# Challenges in Non-Commutative Information Geometry 

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## Introduction

My goal: Find alternative description of quantum information theory

- Standard: a density matrix $\rho$ describes a statistical mixture of quantum states
- Alternative: use the GNS-representation of mixed quantum states

Why? Find a more general theory

- not relying on the properties of the trace
- using elements of information geometry
- using elements of non-commutative geometry
J. Naudts, Quantum Statistical Manifolds, Entropy 20(6), 472 (2018);

Correction submitted

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## Standard theory

See for instance Dénes Petz, Quantum Information Theory and Quantum Statistics (Springer, 2008)

The density matrix $\rho$ is a complex $n$-by- $n$ matrix satisfying

$$
\rho \geq 0 \text { and } \operatorname{Tr} \rho=1
$$

A special role is played by the density matrix

$$
\rho_{0}=\frac{1}{n} \mathbb{I}=\frac{1}{n}\left(\begin{array}{ccccc}
1 & 0 & \cdots & 0 & 0 \\
0 & 1 & \cdots & 0 & 0 \\
\cdots & & & & \\
0 & 0 & \cdots & 0 & 1
\end{array}\right)
$$

The space of $n$-by- $n$ matrices forms a Hilbert space $\mathcal{H}^{\text {Hs }}$ for the Hilbert-Schmidt scalar product

$$
\langle A, B\rangle_{\text {Hs }}=\operatorname{Tr} A^{\dagger} B=n \operatorname{Tr} \rho_{0} A^{\dagger} B=n\left\langle A^{\dagger} B\right\rangle_{\rho_{0}}
$$

## The relative modular operator

Fix two strictly positive density matrices $\rho$ and $\sigma$
Note that in general $\rho$ and $\sigma$ do not commute
(Petz 86) introduced the relative modular operator $\Delta_{\rho, \sigma}$ on $\mathcal{H}^{\text {HS }}$ defined by $\quad \Delta_{\rho, \sigma} A=\rho A \sigma^{-1} \quad$ for all $A$

The relative entropy (Umegaki 1962, Araki 1976), defined by

$$
D(\sigma \| \rho)=\operatorname{Tr} \sigma(\log \sigma-\log \rho)
$$

can be written as $\quad D(\sigma \| \rho)=\left\langle\sigma^{1 / 2} \mid\left[\log \Delta_{\sigma, \rho},\right] \sigma^{1 / 2}\right\rangle_{\text {Hs }}$
Proof Write $\Delta_{\sigma, \rho}=L_{\rho} R_{\sigma^{-1}}$ where $L_{\rho} R_{\sigma^{-1}}=R_{\sigma^{-1}} L_{\rho}$

## The metric

From the relative entropy one derives Bogoliubov's scalar product, which can be written as a metric tensor

$$
\begin{aligned}
g_{\sigma, \tau}(\rho)= & \int_{0}^{1} \mathrm{~d} u \operatorname{Tr} \rho^{u}(\log \sigma-\log \rho) \rho^{1-u}(\log \tau-\log \rho) \\
& -D(\rho| | \sigma) D(\rho \| \mid \tau) \\
= & \int_{0}^{1} \mathrm{~d} \boldsymbol{u}\left\langle\rho^{1 / 2}\right|\left[\log \Delta_{\tau, \rho}\right] \rho^{u}\left[\log \Delta_{\sigma, \rho}\right] \rho^{-u}\left|\rho^{1 / 2}\right\rangle_{\mathrm{HS}} \\
& -\left\langle\rho^{1 / 2} \mid\left[\log \Delta_{\tau, \rho}\right] \rho^{1 / 2}\right\rangle_{\mathrm{HS}}\left\langle\left[\log \Delta_{\sigma, \rho}\right] \rho^{1 / 2}\right\rangle_{\mathrm{HS}} .
\end{aligned}
$$

$\rho, \sigma, \tau$ are strictly positive density matrices

## The metric cont'd

$\rho$ is a point $A$ in the manifold $\mathbb{M}$ of strictly positive density matrices

With $\sigma, \tau$ correspond two points $B, C$ in the tangent plane $T_{\rho} \mathbb{M}$

$g_{\sigma, \tau}(\rho)$ is the scalar product between the vectors $\overrightarrow{A B}, \overrightarrow{A C}$

## The GNS representation (Gelfand, Naimark, Segal)

$n \times n$ matrices represented as $n^{2} \times n^{2}$ block matrices
In this representation the density matrix $\rho$ is replaced by a 'wave function' $\Omega_{\rho}$.

- Diagonalize $\rho: \quad \rho=\sum p_{n}\left|\psi_{n}\right\rangle\left\langle\psi_{n}\right| \quad$ with $p_{n}>0, \quad \sum_{n} p_{n}=1$.
- Let $\quad \Omega_{\rho}=\sum_{n} \sqrt{p_{n}} \psi_{n} \otimes \psi_{n}$
- Then

$$
\begin{aligned}
\left\langle\Omega_{\rho}\right| \boldsymbol{A} \otimes \mathbb{I}\left|\Omega_{\rho}\right\rangle & =\sum_{m, n} \sqrt{p_{m} p_{n}}\left\langle\psi_{m} \otimes \psi_{m} \mid\left(\boldsymbol{A} \psi_{n}\right) \otimes \psi_{n}\right\rangle \\
& =\sum_{m, n} \sqrt{p_{m} p_{n}}\left\langle\psi_{m} \mid \boldsymbol{A} \psi_{n}\right\rangle\left\langle\psi_{m} \mid \psi_{n}\right\rangle \\
& =\sum_{n}^{n} p_{n}\left\langle\psi_{n} \mid \boldsymbol{A} \psi_{n}\right\rangle \\
& =\operatorname{Tr} \rho \boldsymbol{A}=\langle\boldsymbol{A}\rangle_{\rho} .
\end{aligned}
$$

## The GNS representation Cont'd

In particular, with the tracial density matrix $\rho_{0}=\frac{1}{n} \mathbb{I}$ corresponds

$$
\Omega_{0}=\frac{1}{n} \sum_{i=1}^{n} \psi_{i} \otimes \psi_{i}
$$

where $\psi_{1}, \cdots, \psi_{n}$ is any orthonormal bazis.

Let $\mathcal{A}$ denote the space of all 'operators' of the form $A \otimes \mathbb{I}$
The commutant $\mathcal{A}^{\prime}$ consists of all operators of the form $\mathbb{I} \otimes A$

## The chart centered at $\rho$

The metric tensor $g_{\sigma, \tau}(\rho)$ involves 3 density matrices $\rho$ fixes a point $A$ in the manifold $\mathbb{M}$, $\sigma, \tau$ fix the points B and C in the tangent plane $T_{\rho} \mathbb{M}$
$\rho$ is described by the 'wave function' $\Omega_{\rho}$ How to describe $\sigma$ and $\tau$ ?

Theorem Given $\rho, \sigma$ there exists a unique $K=K^{\dagger}$ in $\mathcal{A}^{\prime}$
such that $K \Omega_{\rho}=\int_{0}^{1} \mathrm{~d} u \rho^{u}\left[\log \Delta_{\sigma, \rho}\right] \rho^{-u} \Omega_{\rho}+D(\rho \| \sigma) \Omega_{\rho}$ and $\quad\left\langle\Omega_{\rho} \mid K \Omega_{\rho}\right\rangle=0$

The map $\sigma \mapsto K$ is a chart for $\mathbb{M}$
It is centered at $\rho$ : Indeed, $\sigma=\rho$ implies $K=0$

## The metric

Proposition Consider the chart $\chi_{\rho}: \sigma \mapsto K$ centered at $\rho$. Then

1) There exists a strictly positive operator $G_{\rho}$ in $\mathcal{A}$ such that

$$
G_{\rho} K \Omega_{\rho}=\left[\left(\Delta_{\rho, \sigma}+D(\rho \| \sigma)\right] \Omega_{\rho}\right.
$$

2) For each pair $K=\chi_{\rho}(\sigma)$ and $L=\chi_{\rho}(\tau)$ is

$$
g_{\sigma, \tau}(\rho)=\left(K \Omega_{\rho}, G_{\rho} L \Omega_{\rho}\right)
$$

Note: $G_{\rho}$ is in $\mathcal{A}$ while $K$ and $L$ belong to the commutant $\mathcal{A}^{\prime}$
Positivity of the metric follows immediately from $G_{\rho}>0$

## The exponential connection

The geodesics $t \mapsto \rho_{t}$ are such that

$$
\begin{aligned}
\log \rho_{t} & =(1-t) \log \rho_{0}+t \log \rho_{1}-\zeta(t) \\
& =\log \rho_{0}+t H-\zeta(t)
\end{aligned}
$$

$\zeta(t)$ is a normalizing function, $H$ is defined by $H=\log \rho_{1}-\log \rho_{0}$

## Proposition

1) $\frac{\mathrm{d} \zeta}{\mathrm{d} t}=\langle H\rangle_{t}=D\left(\rho_{t} \| \rho_{0}\right)-D\left(\rho_{t} \| \rho_{1}\right)$
2) $t \mapsto \zeta(t)$ is convex
3) $\chi_{\rho}\left(\rho_{t}\right)=(1-t) \chi_{\rho}\left(\rho_{0}\right)+t \chi_{\rho}\left(\rho_{1}\right)$

The latter shows that the chart $\chi_{\rho}$ is an affine coordinate

## Summary

- It is possible to eliminate references to the trace operation using the GNS representation
- Label $\sigma$ relative to $\rho$ with an operator $K=K^{\dagger}$ in $\mathcal{A}^{\prime}$ This gives a chart $K=\chi_{\rho}(\sigma)$ of $\mathbb{M}$, centered at $\rho$
- Scalar product of Bogoliubov $\Rightarrow G_{\rho}>0$ in $\mathcal{A}$
- $K=\chi_{\rho}(\sigma)$ is an affine coordinate for the exponential connection

Challenges How much of this can be generalized?

