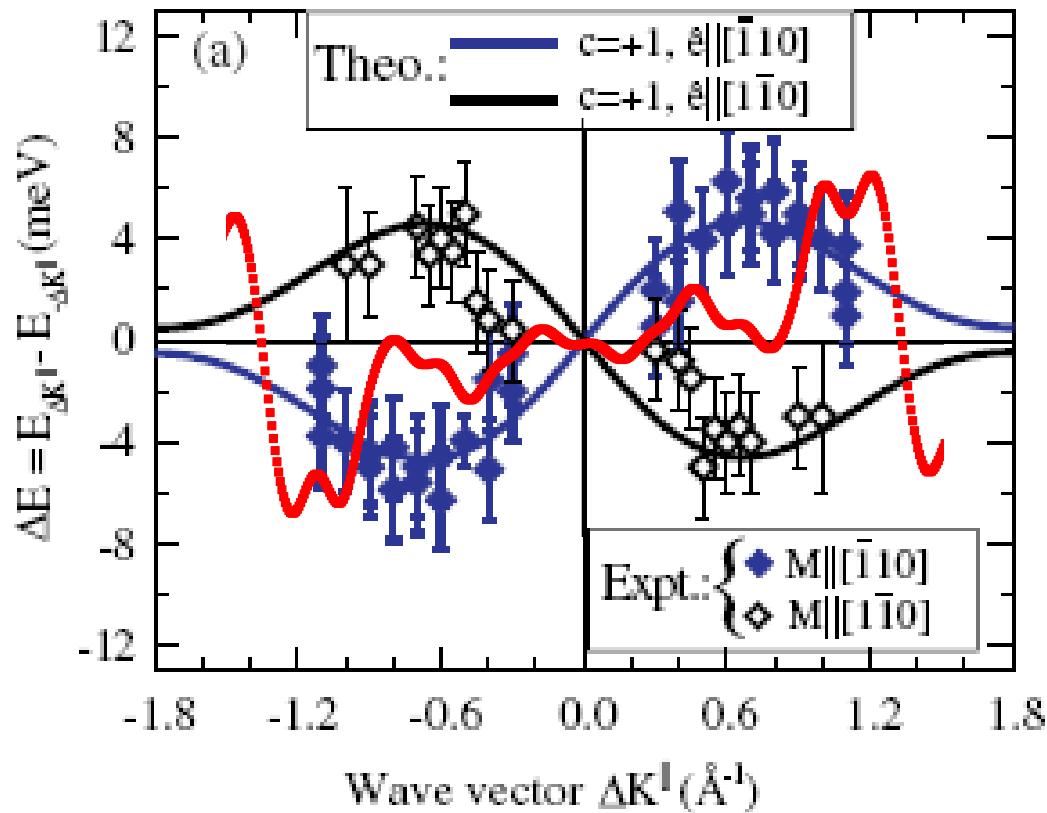
$$\Delta E(q) = \frac{16\mu_B}{M_0} D_1^x \sin\left(\frac{1}{2}qa\right) - \frac{8\mu_B}{M_0} D_2^x \sin(qa)$$

Possibility for a direct measurement of the DM interactions!

Chiral asymmetry of the spin-wave spectrum

Experiment: Fe₂/W(110)

Kh. Zakeri et al. PRL 104, 137203 (2010)



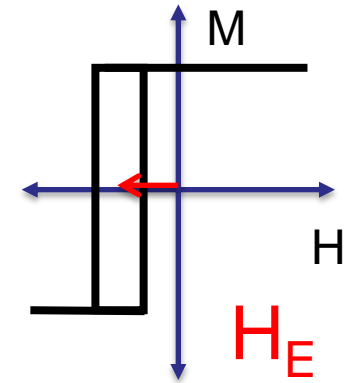
*L.U. & L.S. (unpublished)
SCE spin-model*

Exchange bias

Effective unidirectional anisotropy
at an FM/AFM interface



Shift in the
hysteresis loop

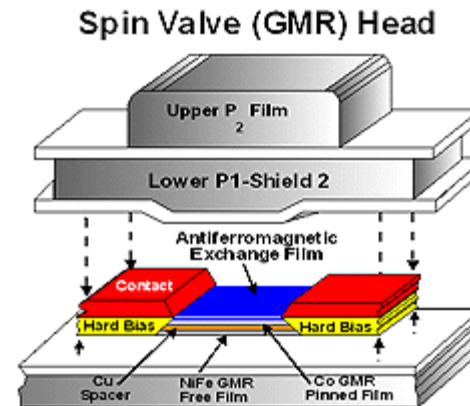
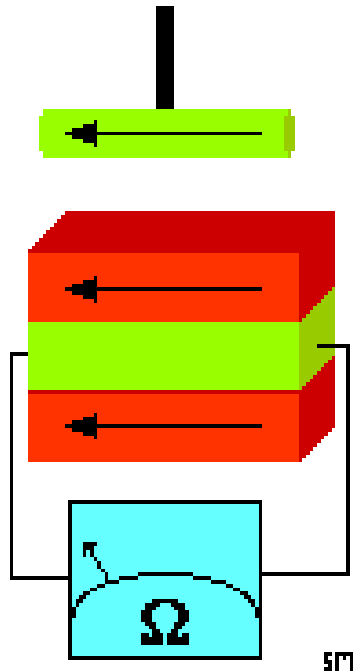


W.H. Meiklejohn *et.al.*, Phys. Rev. 105, 904 (1957)

A.P. Malozemoff, J. Appl. Phys. 63,3874 (1988)

T.C. Schulthess and W.H. Butler, Phys. Rev. Lett. 81,4516 (1998)

P. Miltényi *et.al.*, Phys. Rev. Lett, 84, 4224 (2000)



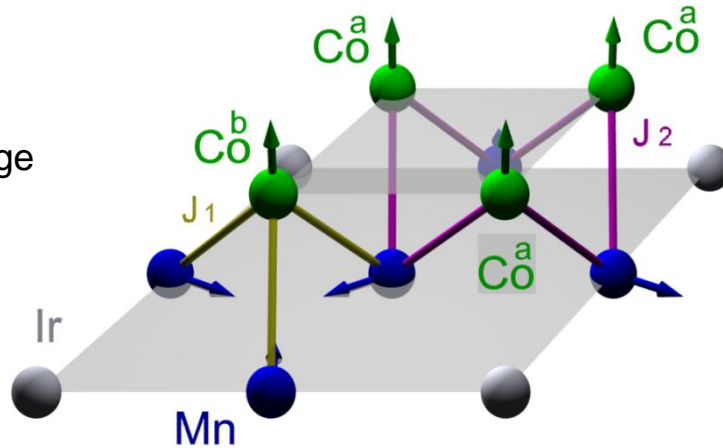
Perpendicular exchange bias (PEB) at an IrMn₃/Co interface

L. Szunyogh *et al.*, Phys. Rev. B 79, 020403(R) (2009)

L. Szunyogh *et al.*, Phys. Rev. B 83, 024401 (2011)

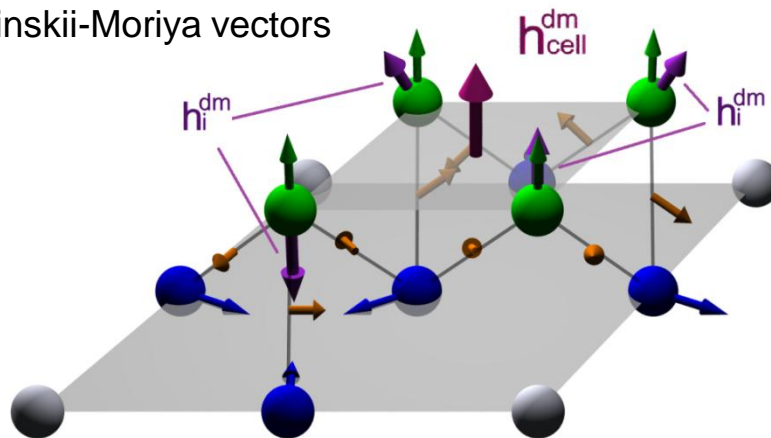
R. Yanes *et al.*, Phys. Rev. Lett. in print (2013)

Isotropic exchange



$J_1 = -6.15 \text{ meV}$ (AFM)
 $J_2 = 1.2 \text{ meV}$ (FM)

Dzyaloshinskii-Moriya vectors



$\text{DM}_{\text{Co-Mn}}$ vectors are in-plane



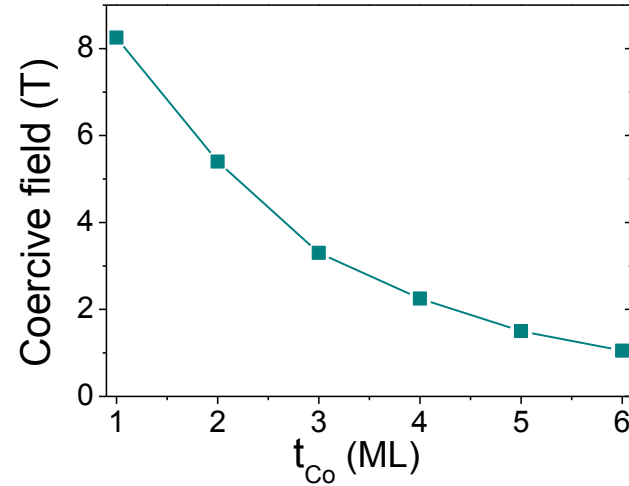
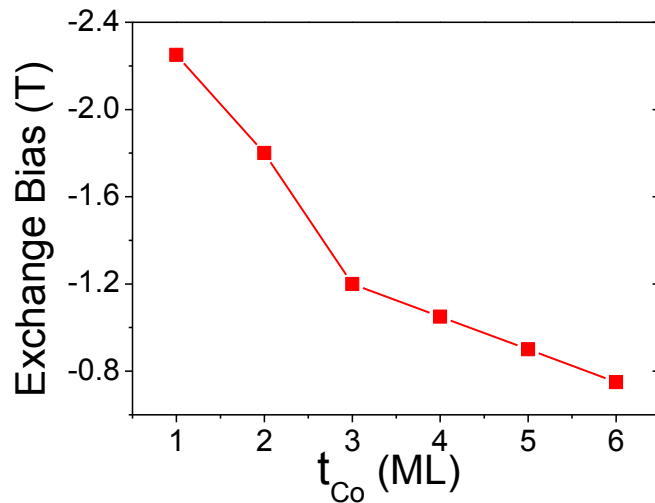
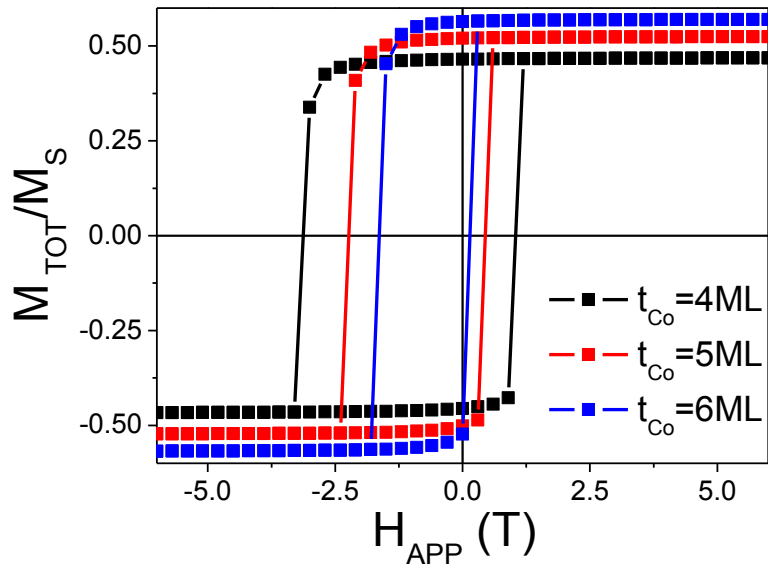
Effective DM field acting on the Co moments::

$$\vec{h}_{\text{cell}}^{\text{dm}} = \sum_{i=1}^4 \sum_j' (\vec{m}_j \times \vec{D}_{ij}) \approx 1.05 \hat{k} \text{ meV}$$

perpendicular to the interface

Perpendicular exchange bias (PEB) at IrMn₃/Co interface

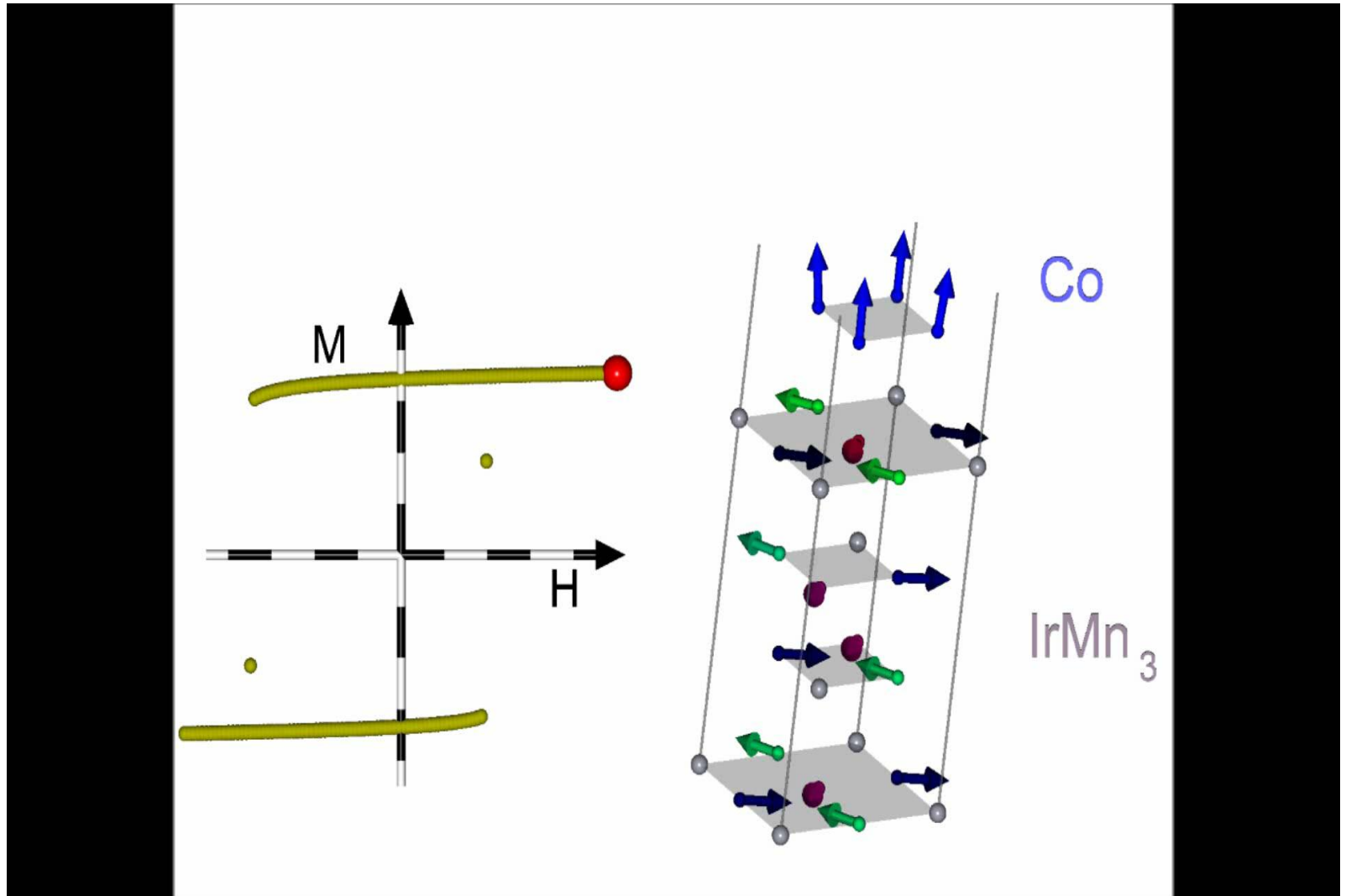
Hysteresis loops as a function of t_{Co}



Perpendicular exchange bias (PEB) at IrMn₃/Co interface

The switching process: spin-flop coupling

[T.C. Schulthess and W.H. Butler, Phys. Rev. Lett. 81,4516 (1998)]



Conclusions

- Relativistic (spin-orbit) effects play a pronounced role in nanomagnetism
- Dzyaloshinskii-Moriya interactions give rise to
 - spin-spirals
 - magnetic skyrmions
 - asymmetry of the spin-wave spectra
 - perpendicular exchange bias
 - *homochirality of domain walls, SS's, SK's*
 - *effects in finite magnetic nanoparticles ...*
- Spin-models are often useful, however, they (mostly) fail in the presence of induced moments → constrained DFT; RDLM at finite temperature

Thank you for attention!